

# A sensitive temperature probe based on $\text{Er}^{3+}$ -doped $\text{GeO}_x$ films

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## ABSTRACT

The light emission due to  $\text{Er}^{3+}$  ions when inserted in a  $\text{GeO}_x$  film was analyzed at different temperatures. The film was prepared by sputtering a combined Ge + Er solid target and the incorporation of oxygen took place by thermal annealing the film under a continuous flow of  $\text{O}_2$ . According to the experimental results, the  $\text{Er}^{3+}$ -related  $^2\text{H}_{11/2} \rightarrow ^4\text{I}_{15/2}$  and  $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$  optical transitions – corresponding to the green region of the visible spectrum of electromagnetic radiation – were highly influenced by the local temperature. Compared with other  $\text{Er}^{3+}$ -doped materials, the present  $\text{GeO}_x$ :Er film exhibited a rather large  $^2\text{H}_{11/2}$ -to- $^4\text{S}_{3/2}$  energy levels separation, suggesting its potential to act as a very sensitive temperature probe in the  $\sim 300$ – $750$  K range. Along with the  $\text{Er}^{3+}$ -related optical emission characteristics, the atomic composition and structure of the  $\text{GeO}_x$ :Er film were also investigated in detail.

## Introduction

At ambient conditions, crystalline germanium dioxide  $\text{GeO}_2$  adopts either the  $\alpha$ -quartz-like (hexagonal) or the rutile-like (tetragonal) atomic structures [1]. Also,  $\text{GeO}_2$  is a thermal stable ( $\sim 1400$  K melting point), wide optical bandgap ( $\sim 6$  eV), high dielectric constant (5.6–8.7) material, that presents a refractive index (1.65) slightly higher than that of  $\text{SiO}_2$  (1.55) [2]. Accordingly, films of  $\text{GeO}_2$  were considered in optical applications ranging from, for example, infrared filters and detectors to waveguides. Special attention has also been given to sub-stoichiometric  $\text{GeO}_2$  (or  $\text{GeO}_x$ ) to function as insulating or passivating layers in electronic components, as well as for storage information purposes [3]. Considering the power exerted by the atomic structure–composition onto the properties of the  $\text{GeO}_x$  films, different methods of preparation has been reviewed, which included: evaporation (e-beam or thermal), sputtering (dc or rf), laser ablation, plasma-assisted, sol–gel etc. [4]. In all of them, the insertion of oxygen came about during the film deposition itself and/or after preparation by means of special sample processing. Some of these features inspired the study of  $\text{GeO}_x$  films combined with Er, in which case the main focus was the  $\text{Er}^{3+}$ -related optical emission at  $\sim 1540$  nm [5,6]. However, in wide-bandgap materials,  $\text{Er}^{3+}$  ions are known to present well-defined and rather intense optical transitions in the visible region – just as required in applications such as light display devices or temperature sensors [7,8].

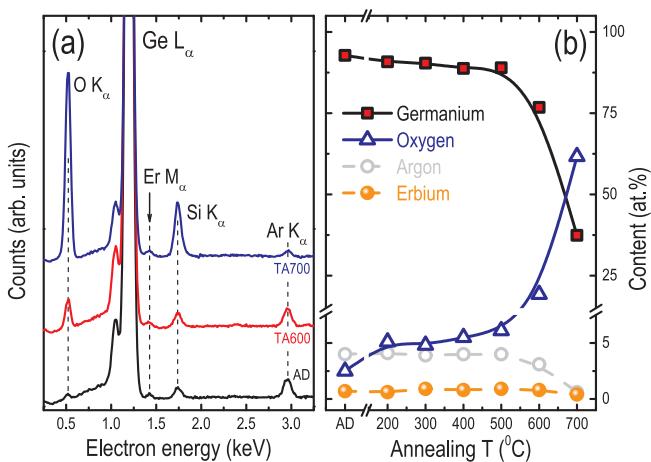
The chemical–physical properties of  $\text{GeO}_x$  allied to the unique optical characteristics of  $\text{Er}^{3+}$  ions gave rise to this work that is mainly concerned with the production of Er-doped  $\text{GeO}_x$  films and their visible  $\text{Er}^{3+}$ -related light emission. Besides, contrary to most of the studies devoted to optical thermometric purposes (i.e., involving the

association of  $\text{Er}^{3+}$ - $\text{Yb}^{3+}$  and non-linear up-conversion processes [9,10]) the present  $\text{Er}^{3+}$ -doped  $\text{GeO}_x$  films are easy to produce, stable, and very sensitive. The near-infrared light emission due to the  $\text{Er}^{3+}$  ions, as well as the main compositional-structural aspects of the  $\text{GeO}_x$ :Er system are also presented and comprehensively discussed in view of the existing literature.

## Experimental details

The starting material consisted in one Er-doped Ge film deposited onto a mirror polished crystalline Si substrate. The film was deposited by  $\text{Ar}^+$  ion sputtering a Ge target (99.999% pure,  $126\text{ cm}^2$  of area) partially covered with Er pieces (99.9%,  $8\text{ cm}^2$ ), so that its final composition was decided by the relative Er-to-Ge target areas and corresponding sputtering yields [11,12]. During deposition, the Si substrate was kept at  $125 \pm 10$  °C and, after 75 min, a  $1\text{ }\mu\text{m}$  thick (uniform) film was achieved. Also, a deposition pressure of  $1.5 \times 10^{-4}$  Torr (preceded by a base pressure of  $2 \times 10^{-6}$  Torr), and a  $0.8\text{ W cm}^{-2}$  radio frequency (13.56 MHz) power density were adopted. After deposition, the “Er-doped Ge film + Si substrate” set was cut into various pieces of  $\sim 1 \times 1\text{ cm}^2$  and annealed under a continuous flow of oxygen gas (99.99% pure), giving rise to the Er-doped  $\text{GeO}_x$  (or  $\text{GeO}_x$ :Er) films. The thermal annealing treatments were cumulative, 30 min long, and were carried out at 200, 300, 400, 500, 600, and  $700 \pm 10$  °C.

All  $\text{GeO}_x$ :Er films – as-deposited (AD) and after thermal annealing (TA) treatments – were investigated by means of energy-dispersive X-ray (EDX), Raman scattering, and photoluminescence (PL) techniques.



**Fig. 1.** (a) EDX spectra (15 keV electron energy and  $30 \times 30 \mu\text{m}^2$  analyzed area) of some Er-doped  $\text{GeO}_x$  films: as-deposited (AD), and after thermal annealing at 600 °C (TA600) and 700 °C (TA700). The spectra were vertically shifted for clarity reasons. (b) Atomic composition (as determined from the EDX analysis) of the Er-doped  $\text{GeO}_x$  films, as a function of the annealing temperature. The experimental data error is typically the size of the symbols.

## Results and discussion

### Composition & structure

As a result of the deposition method and conditions, the  $\text{GeO}_x:\text{Er}$  films presented ~0.5 at.% of Er and oxygen concentrations that scaled with the annealing temperature. The EDX spectra also indicated the presence of some embedded argon atoms (ranging from ~4 to 1 at.%) and the contribution from the Si substrate – this one being more evident owing to the different electron penetration depths in Ge and  $\text{GeO}_x$  [13]. The EDX spectra and corresponding data analysis are shown in Fig. 1. Based on these results it is clear that all  $\text{GeO}_x:\text{Er}$  films are sub-stoichiometric (i.e.,  $[\text{O}] < 66.6$  at.%) and, hence, their generic  $\text{GeO}_x$  designation.

In addition to the compositional changes, thermal annealing at increasing temperatures modified the atomic structure of the films as well. The Raman spectra of Fig. 2(a) show that the as-deposited film is

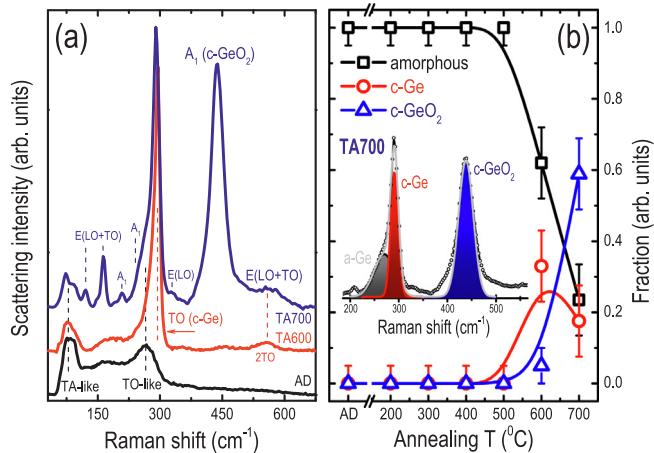
essentially amorphous, as denoted by the broad Ge–Ge scattering signals at:  $\sim 85 \text{ cm}^{-1}$  (corresponding to the transverse-acoustic TA-like phonon mode);  $\sim 265 \text{ cm}^{-1}$  (transverse-optic TO-like); and  $\sim 167$  and  $\sim 230 \text{ cm}^{-1}$  (longitudinal-acoustic and optic-like modes, respectively) [14]. The amorphous Ge structure remained up to 700 °C, but after annealing at 600 °C the Raman spectra indicated the presence of crystalline (c.)Ge – as denoted by the well-defined signal at  $\sim 300 \text{ cm}^{-1}$  (TO) and its overtone at  $560 \text{ cm}^{-1}$  (2TO) [15].

After treatment at 700 °C, new Raman scattering signals developed – all of them related to crystalline ( $\alpha$ -quartz)  $\text{GeO}_2$ : E(LO + TO) modes at  $\sim 120$ , 160, and  $570 \text{ cm}^{-1}$ , E(LO) at  $\sim 320 \text{ cm}^{-1}$ , and  $A_1$  at  $\sim 205$ , 245, and  $435 \text{ cm}^{-1}$  [16]. A proper way to evaluate these structural changes relies on the analysis of the most intense Raman signals by means of the relative amount (or fraction) of Ge atoms: in the amorphous (TO-like mode at  $265 \text{ cm}^{-1}$ ) and crystalline (TO at  $300 \text{ cm}^{-1}$ ) forms, and bonded to oxygen ( $A_1$  at  $435 \text{ cm}^{-1}$ ), for example. These Ge fractions were estimated according to their corresponding Raman signals [17], and are displayed in Fig. 2(b). The results resemble very much those obtained from EDX (Fig. 1) and despite the evident data error (mainly because of the Raman fitting process and related scattering cross-sections) they indicate that, after treatment at 700 °C, the  $\text{GeO}_x:\text{Er}$  film is basically made of ~40% of Ge and ~60% of oxygen. Moreover, temperatures above 500 °C were needed to induce Ge crystallization and the formation of Ge–O bonds.

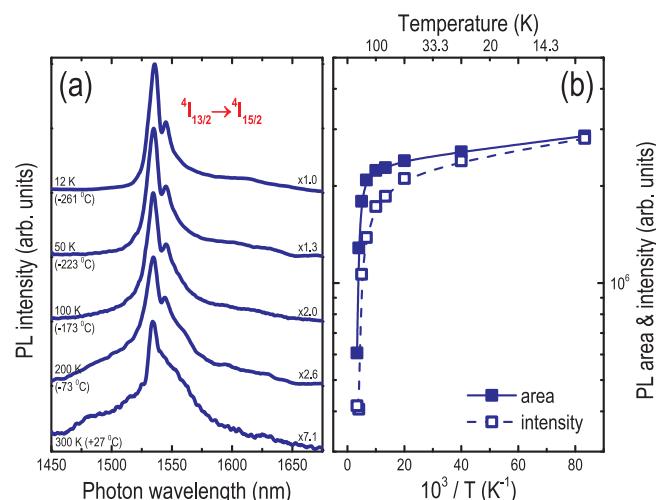
### $\text{Er}^{3+}$ -related light emission

Traditionally, most of the studies involving the doping of semiconductors with  $\text{Er}^{3+}$  ions focused on their characteristic near-infrared light emission at  $\sim 1540 \text{ nm}$ . At this photon wavelength, which coincides with the lowest transmission loss region (C-band) of the silica-based optical fibers, the combination of  $\text{Er}^{3+}$  ions with semiconductor materials was expected to produce great advances in the field of optoelectronics [18]. An interesting outcome of these studies was to verify that the light emission quenching at  $\sim 1540 \text{ nm}$  tends to occur at higher temperatures in wide-bandgap materials [19], in which case the  $\text{Er}^{3+}$  ions can also exhibit emission in the visible range [20].

The PL spectra, in the 1450–1675 nm wavelength range, of the  $\text{GeO}_x:\text{Er}$  film annealed at 700 °C are displayed in Fig. 3(a). The measurements were performed at different temperatures and clearly show the  $4\text{I}_{13/2} \rightarrow 4\text{I}_{15/2}$   $\text{Er}^{3+}$ -related transition and respective Stark splitting



**Fig. 2.** (a) Room-temperature, non-polarized, Raman scattering spectra (488.0 nm photon excitation) of some Er-doped  $\text{GeO}_x$  films: as-deposited (AD), and after thermal annealing at 600 and 700 °C. The spectra were vertically shifted for clarity. (b) Fraction of amorphous Ge, c-Ge, and c-GeO<sub>2</sub> Raman signals as the Er-doped  $\text{GeO}_x$  films were annealed at increasing temperatures. The inset shows the deconvolution of a typical Raman spectrum (sample TA700) illustrating the contributions due to amorphous Ge, c-Ge, and c-GeO<sub>2</sub>.



**Fig. 3.** (a) Photoluminescence spectra (488.0 nm photon excitation) of one Er-doped  $\text{GeO}_x$  film (TA700) at different temperatures of measurement. The spectra were normalized (see the multiplying factors) and vertically shifted for clarity reasons. (b) PL intensity at  $1450 \pm 5 \text{ nm}$  and PL integrated area (1475–1650 nm range) as a function of the inverse of temperature. PL data error is typically smaller than the size of the symbols.

features around 1540 nm [21]. The influence exerted by the temperature of measurement onto the  $\text{Er}^{3+}$ -related PL intensity is presented in Fig. 3(b), that also shows the corresponding PL area i.e., PL signal integrated over a predefined spectral band. According to the data, it is clear the almost one order of magnitude variation experienced by the PL signal in the 12–300 K (or  $-260$  to  $27^\circ\text{C}$ ) range, and that a very low quenching temperature  $T_q$  applies. In fact,  $T_q = 25 \pm 5$  K – as defined by the  $I_{\text{PL}}(T_q) = 0.95 I_{\text{PL}}(\text{lowest})$  variation in the near-infrared PL signal [19] – is perfectly consistent with the presence of Ge–Ge bonds (Fig. 2) in the  $\text{GeO}_x$ :Er film.

As pointed out by literature, the optical bandgap of  $\text{GeO}_x$  stays between  $\sim 1$  eV [22] and 6 eV [23], and it is known to be highly dependent on the deposition conditions and sample processing details that, ultimately, determine the oxygen content and structure of the material [4]. No optical absorption measurements were carried out on the present  $\text{GeO}_x$ :Er (TA700) film, but based on its atomic composition [Fig. 1(b)] it is reasonable to ascribe an optical bandgap  $E_{\text{gap}}$  of  $\sim 4.5 \pm 0.5$  eV [24].

Analogous to other  $\text{Er}^{3+}$ -doped wide-bandgap materials like  $\text{GeN}$  ( $E_{\text{gap}} \sim 2.5$  eV [7]),  $\text{TiO}_2$  ( $E_{\text{gap}} \sim 3$  eV [8]),  $\text{AlN}$  ( $E_{\text{gap}} \sim 4$  eV [25]), and  $\text{SiN}$  ( $E_{\text{gap}} \sim 4.5$  eV [26]), in addition to the PL signal at  $\sim 1540$  nm, the  $\text{GeO}_x$ :Er film also exhibited light emission in the visible range.

As displayed in Fig. 4(a), one of these  $\text{Er}^{3+}$ -related emissions took place in the green region of the electromagnetic spectrum and correspond to the  $^2\text{H}_{11/2} \rightarrow ^4\text{I}_{15/2}$  (at  $\sim 525$  nm) and  $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$  (at  $\sim 550$  nm) optical transitions [21]. As can be seen, the transitions exhibit typical Stark splitting and, most importantly, their overall shape and intensity were affected by temperature [Fig. 4(b)]. It happens because of the proximity of the  $^2\text{H}_{11/2}$  and  $^4\text{S}_{3/2}$  levels such that, depending on the local temperature, electrons can migrate from one level to another [27,28]. More specifically: at higher temperatures electrons recombine preferentially from the upper level at  $\sim 525$  nm, and *vice-versa*. Part of this phenomenon can be described by an Arrhenius law [Eq. (1)] that relates the PL intensity  $I_{\text{PL}}$  (or integrated PL area) ratio  $R(T)$  of the  $\text{Er}^{3+}$ -related green transitions and the temperature of measurement. It involves the energy gap between the  $^2\text{H}_{11/2}$  and  $^4\text{S}_{3/2}$  levels ( $\Delta\epsilon$ ), details regarding their degeneracy and emission probability (exponential pre-factor  $K$ ), and the Boltzmann constant ( $k_B$ ):

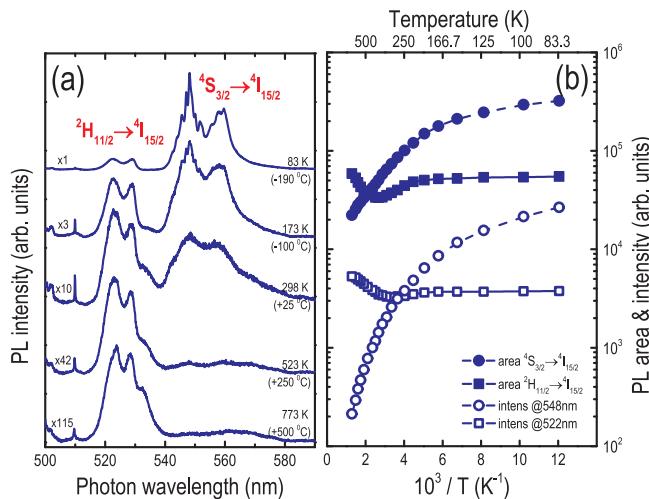


Fig. 4. (a) Photoluminescence spectra (488.0 nm excitation), in the 500–590 nm range, of a  $\text{GeO}_x$ :Er film (TA700) at different temperatures of measurement. The spectra were normalized (see the multiplying factors) and vertically shifted for clarity. The contribution at  $\sim 510$  nm (and below) corresponds to the Raman signal due to c- $\text{GeO}_2$ . (b) PL intensity (at  $525 \pm 3$  and  $550 \pm 3$  nm) and PL integrated area (515–540 and 540–580 nm range) as a function of the inverse of temperature.

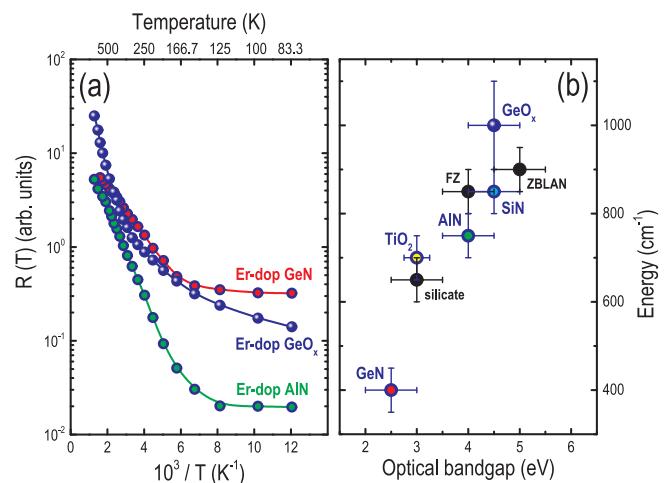


Fig. 5. (a) Photoluminescence intensity ratio  $R(T)$  as a function of the inverse of temperature, of  $\text{Er}^{3+}$ -doped  $\text{GeN}$ ,  $\text{GeO}_x$ , and  $\text{AlN}$  films. (b) Energy separation  $\Delta\epsilon$  (between the  $^2\text{H}_{11/2}$  and  $^4\text{S}_{3/2}$  levels) as a function of the optical bandgap of various  $\text{Er}^{3+}$ -doped materials:  $\text{GeN}$  film [7], lead-silicate glass [36],  $\text{TiO}_2$  film [8],  $\text{AlN}$  film [25], fluorozirconate FZ glass [27],  $\text{SiN}$  film [20,26],  $\text{ZrF}_4\text{-BaF}_2\text{-LaF}_3\text{-AlF}_3\text{-NaF}$  ZBLAN glass [28], and  $\text{GeO}_x$  film (this work). Error bars comprise both experimental data dispersion and uncertainties from literature.

$$R(T) = \frac{I_{\text{PL}}(^2\text{H}_{11/2} \rightarrow ^4\text{I}_{15/2})}{I_{\text{PL}}(^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2})} = K \exp\left(\frac{-\Delta\epsilon}{k_B T}\right) \quad (1)$$

In spite of its convenience, Eq. (1) considers that all electrons reaching the  $^2\text{H}_{11/2}$  level originated from the  $^4\text{S}_{3/2}$  state and ignores other (non-radiative) transition paths [8] – that eventually could be influential at extreme temperatures.

Nonetheless, provided that a calibration curve exists – similar to those presented in Fig. 5(a) – any  $\text{Er}^{3+}$ -doped wide-bandgap material can be used to efficiently probe the local temperature [29–33]. In this case, the host material will be chosen based on its physical properties and in view of the temperature range of interest. This form the basis of an optically-based temperature sensor, according to which the temperature can be determined, almost instantly, with minimum (or no) physical interference by optical means [34,35]. In common, most of the  $\text{Er}^{3+}$ -doped materials considered for optical thermometry purposes seem to strictly follow Eq. (1) only above  $\sim 150$  K, and to present  $\Delta\epsilon$  energies that depend on the main characteristics of the host [8,20,26–29,34]. Part of these features are present in Fig. 5: (a) in which a single exponential behavior applies at  $T \geq 150$  K (or even  $T > 300$  K for the  $\text{GeO}_x$ :Er film); and (b) that shows the relationship between  $\Delta\epsilon$  and the optical bandgap of some  $\text{Er}^{3+}$ -doped materials. In the former case, other optical excitation–recombination mechanisms can be competing with the  $\text{Er}^{3+}$ -related green emission and, hence, the different temperature ranges in which Eq. (1) applies. In the latter, the atomic structure–composition leading to wider optical bandgap materials can also be at the origin of larger  $\Delta\epsilon$  values [19,29]. At this point, one can only conclude about the influence exerted by the local environment (or crystal field effects) onto the  $\text{Er}^{3+}$ -related green light emission, and that more research is needed. However, allied to an optical bandgap that can reach  $\sim 6$  eV, the results of Fig. 5 are very promising, suggesting the suitability of the  $\text{GeO}_x$ :Er system to act as an efficient optically-based temperature sensor (in the  $\sim 300$ – $750$  K range).

## Concluding remarks

The atomic structure and composition of  $\text{Er}^{3+}$ -doped  $\text{GeO}_x$  films were studied by different spectroscopic techniques. The films were prepared by  $\text{Ar}^+$  ion sputtering a solid  $\text{Ge} + \text{Er}$  target and, subsequently, thermal annealed under a flow of  $\text{O}_2$ . Following this approach,

different films were achieved, in which their main properties scaled with the temperature of anneal – ranging from amorphous Ge to sub-stoichiometric crystalline  $\text{GeO}_2$  (or  $\text{c-GeO}_x$ ). In all cases, the Er concentration remained in the low 0.5 at.%. The  $\text{Er}^{3+}$ -related light emission of the film annealed at 700 °C was investigated in detail – both in the near-infrared ( ${}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{15/2}$  transition at  $\sim 1540$  nm) and visible ( ${}^2\text{H}_{11/2} \rightarrow {}^4\text{I}_{15/2}$  and  ${}^4\text{S}_{3/2} \rightarrow {}^4\text{I}_{15/2}$  transitions at  $\sim 525$  and 550 nm) ranges.

According to the experimental results: (1) the originally amorphous Ge films became partially crystallized after annealing above 500 °C; (2) likewise, crystalline Ge–O bonds developed only after treatment at 600 and 700 °C, in which cases the oxygen concentrations reached  $\sim 20$  and 60 at.%; (3) accordingly, the optical bandgap of the  $\text{GeO}_x$ :Er TA700 film was anticipated to be around  $4.5 \pm 0.5$  eV; (4) despite such a wide optical bandgap, the T-quenching experienced by the  ${}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{15/2}$  near-infrared light emission (at  $\sim 1540$  nm) is consistent with the presence of Ge–Ge bonds; (5) regarding the  $\text{Er}^{3+}$ -related emissions taking place in the green visible region ( ${}^2\text{H}_{11/2} \rightarrow {}^4\text{I}_{15/2}$  and  ${}^4\text{S}_{3/2} \rightarrow {}^4\text{I}_{15/2}$  transitions), they proved to be highly susceptible to temperature effects, presenting a single exponential behavior in the  $\sim 300$ –750 K range.

Altogether, this work reports the first results involving the  $\text{Er}^{3+}$ -related green light emissions from the  $\text{GeO}_x$ :Er system, which is being proposed as a very effective medium for optical thermometric applications.

## Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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