

## Research paper

## Short-term and spike-timing-dependent plasticity facilitate the formation of modular neural networks



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## ABSTRACT

The brain has the phenomenal ability to reorganise itself by forming new connections among neurons and by pruning others. The so-called neural or brain plasticity facilitates the modification of brain structure and function over different time scales. Plasticity might occur due to external stimuli received from the environment, during recovery from brain injury, or due to modifications within the body and brain itself. In this paper, we study the combined effect of short-term (STP) and spike-timing-dependent plasticity (STDP) on the synaptic strength of excitatory coupled Hodgkin-Huxley neurons and show that plasticity can facilitate the formation of modular neural networks with complex topologies that resemble those of networks with preferential attachment properties. In particular, we use an STDP rule that alters the synaptic coupling intensity based on time intervals between spikes of postsynaptic and presynaptic neurons. Previous work has shown that STDP may induce the emergence of directed connections from high to low frequency spiking neurons. On the other hand, STP is attributed to the release of neurotransmitters in the synaptic cleft of neurons that alter its synaptic efficiency. Our results suggest that the combined effect of STP and STDP with long recovery times facilitates the formation of connections among neurons with similar spike frequencies only, a kind of preferential attachment. We then pursue this further and show that, when starting with all-to-all neural configurations, depending on the STP recovery time and distribution of neural frequencies, modular neural networks can emerge as a direct result of the combined effect of STP and STDP.

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## 1. Introduction

Mathematical models have been employed in neuroscience since early in the 20th century [1] to explain experimental findings and elucidate the inner workings of the brain. In 1907, Lapique [2] proposed the integrate-and-fire model that can

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reproduce the electrical activity of the membrane potential of neurons. It is one of the most popular models for studying the dynamic behaviour of neural systems. Later in 1952, Hodgkin and Huxley [3] explained the ionic mechanisms in the cell membrane of neurons and proposed the so-called Hodgkin-Huxley (HH) neural model that has, since then, been used extensively to study neural networks [4].

In particular, neural network models have been used extensively in computational neuroscience, such as in studies of neural dynamics [5], dynamic range [6–9], neural synchronisation [10–14], flow of information [15–17], and brain plasticity [18,19] to name a few.

Neural plasticity is the ability of the brain to modify its function and structure over different time scales [21,22]. The term was initially used by James [23] in 1890 to propose that phenomena of habit in living organisms are due to plasticity. Cajal [24] reported in the early 1900s his research about regenerative and degeneration changes in brain structure. In 1924, Lashley [25] demonstrated experimental evidence of a malleable brain. Konorski [26] and Hebb [27] proposed in 1948 and 1949, respectively, that neural activities influence the connection among neurons. In 1960, Bennet et al. [28] carried out experiments with rats and observed chemical and anatomical plasticity in the brain. Since then, there have been many theoretical [29] and empirical [30] studies aiming to explain and understand brain plasticity and its effects in brain structure and function.

In synaptic plasticity, the synapses among neurons are potentiated or depressed in time according to the activity of other neurons [31,32]. Recent works [33,34] have shown that short-term (STP) and spike-timing-dependent plasticity (STDP) are different forms of neural processes leading to synaptic modifications. In particular, STDP depends on the relative timing of presynaptic and postsynaptic neural spikes [35,36]. This type of plasticity can lead to various dynamical phenomena and coupling structures, such as stable localised structures [37], stimulation-induced synchronisation or desynchronisation [38], noise-enhanced synchronisation [39,40] and nontrivial topology [20]. The STDP mechanism plays a role in temporal coding of information by spikes [37,41]. On the other hand, STP is attributed to the release of neurotransmitters in the synaptic cleft of neurons that alter its synaptic efficiency and acts on shorter time scales, ranging from milliseconds to hundreds to thousands of milliseconds [42–45]. As in the case with STDP, STP influence significantly the dynamical behaviour in a network. For instance, it may stabilise the parametric working memory [46], contribute to the emergence of spontaneous travelling waves [47], or induce phase changes in neural postsynaptic spiking [33].

Here, we extend the work in Borges et al. [20] which was focused on STDP only and study the combined effect of STP and STDP in neural networks of excitatory coupled HH neurons. The plasticity terms that model STDP in the equations in Section 2 are based on the experimental results by Bi and Poo [48,49] that were performed on excitatory synapses and on theoretical results by Abbott et al. [43] and Popovych et al. [39] (STDP). The results in Bi and Poo [48,49] show that STDP is a function of the relative timing of postsynaptic and presynaptic spikes and is theoretically backed by the Hebbian synaptic learning rule [27]. Instead, STP depends on the neural recovery dynamics [50,51]. McDonnell and Graham [33] used mathematical analysis and numerical simulations to show that STP induces phase changes in neural postsynaptic spiking. In our work, we start by studying the simplest case of a pair of HH neurons for a range of spike frequencies, aiming to understand how connectivity between them changes by the combined effect of STP and STDP. Next, we build an initially, all-to-all (globally) connected network of HH neurons and consider the simultaneous effect of STP and STDP for a range of coupling strengths. We show that STP plays an important role in changes of connectivity in neural networks with STDP. Indeed, we find that for high STP recovery time, only neurons with similar spike frequencies tend to connect, a form of preferential attachment. More importantly, our results show that, when starting with all-to-all networks, depending on the STP recovery time and distribution of neural frequencies, modular neural networks can emerge as a direct result of the combined effect of STP and STDP, a structure depicted by neurophysiological and experimental studies [52,53]. For the considered setup, STP plays a balancing role: while STDP tends to synchronise all neurons in a cluster, STP destroys the strong synchronisation and leads to a modular structure.

The paper is organised as follows: In Section 2, we introduce the general mathematical model of HH neural networks with STP and STDP and, in Section 3, we present our analysis and results based on numerical simulations, that show the effects of both types plasticity, initially on a pair of neurons and then, on a network of 100 HH neurons. Finally, we present the conclusions of our work in the last section.

## 2. A Hodgkin-Huxley neural network with STP and STDP

In our work, we build use a neural network model of  $N$  HH neurons coupled with excitatory chemical synapses, equipped with STP and STDP rules based on experimental results in Bi and Poo [48,49] and on theoretical models proposed by Popovych et al. [39] (STDP) and McDonnell and Graham [33,43,50] (STP).

Specifically, the HH neural network model considered here is given by

$$\begin{aligned} C\dot{V}_i = & I_i - g_K n_i^4 (V_i - E_K) g_{Na} m_i^3 h_i (V_i - E_{Na}) \\ & - g_L (V_i - E_L) + (V_r - V_i) \sum_{j=1}^N \varepsilon_{ij} f_j D_j, \end{aligned} \quad (1)$$

$$\dot{n}_i = \alpha_{n_i}(V_i)(1 - n_i) - \beta_{n_i}(V_i)n_i, \quad (2)$$

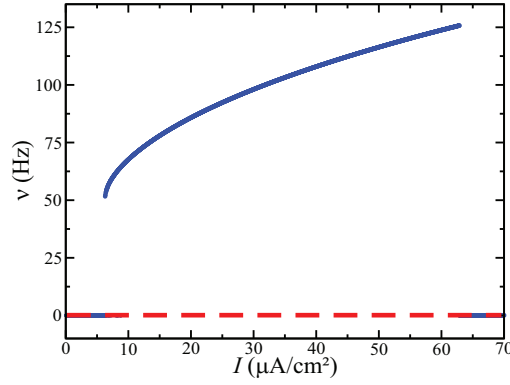


Fig. 1. Natural frequency  $\nu$  as a function of the current  $I$  for a single HH neuron.

$$\dot{m}_i = \alpha_{m_i}(V_i)(1 - m_i) - \beta_{m_i}(V_i)m_i, \quad (3)$$

$$\dot{h}_i = \alpha_{h_i}(V_i)(1 - h_i) - \beta_{h_i}(V_i)h_i, \quad (4)$$

$$\dot{f}_i = -\frac{f_i}{\tau_s}, \quad (5)$$

$$\dot{D}_i = \frac{1 - D_i}{\tau_D}, \quad (6)$$

where  $C$  (in  $\mu\text{F}/\text{cm}^2$ ) is the membrane capacitance and  $V_i$  (in mV) the membrane potential of neuron  $i$  at time  $t$ , with  $i = 1, \dots, N$ .  $I_i$  (in  $\mu\text{A}/\text{cm}^2$ ) is the constant current density of neuron  $i$  and  $\varepsilon_{ij}$  represents the matrix of coupling weights between neurons  $i$  and  $j$ . Parameters  $n_i$  and  $m_i$  are the activation of potassium and sodium functions, respectively, and  $h_i$  the inactivation of sodium function. Parameters  $g$  and  $E$  are associated with the conductance and reversal potential of each ion, respectively, and  $V_r$  is the excitatory reversal potential.

The various rate functions in Eqs. (1)–(6) are given by

$$\alpha_n(V) = \frac{0.01V + 0.55}{1 - \exp(-0.1V - 5.5)}, \quad (7)$$

$$\beta_n(V) = 0.125 \exp\left(\frac{-V - 65}{80}\right), \quad (8)$$

$$\alpha_m(V) = \frac{0.1V + 4}{1 - \exp(-0.1V - 4)}, \quad (9)$$

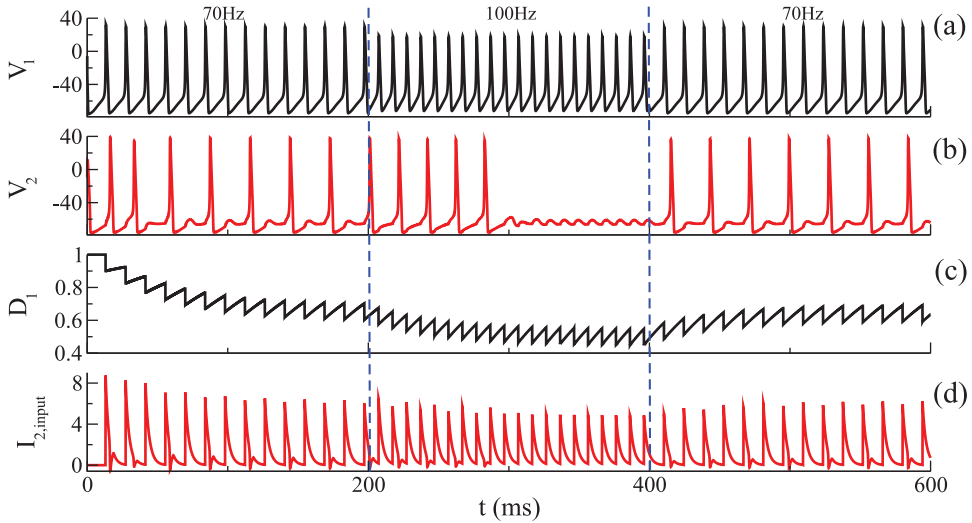
$$\beta_m(V) = 4 \exp\left(\frac{-V - 65}{18}\right), \quad (10)$$

$$\alpha_h(V) = 0.07 \exp\left(\frac{-V - 65}{20}\right), \quad (11)$$

$$\beta_h(V) = \frac{1}{1 + \exp(-0.1V - 3.5)}, \quad (12)$$

We consider  $C = 1 \mu\text{F}/\text{cm}^2$ ,  $g_K = 36 \text{ mS}/\text{cm}^2$ ,  $g_{Na} = 120 \text{ mS}/\text{cm}^2$ ,  $g_L = 0.3 \text{ mS}/\text{cm}^2$ ,  $E_K = -77 \text{ mV}$ ,  $E_{Na} = 50 \text{ mV}$ ,  $E_L = -54.4 \text{ mV}$  and  $V_r = 20 \text{ mV}$ .

Fig. 1 shows the spike (or natural) frequency  $\nu$  (Hz) of a single HH neuron as a function of current  $I$ . The spikes were numerically computed when the voltage  $V$  crosses the threshold of  $0\text{mV}$ , increasing from negative to positive values. In the simulations of 100 coupled HH neurons in Section 4, we consider  $I_i$  randomly distributed in the interval  $[10, 30] \mu\text{A}/\text{cm}^2$ ,



**Fig. 2.** The effect of STP on a pair of unidirectionally connected HH neurons, where neuron 1 is connected to neuron 2 but not vice versa: Temporal evolution of (a)  $V_1$ , (b)  $V_2$ , (c)  $D_1$  and (d)  $I_{2,\text{input}} = (V_r - V_2)\varepsilon_{21}f_1D_1$  with  $\varepsilon_{12} = 0$ ,  $\varepsilon_{21} = 0.1$ ,  $I_2 = 0$  and  $\tau_D = 50$  ms. Note that, for  $\nu_1 = 70$  Hz and  $\nu_1 = 100$  Hz, we used  $I_1 = 10.97 \mu\text{A}/\text{cm}^2$  and  $I_1 = 31.8 \mu\text{A}/\text{cm}^2$ , respectively.

leading to spike frequencies  $\nu_i$  in the interval [70,100] Hz. This choice of interval also allows for the monotonic increase of the natural frequencies  $\nu_i$  without reaching currents  $I$  bigger than  $60 \mu\text{A}/\text{cm}^2$  that correspond to a non-spike regime.

In Eq. (5),  $f_i$  is the strength of the effective synaptic output current from neuron  $i$  to neuron  $j$  and  $\tau_s$  (in ms) the synaptic time constant, fixed at  $\tau_s = 2.728$  ms. When neuron  $i$  spikes,  $f_i$  is updated ( $f_i \rightarrow 1$ ) before it starts to decay exponentially.

Eq. (6) models STP [33,50] with  $\tau_D$  (in ms) being the recovery time constant, related to biological mechanisms such as the depletion of release-ready neurotransmitter vesicles at the presynaptic terminal [44,51]. We assume that every time neuron  $i$  spikes, the update rule  $D_i \rightarrow D_i - d$  is applied. Biologically,  $D_i$  could represent the vesicles that can be used to transmit a signal from the presynaptic to the postsynaptic neuron. Therefore, the update rule means that the amount  $D_i$  of available vesicles is decreased by  $d = 0.1$  at every spike of neuron  $i$ , and then it recovers according to Eq. (6). We verify that different values of  $d$  lead qualitatively to the same behaviour. In that framework,  $D_i$  lies in [0,1] as if it happens  $D_i$  to be negative, it means that the neuron used all stored vesicles and when this happens,  $D_i$  is reset to 0. On the other hand, when  $D_i$  is equal to 1, it means that all neurotransmitter vesicles are restored.

Fig. 2 shows the effect of STP on a pair of neurons coupled with a unidirectional connection from neuron 1 to neuron 2. This is implemented by fixing  $\varepsilon_{12} = 0$  (i.e. the connectivity strength from neuron 2 to neuron 1 is 0) and  $\varepsilon_{21} = 0.1$  for the connectivity strength from neuron 1 to 2. In this study, we have set the STP recovery time  $\tau_D$  at 50 ms and  $I_2$  at 0, so that neuron 2 spikes only when it receives a strong enough input  $I_{2,\text{input}} = (V_r - V_2)\varepsilon_{21}f_1D_1$  from neuron 1. When the spiking frequency  $\nu_1$  in Fig. 2(a) changes from 70 Hz to 100 Hz,  $\nu_2$  in Fig. 2(b) exhibits a delayed alteration in its dynamic behaviour as its amplitude plummets at about 300 ms. We investigate further this phenomenon through the temporal evolution of  $D_1$  in Fig. 2(c) and the input current  $I_{2,\text{input}}$  received by neuron 2 in Fig. 2(d). Comparing the 70 Hz-regime with the 100 Hz-regime in Fig. 2(c) and (d), one can see that  $D_1$  decreases with the increase of the spike-frequency of neuron 1 to 100 Hz, and consequently  $I_{2,\text{input}}$  becomes less intense as it is not strong enough to cause spikes in the activity of neuron 2. When neuron 1 returns to 70 Hz spike frequency, there is more time for  $D_1$  to recover, thereby to increase the intensity of  $I_{2,\text{input}}$ , which triggers again spikes in the activity of neuron 2. Generalising this, one might say that STP increases the variability of postsynaptic responses to presynaptic frequency changes.

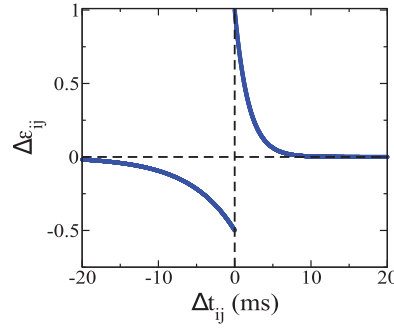
Moving now to the other form of plasticity introduced to the model, STDP gives rise to changes in the synaptic strength by means of the update function [39,48]

$$\varepsilon_{ij} \rightarrow \varepsilon_{ij} + 10^{-3} \Delta \varepsilon_{ij}, \quad (13)$$

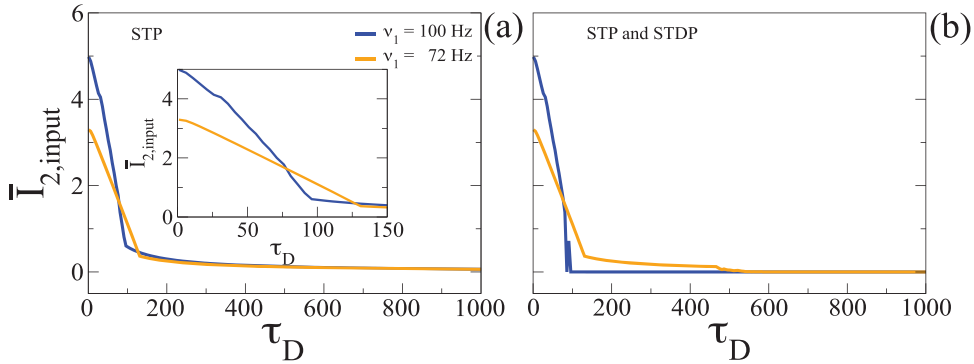
where

$$\Delta \varepsilon_{ij} = \begin{cases} \varepsilon_+ = A_1 e^{(-\Delta t_{ij}/\tau_1)}, & \text{if } \Delta t_{ij} > 0 \\ \varepsilon_- = -A_2 e^{(\Delta t_{ij}/\tau_2)}, & \text{if } \Delta t_{ij} < 0. \\ 0, & \text{if } \Delta t_{ij} = 0 \end{cases}$$

Here,  $\Delta t_{ij} = t_i - t_j$  is the time difference between the spike time  $t_i$  of the postsynaptic neuron  $i$  and the time  $t_j$  of the presynaptic neuron  $j$ , where  $t_i > t_j$ , leading to  $\Delta t_{ij} > 0$ . Fig. 3 shows the plot of the plasticity function  $\Delta \varepsilon_{ij}$  calculated from Eq. (14) for  $A_1 = 1$ ,  $A_2 = 0.5$ ,  $\tau_1 = 1.8$  ms,  $\tau_2 = 6$  ms, and  $\Delta t_{ij}$  varying from  $-20$  ms to  $20$  ms. This update rule is applied every time the postsynaptic neuron  $i$  spikes.



**Fig. 3.** Plot of  $\Delta\epsilon_{ij}$  that models STDP as a function of the difference between the spike times of the postsynaptic neuron  $i$  and presynaptic neuron  $j$ . Note the discontinuity at  $\Delta t_{ij} = 0$  ms.



**Fig. 4.** The average input current  $\bar{I}_{2,\text{input}}$  neuron 2 receives from neuron 1 as a function of the recovery time  $\tau_D$  for the following cases: (a) only STP and (b) STP and STDP. Neuron 2 has a fixed spike frequency  $\nu_2 = 70$  Hz ( $I_2 = 10.97 \mu\text{A}/\text{cm}^2$ ). The blue curve represents the case where  $\nu_1 = 100$  Hz ( $I_1 = 31.8 \mu\text{A}/\text{cm}^2$ ) and the orange curve, the case where  $\nu_1 = 72$  Hz ( $I_1 = 11.88 \mu\text{A}/\text{cm}^2$ ). We use  $\epsilon_{21} = 0.3$  and  $\epsilon_{12} = 0$ , to implement the case where  $\epsilon_{21}$  changes when STDP is considered. Note that the simulations run for  $200 \times 10^3$  ms and that  $\bar{I}_{2,\text{input}}$  was calculated over the last  $10 \times 10^3$  ms.

To understand better the effects of the recovery time  $\tau_D$ , we consider in the next the case of two, unidirectionally connected neurons where neuron 1 (presynaptic) is connected to neuron 2 (postsynaptic), but not vice versa. We evaluate how the average input current  $\bar{I}_{2,\text{input}}$  change with the increase of the recovery time  $\tau_D$  by comparing cases where the neurons have similar and dissimilar spike frequencies.

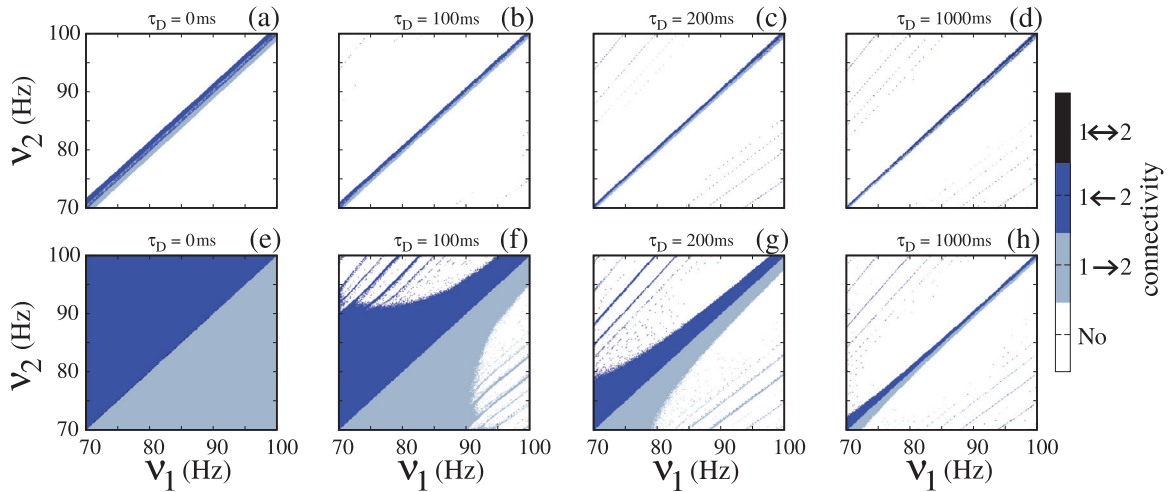
Specifically, we consider the cases where the system is equipped with only STP (see Fig. 4(a)) and, STP and STDP (see Fig. 4(b)). In both cases, we consider  $\epsilon_{21} = 0.3$  and  $\epsilon_{12} = 0$ , but for the case of STP and STDP (Fig. 4(b)), STDP acts only on  $\epsilon_{21}$  (unidirectional connection) ensuring that the spike frequency of presynaptic neuron 1 remains unchanged. The spike frequency of neuron 2 was fixed at  $\nu_2 = 70$  Hz ( $I_2 = 10.97 \mu\text{A}/\text{cm}^2$ ) and, for the blue curves in Fig. 4, neuron 1 has  $\nu_1 = 100$  Hz ( $I_1 = 31.8 \mu\text{A}/\text{cm}^2$ ) and for the orange curves,  $\nu_1 = 72$  Hz ( $I_1 = 11.88 \mu\text{A}/\text{cm}^2$ ). For the calculation of  $\bar{I}_{2,\text{input}}$ , we consider the last  $10 \times 10^3$  ms of the simulations.

Observing Fig. 4(a) (STP), one can see that for recovery times  $\tau_D < 75$  ms, the average current  $\bar{I}_{2,\text{input}}$  for  $\nu_1 = 100$  Hz (blue curve) is bigger than  $\bar{I}_{2,\text{input}}$  for  $\nu_1 = 72$  Hz (orange curve) (see also the inset in Fig. 4(a)). Interestingly, this change in the region  $75 \text{ ms} < \tau_D < 125$  ms where  $\bar{I}_{2,\text{input}}$  for  $\nu_1 = 72$  Hz (orange curve) is bigger than  $\bar{I}_{2,\text{input}}$  for  $\nu_1 = 100$  Hz (blue curve). For  $\tau_D > 125$  ms, both  $\bar{I}_{2,\text{input}}$  curves assume similar values and settle asymptotically to zero with further increasing in  $\tau_D$ . In Fig. 4(b) (STP and STDP), we see that the action of STP and STDP causes  $\bar{I}_{2,\text{input}}$  for  $\nu_1 = 100$  Hz to drop to 0 at  $\tau_D \approx 100$  ms, whereas  $\bar{I}_{2,\text{input}}$  for  $\nu_1 = 72$  Hz (orange curve) stays positive until  $\tau_D \approx 480$  ms.

These results show that STP make the influence of slower (or close frequency) neurons to become greater than that of faster (or very different frequency) neurons as  $\tau_D$  increases. In conjunction with STDP, this effect occurs due to the decrease of  $\bar{I}_{2,\text{input}}$  that makes the firing times uncorrelated (less synchronised), which then causes the coupling to disappear, i.e.  $\epsilon_{21} \rightarrow 0$ . This decoupling process tends to occur for smaller  $\tau_D$  values as the difference between  $\nu_1$  and  $\nu_2$  is amplified. We address this further later when we discuss the results in Fig. 5.

### 3. Two HH neurons case

Neural networks with STDP and random synaptic input were studied by Popovych et al. [39]. The authors reported that the mean synaptic coupling depends on noise intensity. Recently, the authors in [20] showed that STDP induces non-trivial



**Fig. 5.** Parameter spaces  $v_1 \times v_2$  for a pair of initially uncoupled neurons for recovery times (a)  $\tau_D = 0$  ms, (b)  $\tau_D = 100$  ms, (c)  $\tau_D = 200$  ms and (d)  $\tau_D = 1000$  ms, and for a pair of initially coupled neurons for (e)  $\tau_D = 0$  ms, (f)  $\tau_D = 100$  ms, (g)  $\tau_D = 200$  ms and (h)  $\tau_D = 1000$  ms. Note that  $v_1$  and  $v_2$  vary in  $[70, 100]$  Hz and that the colour bar shows the direction of synaptic connectivity, where white accounts for the uncoupled case (denoted “No”), grey for the case where neuron 1 is connected to neuron 2 (i.e.  $1 \rightarrow 2$ ), blue for the case where neuron 2 is connected to neuron 1 (i.e.  $1 \leftarrow 2$ ) and black for the bidirectional connection (i.e.  $1 \leftrightarrow 2$ ).

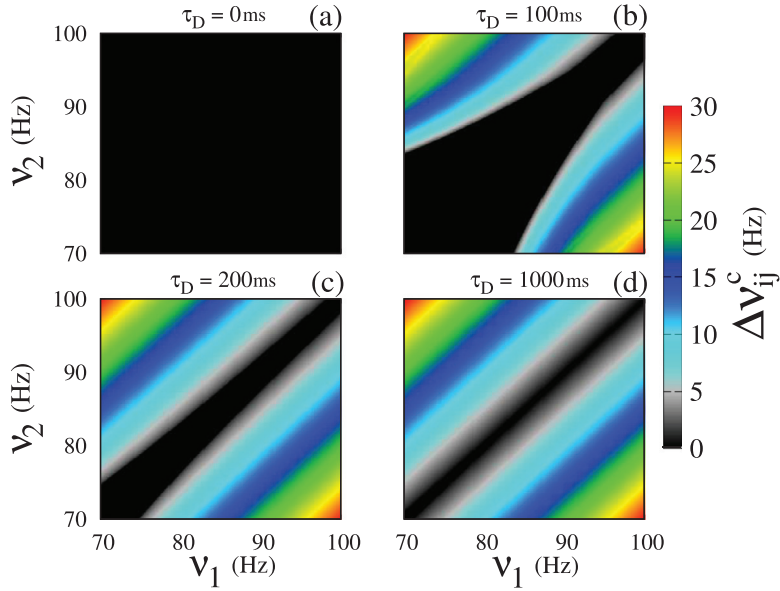
topology in neural networks. Here, we extend this work and build a neural network of  $N = 100$  HH neurons to study the combined effect of STP and STDP on the structure of the network and in particular, on its connectivity.

We start by analysing neural connectivity under the effect of STP and STDP. To this end, we consider a pair of HH neurons with STP and STDP, with the coupling strengths  $\varepsilon_{12}$  and  $\varepsilon_{21}$  varying in  $[0, 0.3]$ , and the connectivity threshold set at 0.01. This threshold is chosen because the coupling weights that should vanish in time actually oscillate around 0 assuming very small values as STDP is always present and thus, affecting them. As a consequence, the time averages of the coupling weights are not 0 but very close to 0. Moreover, for coupling weights smaller than this threshold, we notice that neurons influence each other only slightly.

Initially, the pair of neurons is uncoupled (i.e.  $\varepsilon_{12} = \varepsilon_{21} = 0$ ) or bidirectionally coupled with  $\varepsilon_{12} = \varepsilon_{21} = 0.3$ . Fig. 5 shows the direction of connectivity, after a transient time, for different natural frequencies  $v_1$  and  $v_2$  in  $[70, 100]$  Hz. The direction is coloured according to the connection from the output of one neuron to the input to the other neuron (directions of arrows in the colour bar): white accounts for the uncoupled case denoted in the colour bar as “No” ( $\bar{\varepsilon}_{ij} < 0.01$ ), grey for the case where neuron 1 is connected to neuron 2 (i.e.  $1 \rightarrow 2$ ), blue for the case where neuron 2 is connected to neuron 1 (i.e.  $1 \leftarrow 2$ ), and black for the bidirectional connection (i.e.  $1 \leftrightarrow 2$ ). For initially uncoupled neurons, the connections remain only for neurons with very similar frequencies, from those with faster spike frequencies to those with slower, as depicted in Fig. 5(a) ( $\tau_D = 0$  ms), Fig. 5(b) ( $\tau_D = 100$  ms), Fig. 5(c) ( $\tau_D = 200$  ms) and Fig. 5(d) ( $\tau_D = 1000$  ms). We note the absence of bidirectional connections and that the increase of the recovery time  $\tau_D$  gives rise to a narrower region of directed connections. With regard to neurons starting with bidirectional connections, we observe that for  $\tau_D = 0$  ms, neurons still remain unidirectionally connected (for all natural frequencies) from the high to the low frequency neurons (see Fig. 5(e)). Increasing  $\tau_D$  to 100 ms (see Fig. 5(f)), 200 ms (Fig. 5(g)) and 1000 ms (Fig. 5(h)), we observe a decrease in the area that represents connectivity, implying that the region of high spike frequencies is more affected by the influence of STP. Therefore, STDP makes the connections increase from faster spiking to slower spiking neurons and STP decreases the influence of high frequency neurons, allowing connections only for those neurons with similar spike frequencies.

The reason for the initially, bidirectionally coupled system to have larger areas of connectivity is related to the difference in frequency  $\Delta v_{ij}^c$  of neurons when coupled (or the synchronisation level). We note that this difference is not equal to the difference between their natural frequencies  $v_i - v_j$ , i.e.  $\Delta v_{ij}^c \neq v_i - v_j$ . In Fig. 6, we present the calculation of the average  $\Delta v_{ij}^c$  (colour bar) for the first  $t = 4000$  ms for the system with STP only. We consider  $N = 2$ ,  $\varepsilon_{12} = \varepsilon_{21} = 0.3$  and vary  $v_1$  and  $v_2$  in  $[70, 100]$  Hz. Fig. 6(a) shows the result for  $\tau_D = 0$  ms where, despite the natural frequencies of the neurons,  $\Delta v_{ij}^c$  is approximately equal to 0 (black region). When comparing these results with the results in Fig. 5(e), once can see that this corresponds to a connected region. For increasingly bigger recovery times  $\tau_D$  (see Fig. 6(b) for  $\tau_D = 100$  ms, Fig. 6(c) for  $\tau_D = 200$  ms and Fig. 6(d) for  $\tau_D = 1000$  ms), one can observe a decrease in the size of the black area, similar to the size of the area seen in Fig. 5(f), (g) and (h). Thus, one can infer that what defines a connected configuration in the case of STDP is the difference in spike frequency among neurons or how synchronised they become. The non-homogeneous distribution of connected areas in relation to the main diagonal in Fig. 5 occurs due to the non-linear variation of neural frequencies in relation to the received external current. In Fig. 1, we see that for low currents (i.e.  $I \in [10, 25]$ ), a small increase in  $I$  can cause a bigger variation in the frequencies when compared to the interval where  $I > 30$ . Therefore, low frequency





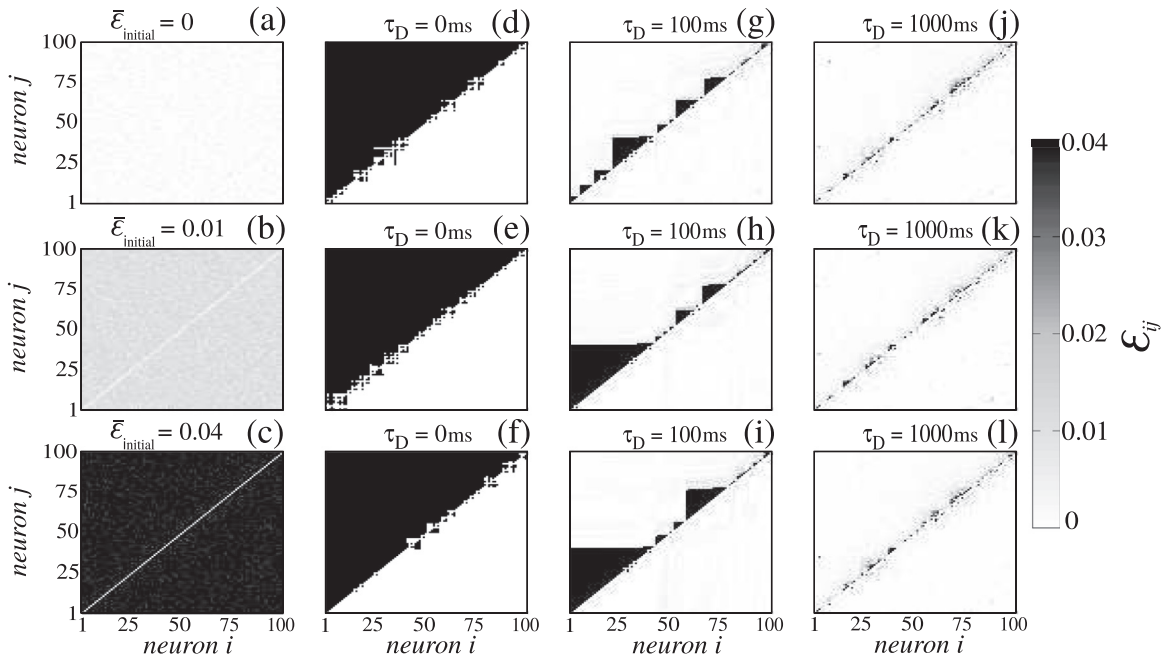
**Fig. 6.** Parameter space  $v_1 \times v_2$  for increasing values of  $\tau_D$ : (a) for  $\tau_D = 0$  ms, (b) for  $\tau_D = 100$  ms, (c) for  $\tau_D = 200$  ms and (d) for  $\tau_D = 1000$  ms. Note that  $v_1$  and  $v_2$  vary in [70,100] Hz and the colour bar represents  $\Delta v_{ij}^c$  values.

neurons are more sensitive to changes in their external currents, which facilitates their synchronisation with neurons with similar frequencies. The more synchronised they become, the more connected they remain. Coming back to the effect of STP, we find that it decreases the influence of faster neurons on slower ones, and consequently leads to the increase of their frequency differences and to the suppress of their synchronisation. That then leads to the depression of the connectivity via STDP.

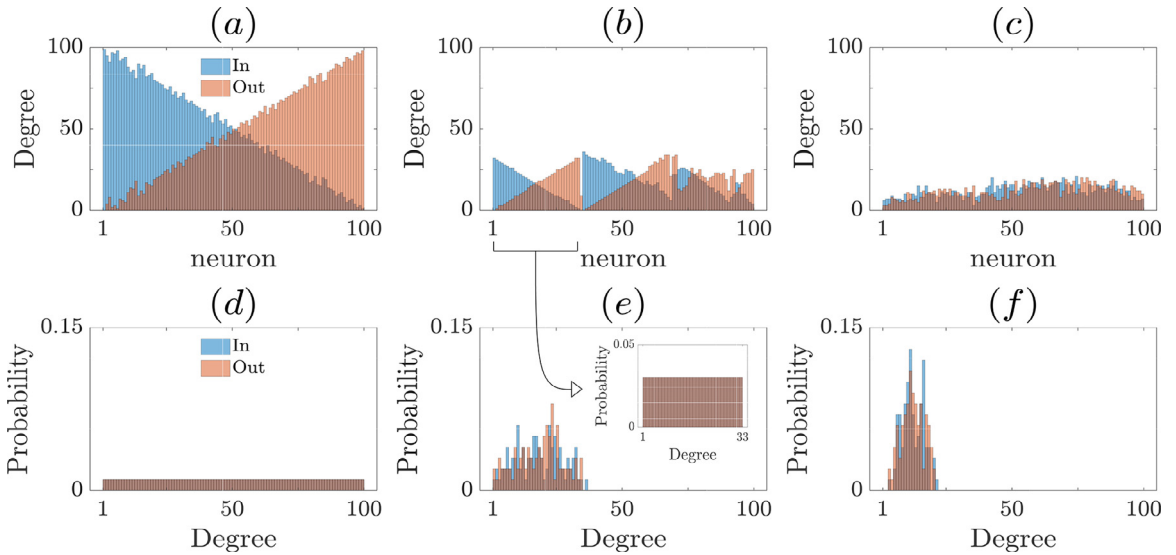
#### 4. Effects of STP on HH neural networks with STDP

We performed a similar study for a HH neural network and analysed how the connections evolve with the combined application of STP and STDP. In particular, we start with an all-to-all (globally connected) network of  $N = 100$  excitatory HH neurons. The coupling weights vary in  $[0,0.04]$  and  $I_i$  is randomly distributed so that the natural frequencies are in the range [70,100] Hz. At each time step, we check which neuron spikes and update its couplings according to the plasticity rule. We choose three initial coupling weight averages  $\bar{\varepsilon}_{\text{initial}}$ . The coupling matrices for the connectivity strengths  $\varepsilon_{ij}$  are shown in Fig. 7, where  $\bar{\varepsilon}_{\text{initial}} = 0$  in Fig. 7(a),  $\bar{\varepsilon}_{\text{initial}} = 0.01$  in Fig. 7(b) and  $\bar{\varepsilon}_{\text{initial}} = 0.04$  in Fig. 7(c). For all initial conditions studied, we kept the standard deviation of  $\bar{\varepsilon}_{\text{initial}}$  fixed at 0.002. In the coupling matrix  $(\varepsilon_{ij})$ , the presynaptic neurons  $j$  and postsynaptic neurons  $i$  are sorted and plotted in ascending frequency-order (i.e. from the smallest to the largest spike frequency). The final coupling matrices  $\varepsilon_{ij}$  for  $\tau_D = 0$  ms are shown in Fig. 7(d), (e) and (f) for  $\bar{\varepsilon}_{\text{initial}} = 0$ ,  $\bar{\varepsilon}_{\text{initial}} = 0.01$  and  $\bar{\varepsilon}_{\text{initial}} = 0.04$ , respectively. In all cases, the final coupling matrices have, predominantly, connections from faster to slower spiking neurons. This behaviour was also reported in Borges et al. [20] for a neural network with STDP. Fig. 7(g), (h) and (i) present our results for  $\bar{\varepsilon}_{\text{initial}} = 0$ ,  $\bar{\varepsilon}_{\text{initial}} = 0.01$ , and  $\bar{\varepsilon}_{\text{initial}} = 0.04$ , respectively, where  $\tau_D = 100$  ms. Due to the effect of STP on the dynamics of the neurons in the network, we observe the formation of different modules of directly connected neurons. Again, the effect of STP leads to a decrease on the influence of the fastest neurons to the slowest ones, allowing for the formation of connections among those neurons with similar spike frequencies. The size of these modules increases according to the intensity of the initial coupling and their different sizes can be explained by the analysis made based on the results in Figs. 5 and 6, where the coupling shortened the frequency differences, leading to the formation of connections especially among neurons with close frequencies. This also explains why the bigger-size modules are composed of the smallest-frequency spiking neurons. Increasing  $\tau_D$  to 1000 ms, leads to the disappearance of big modules and to the decrease of the number of coupled neurons, as shown in Fig. 7(j), (k) and (l). Therefore, by varying the STP recovery time  $\tau_D$ , one can control the formation of modules in neural networks with STP and STDP.

We also analyse the neural in-out degree distributions of our network. Fig. 8 shows the neural in-degree (blue) and out-degree (red) distributions for  $\bar{\varepsilon}_{\text{initial}} = 0.01$  with (a)  $\tau_D = 0$  ms, (b)  $\tau_D = 100$  ms, and (c)  $\tau_D = 1000$  ms with their respective network degree probability distributions shown in panels (d), (e) and (f). We calculate the in-out degree considering the connections among neurons and not their synaptic weight. In Fig. 8(a), we observe a decrease and an increase in the in-degree and out-degree distributions, respectively, which is in accordance with the one cluster formation shown in Fig. 7(e). This configuration also exhibits a uniform in-out degree probability distribution as displayed in Fig. 8(d). In Fig. 8(b), we see



**Fig. 7.** The combined effect of STP and STDP on connectivity in networks of  $N = 100$  excitatory coupled HH neurons and emergence of modular neural networks. The initial couplings used are: In (a)  $\bar{\varepsilon}_{\text{initial}} = 0$ , in (b)  $\bar{\varepsilon}_{\text{initial}} = 0.01$  and in (c)  $\bar{\varepsilon}_{\text{initial}} = 0.04$ . We consider  $\tau_D = 0$  ms for (d)  $\bar{\varepsilon}_{\text{initial}} = 0$ , (e)  $\bar{\varepsilon}_{\text{initial}} = 0.01$  and (f)  $\bar{\varepsilon}_{\text{initial}} = 0.04$ .  $\tau_D = 100$  ms for (g)  $\bar{\varepsilon}_{\text{initial}} = 0$ , (h)  $\bar{\varepsilon}_{\text{initial}} = 0.01$  and (i)  $\bar{\varepsilon}_{\text{initial}} = 0.04$  and  $\tau_D = 1000$  ms for (j)  $\bar{\varepsilon}_{\text{initial}} = 0$ , (k)  $\bar{\varepsilon}_{\text{initial}} = 0.01$  and (l)  $\bar{\varepsilon}_{\text{initial}} = 0.04$ . Note that the synaptic weights  $\varepsilon_{ij}$  (with  $i, j = 1, \dots, 100$ ) of the coupling matrices are encoded in grey scale in the colour bar.

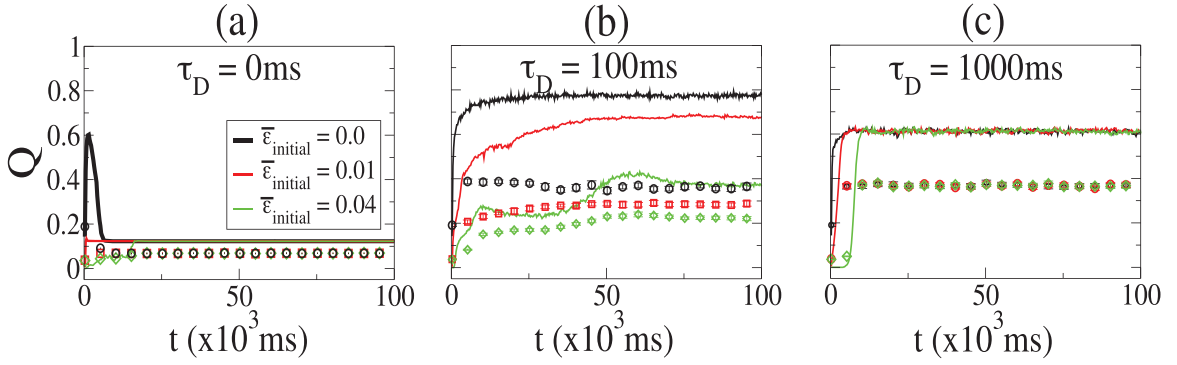


**Fig. 8.** Neuron in-out degree distributions for  $\bar{\varepsilon}_{\text{initial}} = 0.01$  for (a)  $\tau_D = 0$  ms, (b)  $\tau_D = 100$  ms, and (c)  $\tau_D = 1000$  ms with their respective network degree probability distributions shown in (d), (e), and (f). We calculate the in-out degree considering the connection between neurons and not the synaptic weight.

the same neural degrees patterns restricted to the neurons within the small clusters. Fig. 8(e) shows that the network in-out degree probability distribution is similar to a Gaussian distribution, however, it does not reflect the distribution inside the clusters, as shown in the inset of Fig. 8(e). In particular, the inset displays the uniform distribution among neurons within the first cluster. Fig. 8(c) and (f) show the degree and probability, respectively, the results for the in-degree is similar to the out-degree. The decrease in both degrees is expected as the network becomes sparse for higher  $\tau_D$ .

To study further the observed modules, we consider connections with  $\varepsilon_{ij} > 0.002$  (i.e. with connectivity strength bigger than 5% of the maximal coupling strength). Thus, weaker connections are not considered in the resulting network analysis.





**Fig. 9.** Time evolution of modularity  $Q$  for (a)  $\tau_D = 0$  ms, (b)  $\tau_D = 100$  ms and (c)  $\tau_D = 1000$  ms.  $Q$  was calculated over the time evolution of the coupling matrix and the coloured curves represent the network average initial coupling: the black curve is for  $\bar{\varepsilon}_{\text{initial}} = 0$ , the red for  $\bar{\varepsilon}_{\text{initial}} = 0.01$  and the green for  $\bar{\varepsilon}_{\text{initial}} = 0.04$ . The coloured symbols represent the average  $Q$  for random networks computed by rewiring the connections in the corresponding original networks. For example, the black symbols represent the average  $Q$  calculated over 20 random networks generated by rewiring randomly the connections in the original networks (black curve).

This procedure also avoids measurement errors that might be caused by connections whose weights fluctuate closely to zero over time.

To evaluate how modular structures evolve over time, we compute the modularity  $Q$  by using the Louvain method [55].  $Q$  is measured across network partitions in densely connected communities. In particular, the modularity assumes values in the range  $[-1, 1]$  comparing the density of connections within communities with the density among communities. The best network partition in modules is one that maximises modularity.  $Q$  is defined as [56]

$$Q = \frac{1}{W} \sum_i^N \sum_j^N \left( \varepsilon_{i,j} - \frac{\omega_i \omega_j}{W} \right) \delta(c_i, c_j), \quad (14)$$

where  $\omega_i = \sum_j^N \varepsilon_{ij}$  is the sum of the connection weights received by node  $i$  and  $W = \sum_i^N \sum_j^N \varepsilon_{ij}$  the sum of all weights in the coupling matrix. The term  $c_i$  represents the community that neuron  $i$  has been allocated to and  $\delta(c_i, c_j)$  is given by

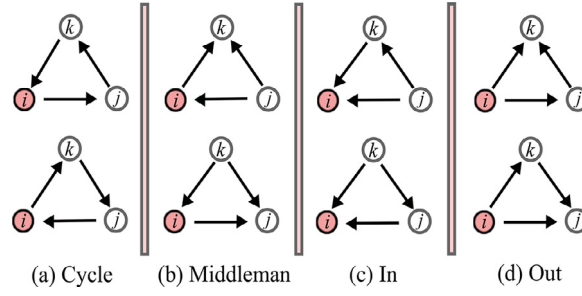
$$\delta(c_i, c_j) = \begin{cases} 1, & \text{if } c_i = c_j, \\ 0, & \text{otherwise.} \end{cases} \quad (15)$$

The Louvain [55] method is defined in two steps: At first, each node in the network is considered as a community in itself, and thus initially, there will be as many communities as nodes in the network. At this stage, each node  $i$  is reassigned to the community of each of its neighbours  $j$ , then  $i$  will be permanently fixed in the community that promotes the largest gain in modularity  $Q$  (positive gain). This process is applied repeatedly to all network nodes until there are no more gains in  $Q$ . The second step amounts to taking the defined communities at the end of the first step and consider them as the nodes of a new network. The weight of the connections between these new nodes is given by the sum of the weights of the connections between the nodes present in the communities (defined in the first step). Once the new network is computed, the first step can be applied again to its nodes. This process occurs repeatedly until no further changes to  $Q$  occur and a maximum value is thus obtained.

Fig. 9 shows the time evolution of modularity  $Q$  for the three recovery times,  $\tau_D = 0$  ms,  $\tau_D = 100$  ms and  $\tau_D = 1000$  ms. The colours represent the initial connectivity strengths in the coupling matrix: the black curve is for  $\bar{\varepsilon}_{\text{initial}} = 0$ , the red curve for  $\bar{\varepsilon}_{\text{initial}} = 0.01$  and the green curve for  $\bar{\varepsilon}_{\text{initial}} = 0.04$ . The coloured symbols represent the average  $Q$  values calculated over 20 random networks obtained by rewiring all connections of the original networks. We did this to compare our results with those obtained for random networks with the same connections and number of nodes (which we call random variants). Unless stated otherwise, the symbols in Fig. 9 represent the measurements taken on the random variants.

In Fig. 9(a), we observe that for  $\tau_D = 0$  ms (instantaneous recovery),  $Q$  is very low for all  $\bar{\varepsilon}_{\text{initial}}$  connectivity strengths, remaining constant after a small transient. These results confirm those in Fig. 7(d), (e) and (f). In Fig. 9(b) for  $\tau_D = 100$  ms,  $Q$  assumes its largest value for  $\bar{\varepsilon}_{\text{initial}} = 0$ , confirming what we observed in Fig. 7(g), (h) and (i) by comparing the number of modules. It is also apparent that  $Q$  stabilises for simulation times  $t$  greater than 80000 ms as there are no changes in the modular structures occurring in the networks anymore. Interestingly, Fig. 9(c) shows that for  $\tau_D = 1000$  ms,  $Q$  converges to the same value (i.e.  $Q \approx 0.6$ ) for all  $\bar{\varepsilon}_{\text{initial}}$  coupling strengths, again in accordance with the results in Fig. 7(j), (k), (l), which show networks with similar configurations regardless of the initial coupling. Finally, in all cases considered, we find that the modularity of the networks is bigger than their random variants.

Next, we compute other quantities that characterise the structure of the networks considered previously, such as the mean path-length, clustering coefficient and assortativity. In particular, a path is defined as the route that passes through network connections, connecting nodes  $i$  and  $j$ . The path with the shortest number of connections is called the shortest path



**Fig. 10.** The 8 distinct triangular motifs in directed networks. The motifs are considered with respect to node  $i$  depicted in red and are split into the cycle, middleman, and in out motifs. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

$l_{ij}$  [57]. The networks we consider here are directed, so  $l_{ij}$  is not necessarily equal to  $l_{ji}$ . In general, the shortest average path length is given by

$$\bar{L} = \frac{1}{N(N-1)} \sum_i \sum_{j \neq i} l_{ij}, \quad (16)$$

where  $N$  is the number of nodes in the network. In our approach, we calculated  $\bar{L}$  via a breadth-first search approach [57] and do not consider the weights of the connections. If there is no possible directed path  $l_{ij}$  between nodes  $i$  and  $j$ , then it is not considered in the calculations in Eq. (16).

The second quantity considered is the clustering coefficient  $CC_i$ , which measures the degree to which the neighbours of a node  $i$  are connected to each other and varies in  $[0,1]$ . It is computed by considering the number of triangular motifs made by node  $i$  and its neighbours compared to all possible triangular motifs of that node [57]. For directed networks, given 3 connected nodes  $i$ ,  $j$  and  $h$ , there are 8 distinct triangular motifs, shown in Fig. 10. These motifs are further organised into 4 groups when considering node  $i$  as the reference node: Fig. 10(a) shows a “cycle” motif, Fig. 10(b) a “middleman” motif, Fig. 10(c) an “in” motif and Fig. 10(d) an “out” motif. For each motif, we calculate  $CC_i^{(\text{cyc,mid,in,out})}$  relative to node  $i$ , as shown in [58].

For directed and weighted networks, there are 4 types of clustering coefficients

$$\begin{aligned} CC_i^{\text{cyc}} &= \frac{\frac{1}{2} \sum_j \sum_h [\varepsilon_{ij}^{1/3} \varepsilon_{jh}^{1/3} \varepsilon_{hi}^{1/3} + \varepsilon_{ih}^{1/3} \varepsilon_{hj}^{1/3} \varepsilon_{ji}^{1/3}]}{d_i^{\text{in}} d_i^{\text{out}} - d_i^{\leftrightarrow}}, \\ CC_i^{\text{mid}} &= \frac{\frac{1}{2} \sum_j \sum_h [\varepsilon_{ih}^{1/3} \varepsilon_{jh}^{1/3} \varepsilon_{ji}^{1/3} + \varepsilon_{ij}^{1/3} \varepsilon_{hi}^{1/3} \varepsilon_{hj}^{1/3}]}{d_i^{\text{in}} d_i^{\text{out}} - d_i^{\leftrightarrow}}, \\ CC_i^{\text{in}} &= \frac{\frac{1}{2} \sum_j \sum_h [\varepsilon_{ij}^{1/3} \varepsilon_{ih}^{1/3} \varepsilon_{jh}^{1/3} + \varepsilon_{ij}^{1/3} \varepsilon_{ih}^{1/3} \varepsilon_{hj}^{1/3}]}{d_i^{\text{in}} (d_i^{\text{in}} - 1)}, \\ CC_i^{\text{out}} &= \frac{\frac{1}{2} \sum_j \sum_h [\varepsilon_{hi}^{1/3} \varepsilon_{ji}^{1/3} \varepsilon_{jh}^{1/3} + \varepsilon_{hi}^{1/3} \varepsilon_{ji}^{1/3} \varepsilon_{hj}^{1/3}]}{d_i^{\text{out}} (d_i^{\text{out}} - 1)}, \end{aligned} \quad (17)$$

where  $d_i^{\text{in}} = \sum_j a_{ij}$  is the in-degree of node  $i$  and  $d_i^{\text{out}} = \sum_j a_{ji}$  its out-degree. The term  $d_i^{\leftrightarrow} = \sum_j a_{ij} a_{ji}$  represents the number of bilateral connections between node  $i$  and its neighbours. For the calculation of  $d_i^{\text{in}}$ ,  $d_i^{\text{out}}$  and  $d_i^{\leftrightarrow}$ , the coupling weights are not considered, only the number of connections, that is  $a_{ij} = 1$  if  $\varepsilon_{ij} > 0.002$ , otherwise  $a_{ij} = 0$ .

Consequently, the clustering coefficient of the network,  $CC^*$ , is calculated by averaging  $CC_i^*$  over all  $N$  nodes in the network

$$CC^* = \frac{1}{N} \sum_i CC_i^*, \quad (18)$$

where  $*$  stands for either of the cyc, mid, in or out motifs.

The last quantity considered in our study is assortativity, which is the correlation coefficient (i.e. the Pearson correlation [59]) between the degrees of nodes on two opposite ends of a connection in a network, for all connections in the network. This correlation varies in  $[-1, 1]$  and, is positive in assortative networks and negative in disassortative networks. Since our networks are directed and weighted, we use the four directed assortativity measures defined in [60].

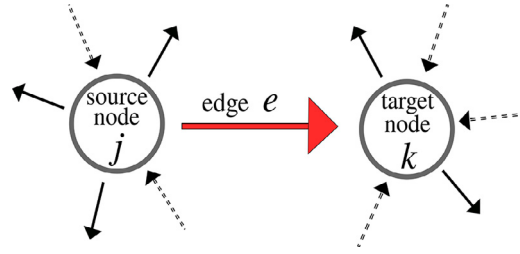


Fig. 11. Representative illustration of an edge  $e$  and, the source  $j$  and target  $k$  nodes considered in Eq. (19).

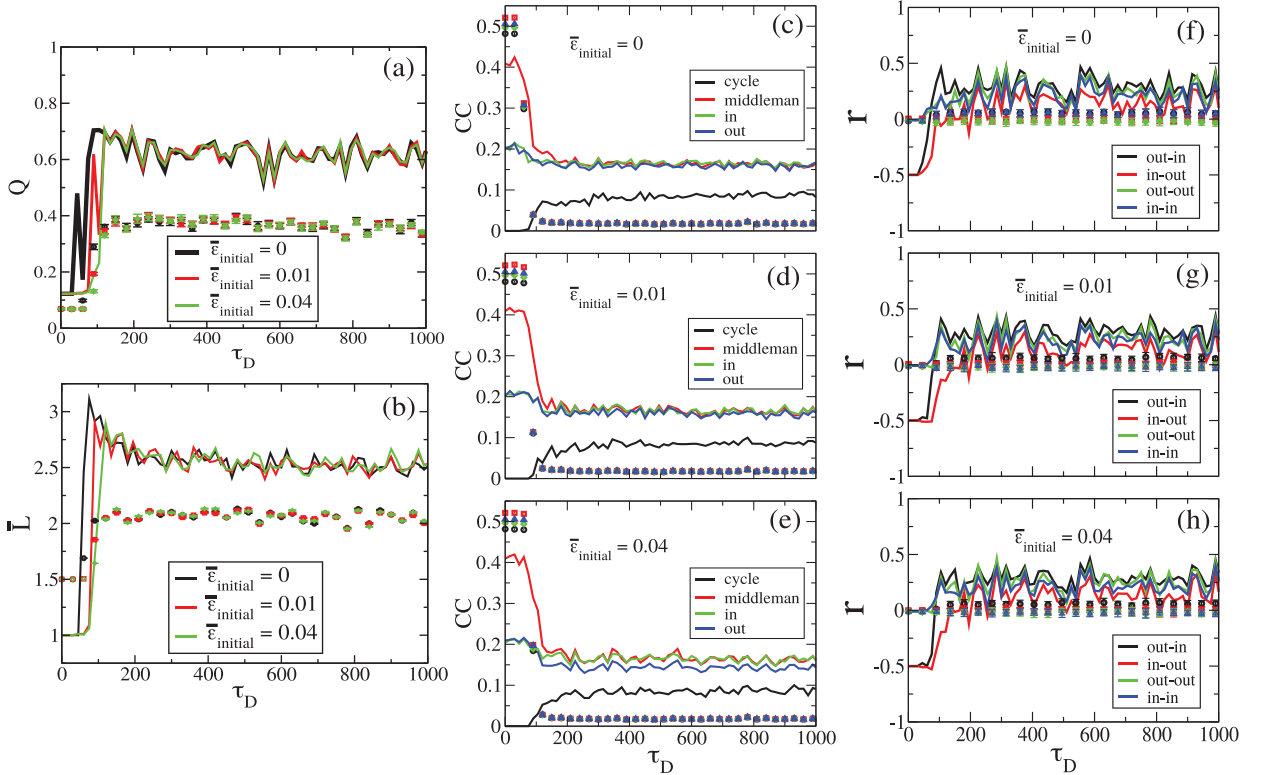


Fig. 12. Structural properties of the final coupling matrix configuration (coloured curves) and their respective random variants (averaged over 20 random networks, coloured symbols) as a function of the recovery time  $\tau_D$ . Panels (a) and (b) show the modularity  $Q$  and average shortest path length  $\bar{L}$  for three different average initial couplings:  $\bar{\epsilon}_{\text{initial}} = 0$  in black,  $\bar{\epsilon}_{\text{initial}} = 0.01$  in red and  $\bar{\epsilon}_{\text{initial}} = 0.04$  in green. In (c), (d) and (e), we show all clustering coefficients for  $\bar{\epsilon}_{\text{initial}} = 0$ ,  $\bar{\epsilon}_{\text{initial}} = 0.01$ , and  $\bar{\epsilon}_{\text{initial}} = 0.04$ , respectively. Panels (f), (g) and (h) show the assortativity  $r$  for  $\bar{\epsilon}_{\text{initial}} = 0$ ,  $\bar{\epsilon}_{\text{initial}} = 0.01$ , and  $\bar{\epsilon}_{\text{initial}} = 0.04$ , respectively.

In particular, let  $a$  and  $b$  be indices representing the type of degree (i.e. in- and out-degree) and  $j_e^a$  and  $k_e^b$  the  $a$ - and  $b$ -degree from source node  $j$  to target node  $k$  of an edge  $e$ . Fig. 11 shows an illustrative representation of edge  $e$  and their respective source nodes  $j$  and target nodes  $k$ . Foster et al. [60] defined the assortativity measure  $r(a, b)$  as

$$r(a, b) = \frac{N_e^{-1} \sum_e^{N_e} [(j_e^a - \bar{j}^a)(k_e^b - \bar{k}^b)]}{\sigma^a \sigma^b}, \quad (19)$$

where  $N_e$  is the number of edges in the network,  $\bar{j}^a = N_e^{-1} \sum_e^{N_e} j_e^a$  the average in- and out-degrees of the source node and  $\sigma^a = \sqrt{N_e^{-1} \sum_e^{N_e} (j_e^a - \bar{j}^a)^2}$  its deviation, calculated for all edges. The equations for  $\bar{k}^b$  and  $\sigma^b$  are similarly defined. In our approach, instead of in- and out-degrees, we use the in- and out-strength from the source and target nodes, given by the sum of their in- and out-coupling weights. A positive assortativity coefficient indicates that nodes tend to connect to other nodes with the same or similar strength.

Fig. 12 shows the results of the computations of the quantities discussed previously as a function of  $\tau_D$  in [0,1000] ms. In all computations, we consider the final coupling matrix obtained after  $t = 200000$  ms of simulated data. Fig. 12(a) shows the modularity  $Q$  for 3 initial couplings:  $\bar{\epsilon}_{\text{initial}} = 0$  in black,  $\bar{\epsilon}_{\text{initial}} = 0.01$  in red and,  $\bar{\epsilon}_{\text{initial}} = 0.04$  in green and their random

variants given by different symbols with their respective colours. We note that for  $\tau_D \lesssim 30$  ms and for all 3  $\bar{\epsilon}_{\text{initial}}$  values,  $Q$  is low. For  $30 \text{ ms} < \tau_D \lesssim 120$  ms,  $Q$  depends on  $\bar{\epsilon}_{\text{initial}}$ . In particular, the lower the initial average coupling  $\bar{\epsilon}_{\text{initial}}$ , the more modular the final configuration of the network. For  $\tau_D \gtrsim 120$  ms, there is no dependence on  $\bar{\epsilon}_{\text{initial}}$  and the resulting networks exhibit similar modularity values. In all cases of  $\bar{\epsilon}_{\text{initial}}$ ,  $Q$  is higher than that calculated for their random variants, which shows that the resulting networks do not have characteristics of random networks.

Fig. 12 (b) shows the results for the average path length  $\bar{L}$  for the same three  $\bar{\epsilon}_{\text{initial}}$  values. For  $\tau_D < 30$  ms, the average path  $\bar{L} = 1$ , in accordance with Fig. 7(d), (e) and (f), where we observe a network where all nodes are connected to all other nodes with unidirectional connections. For  $30 \text{ ms} < \tau_D < 120$  ms, the average path length  $\bar{L}$  is bigger for weaker initial couplings. This is because weak initial couplings result in the formation of a larger number of modular structures (as seen in Fig. 12(a)), which results in a greater path for one node to access other nodes in different modules, possibly crossing through other modules. For  $\tau_D > 120$  ms, we see that  $\bar{L} \approx 2.5$  for all 3  $\bar{\epsilon}_{\text{initial}}$  values. In contrast, the average path lengths  $\bar{L}$  for the random variants are smaller than those for the original networks.

Fig. 12 (c), (d) and (e) present the 4 types of clustering coefficients for  $\bar{\epsilon}_{\text{initial}} = 0$ ,  $\bar{\epsilon}_{\text{initial}} = 0.01$  and  $\bar{\epsilon}_{\text{initial}} = 0.04$ . The black curve represents  $\text{CC}^{\text{cyc}}$ , the red  $\text{CC}^{\text{mid}}$ , the green  $\text{CC}^{\text{in}}$  and the blue  $\text{CC}^{\text{out}}$ . In all cases, the results are similar. It is worth it to note that for  $\tau_D < 120$  ms,  $\text{CC}^{\text{mid}}$  has the highest value and  $\text{CC}^{\text{cyc}}$  the lowest. These results build on what we have already observed: the action of STDP promotes connections from faster to slower spiking neurons and do not permit cyclic connections (see Fig. 10(a)). As  $\tau_D$  increases,  $\text{CC}^{\text{mid}}$  converges to  $\text{CC}^{\text{in}}$ ,  $\text{CC}^{\text{out}} \approx 0.15$  and  $\text{CC}^{\text{cyc}}$  increases to about 0.09. For the random variants of the networks (coloured symbols), we see that for  $\tau_D < 120$  ms, all clustering coefficients are similar and larger than those for the original networks. This is because for such  $\tau_D$  values, the networks from which they were generated (Fig. 7(d), (e), (f)) are densely connected and the high number of connections allows for the formation of triangular motifs without any of the 4 types occurring preferentially. For  $\tau_D > 120$ , all random networks have their clustering coefficients fluctuate near zero. In this case, the generating networks are more sparse (as in Fig. 7(g), (h), (i)), which makes their random variants have low probability in forming triangular motifs. We note that these clustering coefficients are not high enough to claim with certainty that the networks have a small-world topology.

The right column of plots in Fig. 12 shows the 4 assortativity measures,  $r$ , computed for the considered networks with  $\bar{\epsilon}_{\text{initial}} = 0$ ,  $\bar{\epsilon}_{\text{initial}} = 0.01$ , and  $\bar{\epsilon}_{\text{initial}} = 0.04$ . The black curves represent the out-strength/in-strength correlation  $r(\text{out}, \text{in})$  between the source and target nodes, the red curve the in-strength/out-strength correlation  $r(\text{in}, \text{out})$ , the green curve the out-strength/out-strength correlation  $r(\text{out}, \text{out})$  and the blue curve, the in-strength/in-strength correlation  $r(\text{in}, \text{in})$ .

In Fig. 12(f), (g), (h) and for the three  $\bar{\epsilon}_{\text{initial}}$  values, we find that for  $\tau_D < 100$  ms, the networks are disassortative as  $r(\text{out}, \text{in})$  and  $r(\text{in}, \text{out})$  are negative and, at the same time, neither assortative nor disassortative as  $r(\text{out}, \text{out})$  and  $r(\text{in}, \text{in})$  are approximately equal to 0. The predominance of the connections from the fastest to the slowest neurons is responsible for the disassortative properties of the network. Neurons that receive many strong connections (the slowest ones) generally do not send connections to neurons that also send many strong connections (the fastest ones). The values of  $r(\text{out}, \text{out})$  and  $r(\text{in}, \text{in})$  are low due to the fact there are no neurons that send (receive) many strong connections connected with other neurons that send (receive) many strong connections. As  $\tau_D$  increases, all correlations grow, being mostly positive with  $r(\text{out}, \text{in})$  being the largest. This shows a greater correlation in the network of high out-strength nodes to connect with nodes with high in-strength and corroborates the results in Fig. 7, where connections occur preferably from the faster neurons (which have a high out-strength) to the slower ones (with high in-strength). For the random variants of these networks, all correlations are close to zero for the entire  $\tau_D$  range, a result completely different to those for the original networks. We thus conclude that the original networks are far from being purely random networks, exhibiting a type of preferential attachment in their connectivities.

In the analysis of the structural properties of the networks, the topology may vary greatly depending on  $\tau_D$ . In all cases studied, the networks were different from random or small-world networks, since the average path lengths are bigger than in random networks and their clustering coefficients are small than in random networks. This corroborates further the results obtained previously here that there is a complex process taking place in which modules are formed via preferential attachment.

## 5. Conclusions

In this paper, we studied the effects of plasticity (STP and STDP) on networks of excitatory coupled Hodgkin-Huxley neurons. Neural plasticity is responsible for alterations in the organisation and structure of the brain, and both play an important role in synaptic weights. Besides, STDP has a longer time scale than STP, so it affects differently the structure and function in brain networks.

We started analysing the effect of STP in a pair of neurons that are initially, either uncoupled or bidirectionally coupled. For initially uncoupled neurons, the action of STP and STDP promotes directed connections among neurons with small spike frequency differences, from the faster to the slower spiking neurons. The increase of the recovery-time shortened the interval of frequency differences where connections are formed. When neurons are initially coupled, their frequency difference is smaller and increases the size of the area of directed connections. We found that STP induces uncoupling, depending on the recovery time: the bigger the recovery time, the smaller the interval of frequency difference that allows for the formation of connections.

Next, we build a neural network with an all-to-all topology. Considering only STDP, the coupling matrix exhibits directed connections from neurons with high to neurons with low spike frequencies. We have shown that due to STP, neural networks equipped with STDP facilitate the formation of synapses among neurons with similar spike frequencies only and that modular neural networks can emerge as a direct result of the combined effect of STP and STDP, a phenomenal structure also depicted by neurophysiological and experimental studies. However, by increasing the STP recovery time, the number of connections decreased and as a consequence, the modules disappeared. That is actually a way to control the modular organisation in neural networks. The structure of these modular networks is complex, unlike those in random or small-world networks, resembling more to networks with preferential attachment properties.

In future, we plan to study neural networks with greater diversity in chemical synapses, addressing other STP and STDP rules related to synapses between excitatory-inhibitory and inhibitory-inhibitory neurons [54]. Finally, another interesting aspect of our work would be the introduction of time delay in the synaptic transmission to study how it affects the evolution of the couplings and the modular properties of neural networks.

## Declaration of Competing Interest

We wish to confirm that there are no known conflicts of interest associated with this work and there has been no significant financial support for this work that could have influenced its outcome.

No conflict of interest exists.

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