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Measurement of the cosmic-ray energy spectrum above 2.5 EeV using 19 years of operation of the Pierre Auger Observatory

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Abstract

We present the spectrum of cosmic rays with energies above 2.5 EeV measured at the Pierre Auger Observatory after 19 years of operation, covering the period before the AugerPrime upgrade. Two independent event sets from the surface array of 1500-m spaced detectors are combined, yielding a total exposure of 104 900 km² sr yr. The first set includes events with zenith angles less than 60°, while the second consists of events between 60° and 80°, for which azimuthal asymmetries must be accounted for in the energy estimator. The threshold energy is chosen to ensure a trigger efficiency of the surface detector greater than 97%, thus minimizing composition biases. The energy scale is determined using high-quality fluorescence measurements, providing calorimetric estimates without reliance on simulations.

A statistically successful combination is achieved within the uncorrelated systematic uncertainties of the individual spectra. All spectra are consistent when analyzing potential declination dependences, except for a mild modulation expected from the previously reported dipolar anisotropy. In particular, this statement applies to the northernmost declination band [+24.8°, +44.8°], where only events with zenith angles between 60° and 80° contribute. Beyond the firmly established ankle and suppression spectral features, the combined spectrum across declinations -90° to +45° provides a measurement of the instep feature with more than 5σ confidence.

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Measurement of the cosmic-ray flux from 2.5 EeV by the Pierre Auger Observatory after 19 years of operation

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We present the spectrum of cosmic rays with energies above 2.5 EeV measured at the Pierre Auger Observatory after 19 years of operation, covering the period before the AugerPrime upgrade. Two independent event sets from the surface array of 1500 m-spaced detectors are combined, yielding a total exposure of $104\,900\text{ km}^2\text{ sr yr}$. The first set includes events with zenith angles less than 60° , while the second consists of events between 60° and 80° , for which azimuthal asymmetries must be accounted for in the energy estimator. The threshold energy is chosen to ensure a trigger efficiency of the surface detector greater than 97%, thus minimizing composition biases. The energy scale is determined using high-quality fluorescence measurements, providing calorimetric estimates without reliance on simulations.

A statistically successful combination is achieved within the uncorrelated systematic uncertainties of the individual spectra. All spectra are consistent when analyzing potential declination dependences, except for a mild modulation expected from the previously reported dipolar anisotropy. In particular, this statement applies to the northernmost declination band $[+24.8^\circ, +44.8^\circ]$, where only events with zenith angles between 60° and 80° contribute. Beyond the firmly established ankle and suppression spectral features, the combined spectrum across declinations -90° to $+45^\circ$ provides a measurement of the instep feature with more than 5σ confidence.

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1. Introduction

The Pierre Auger Observatory [1] has accumulated, since its beginning in 2004, the largest dataset of ultra-high-energy cosmic rays currently available. The volume and quality of the acquired data allowed us to measure a dipolar modulation of $\sim 6\%$ of an otherwise isotropic distribution of the arrival directions of cosmic rays with a significance greater than 5σ [2]. Motivated by the observed dipole, we previously searched for a modulation of the spectrum with declination [3, 4]. We recently conducted a more comprehensive search exploiting a $\sim 66\%$ larger exposure with a dataset that also includes events arriving with a zenith angle between 60° and 80° thus extending the declination reach from $+24.8^\circ$ to $+44.8^\circ$ [5] that is also reported here.

We present an energy spectrum using events observed by the water-Cherenkov detectors of the Surface Detector Array spaced at 1500 m during the 19 years of the Phase I data taken between January 1, 2004 and January 1, 2023. After this successful operation, the Observatory was upgraded, within a project named AugerPrime, with the goal of improving the measurement of the mass composition of ultra-high energy cosmic rays [6]. The upgrade includes new scintillator and radio detectors added in each position of the Surface Detector Array. We also upgraded the detector electronics to improve the timing and signal resolution of water-Cherenkov detectors and extended their dynamic range by adding a new PMT, smaller than the previous ones. We installed underground muon detectors at the positions of the 433 m and 750 m Surface Detector Arrays. Phase II data-taking started on April 1, 2023. We show a preliminary spectrum built with events acquired during the first two years of Phase II.

2. Vertical and inclined spectra

The depth of the water-Cherenkov detectors grants them the geometric acceptance needed to measure particles arriving at large zenith angles. However, the structure of the shower changes significantly for more vertical and inclined events. While the more vertical showers are symmetric around their axis and have a significant contribution at ground level of photons and electrons, the more inclined showers are composed predominantly of muons, and their axial symmetry is broken due to the geomagnetic field and other effects. Due to these reasons, we reconstruct the events arriving with a zenith angle up to 60° , denominated here as *vertical* events, with a different method than those arriving between 60° and 80° , the *inclined* events.

We reconstruct vertical events using a function that describes the fall-off of detector signals with the distance to the shower axis [7]. The value of this function at a distance of 1000 m ($S(1000)$) provides a measure of the shower size. We estimated a shower size independent of the zenith angle (S_{38}), equivalent to the size the shower would have had if it had arrived at a zenith angle of 38° , via an attenuation function, $S_{38} = S(1000)/f_{\text{att}}$, calculated with the Constant Intensity Cut method [8]. The attenuation function f_{att} depends on the zenith angle through a variable $x = \sin^2 \theta - \sin^2 38^\circ$ and on S_{38} through a variable $y = \log_{10}(S_{38}/S_0)$, with $S_0 = 40$ VEM, and

$$f_{\text{att}}(x, y) = 1 + \sum_{i=1}^3 \sum_{j=0}^2 a_{ij} x^i y^j, \quad (1)$$

with parameters a_{ij} in the Table 1.

Table 1: Coefficients of the attenuation function used to correct the shower size of vertical events.

	1	y	y^2
x	-0.936	-0.005	0.4
x^2	-1.62	-0.51	-0.13
x^3	0.92	-0.54	-1.75

The shower size S_{38} is calibrated with the energy measured with the telescopes of the Fluorescence Detector (E_{FD}) using high-quality events observed in coincidence with the Surface Detector Array. From calibration data, we establish the relationship between the cosmic ray energy and the shower size, $E = AS_{38}^B$, with $A = (186 \pm 3)$ PeV, $B = 1.021 \pm 0.004$, and a correlation coefficient between A and B , $\rho = -0.98$.

We reconstruct inclined events by fitting a function that describes the signal pattern [9, 10]. This map is scaled with the ratio of the measured and simulated shower size at 10 EeV (N_{19}). The parameter N_{19} attenuates with the zenith angle in a slightly model-dependent way, in particular, due to the residual contribution of photons and electrons. This attenuation is corrected using the Constant Intensity Cut method as in the case of vertical events to derive an equivalent shower size at 68° , $N_{68} = N_{19}/f_{att}$. The attenuation function is described as a simplified model of Eq. 1, with $x = \sin^2 \theta - \sin^2 68^\circ$, $y = \log_{10}(N_{19})$, and $f_{att} = 1 + (0.292 - 0.468 y)x + (-4.96 + 0.79 y)x^2$. The shower size N_{68} is calibrated in energy using events observed in coincidence with the Fluorescence Detector as with vertical events. The energy of inclined events is calculated from $E = AN_{68}^B$, with $A = (5.29 \pm 0.06)$ EeV, $B = 1.046 \pm 0.014$, and a correlation coefficient $\rho_{AB} = -0.66$.

We calibrated the events using a Fluorescence Detector reconstruction updated with an improved estimation of aerosol attenuation and parametrization of the longitudinal profile of light emission [11]. We calibrated the vertical and inclined events with the same fluorescence reconstruction so that both datasets share the same energy scale. The systematic uncertainty of the energy of vertical and inclined events is driven by the absolute energy scale of fluorescence detection (14%) [12] that is correlated between both data sets. In addition, the uncertainty of the calibration parameters propagates as another source of systematic uncertainty in the energy but is not correlated between vertical and inclined events. Although the number of events available in the vertical calibration renders this source of systematic uncertainty negligible, inclined energies with fewer calibration events have a systematic uncertainty of about 2%.

For the spectrum based on vertical events, we selected events that have the detector with the highest signal surrounded by six working detectors to ensure the energy is well reconstructed. We applied a similar condition for inclined events, requiring the reconstructed core to be contained within a hexagon of six working detectors. We choose the energy thresholds of the vertical spectrum at $\log_{10}(E/\text{eV}) = 18.4$ ($E \sim 2.5$ EeV), and of the inclined spectrum at $\log_{10}(E/\text{eV}) = 18.6$ ($E \sim 4$ EeV), limits at which the trigger efficiency exceeds 97% [13]. Above the energy threshold, the exposure is constant and is calculated by monitoring the status of the detectors every 1 s. From this thorough monitoring, we identified periods of unstable acquisition that accounted for less than 3% of the observation time. We excluded the events observed during these periods and penalized the exposure accordingly. We computed an exposure of $(81\,100 \pm 700)$ km² sr yr for the vertical spectrum and $(23\,800 \pm 400)$ km² sr yr for the inclined spectrum.

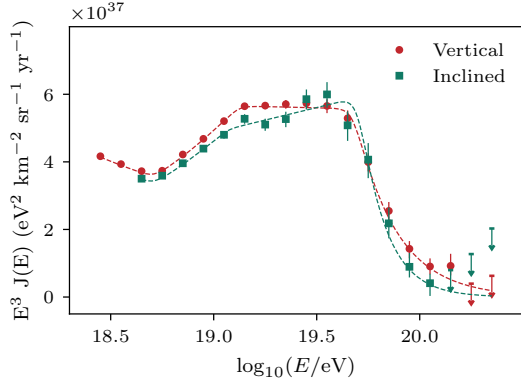


Figure 1: Vertical spectrum using events arriving with a zenith angle less than 60° , and inclined spectrum from events with a zenith angle between 60° and 80° . The spectrum data corrected for the detector response and the fitted flux are shown. No cut in the declination of the arrival direction has been applied.

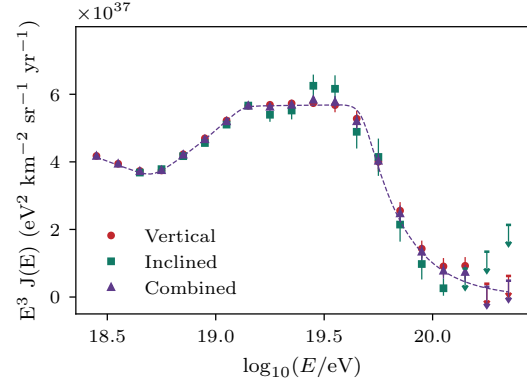


Figure 2: Vertical and inclined spectra in their common declination band $[-84.8^\circ, +24.8^\circ]$. The inclined spectrum is shifted according to the recalibrated event energies calculated in the spectra combination. The resulting combined spectrum is shown.

We fitted the spectrum data with a flux of cosmic rays modeled as a power-law function with smooth transitions,

$$J(E) = J_0 \left(\frac{E}{E_0} \right)^{-\gamma_0} \frac{\prod_{i=1}^3 \left[1 + \left(\frac{E}{E_i} \right)^{\omega_i^{-1}} \right]^{(\gamma_i - \gamma_{i+1}) \omega_i}}{\prod_{i=1}^3 \left[1 + \left(\frac{E_0}{E_i} \right)^{\omega_i^{-1}} \right]^{(\gamma_i - \gamma_{i+1}) \omega_i}}, \quad (2)$$

with the reference energy E_0 set at $10^{0.5} \sim 3.16$ EeV and the three transitions widths fixed at $w_i = 0.05$. We fitted the flux normalization J_0 , the ankle energy E_1 , the instep energy E_2 , the fall-off energy E_3 , and the four spectral indexes γ_i . The spectral index γ_1 corresponds to the region before the ankle, γ_2 between the ankle and the instep, and likewise for the other spectral indexes.

The observed number of events in each energy bin is affected, besides the prediction arising from the cosmic ray flux, by the effect of the detector response that includes the trigger efficiency, and the bias and finite resolution of the reconstructed energy [3]. We incorporated these effects in the spectrum fitting. We correspondingly corrected the observed flux in each bin by factors $c_i \sim [0.9 - 1]$, $J'_i = c_i J_i$, with $J_i = N_i / (\varepsilon \Delta E_i)$, with N_i the observed number of events, and ΔE_i the energy difference between the bin limits. The vertical and inclined spectra corrected by the detector response and the corresponding fitted fluxes are shown in Fig. 1.

3. Dependency of the flux with declination

The latitude of the Observatory at 35.2° S permits the observation of cosmic rays arriving with a declination from the south celestial pole up to $+24.8^\circ$ with vertical events, and the range $[-84.8^\circ, +44.8^\circ]$ with inclined events. We exploited the cumulative exposure of the vertical and

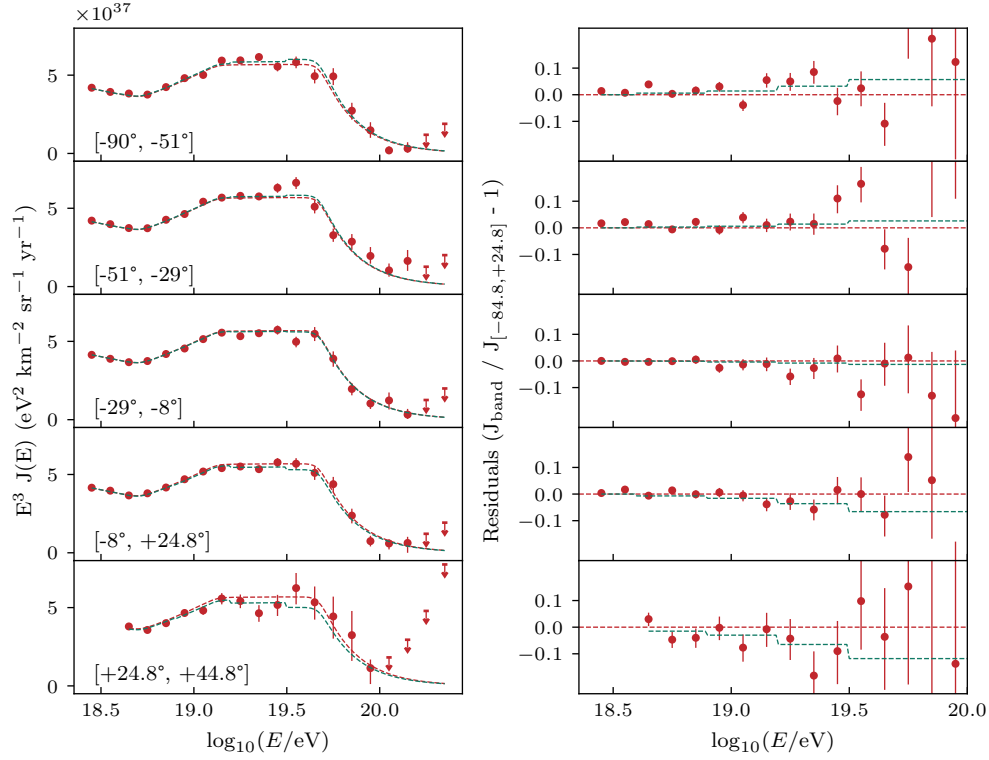


Figure 3: Left panel: Spectra in five declination bands combining vertical and inclined events. The flux fitted in the common declination band $[-84.8^\circ, +24.8^\circ]$ is shown in red, and the flux modulated by the dipole of the arrival direction distribution is shown in green. Right panel: Residuals of the combined spectra and the fitted flux in each declination band clipped to $[-0.25, 0.25]$.

inclined spectra by combining them into a single spectrum. To guarantee the observation of the same sky, we combined them in their common field of view given by the declination band $[-84.8^\circ, +24.8^\circ]$ using the method in reference [14]. We fitted the vertical and inclined data simultaneously with the flux model (2).

During the fit, we included the effect of the uncorrelated systematics on the energy of the vertical and inclined events. We fixed the energies resulting from the calibration for the vertical events as they have negligible systematics. However, we fluctuated the energy of inclined events by varying the parameters A and B around the calibration estimates. We use different B deviations for energies below and above 10 EeV, and denominated them δB and δC , respectively, to account for possible deviations of the inclined calibration from a pure power-law model originated, for example, by the evolution of the primary composition. The combination maximizes a Poisson likelihood that contains contributions predicting the expected number of events in each bin of the vertical and inclined spectra. In addition, during the fit, we penalized the deviations from the calibration estimates (δA , δB , and δC).

We fitted simultaneously the eight spectral parameters and the three calibration parameters. The fitted deviations are $\delta A = (160 \pm 39)$ PeV, $\delta B = 0.003 \pm 0.016$, and $\delta C = -0.02 \pm 0.02$. The recalibration of inclined events increases the energy by 2.9% at 4 EeV, 3.1% at 10 EeV, and decreases it 1% at 100 EeV. For the combination we obtained a deviance $D = 40.5$, for which we

Table 2: Spectral parameters of Phase I spectrum including the energy of the ankle (E_1), the instep (E_2), and the flux suppression (E_3), and the spectral indexes γ .

Parameter	Value	σ_{stat}	σ_{sys}	Parameter	Value	σ_{stat}	σ_{sys}
J_0 ($\text{km}^{-2}\text{sr}^{-1}\text{yr}^{-1}\text{eV}^{-1}$)	1.269	0.003	0.40	γ_1	3.26	0.01	0.10
E_1 (EeV)	5.1	0.1	1.1	γ_2	2.51	0.03	0.05
E_2 (EeV)	13	1	2	γ_3	2.99	0.03	0.10
E_3 (EeV)	48	2	5	γ_4	5.4	0.2	0.1

estimated a p -value $\simeq 0.12$. We show in Fig. 2 the fitted flux and the combined spectrum calculated in each bin as,

$$J_i = \frac{c_i n_i + c'_i n'_i}{(\varepsilon + \varepsilon') \Delta E_i}, \quad (3)$$

with the unprimed and primed symbols corresponding to the vertical and inclined spectra.

We searched for a dependence of the spectrum with declination by dividing the sky into five declination bands. We divided the range observed with vertical events $[-90^\circ, +24.8^\circ]$ in four bands of similar exposure. In addition, we considered a northern band $[+24.8^\circ, +44.8^\circ]$ observed only with inclined events. We compared the spectra in each declination band with the flux in the common declination range $[-84.8^\circ, +24.8^\circ]$. We calculated the combined spectrum in each band as per Eq. 3 using the number of events and exposures of each band. We show the spectra in declination bands and the flux in the left panel of Fig. 3. We also show the flux expected in each band when the modulations given by the dipole in the arrival directions observed by Auger are considered. We show in the right panel of Fig. 3 the residuals of the spectra in declination bands with respect to the fitted flux and the deviations of the flux modulated by the dipole. The residuals follow the trend imprinted by the dipole between 4 EeV and 32 EeV, and statistical uncertainties dominate beyond 32 EeV. The agreement of the dipole modulation with the spectra in each band is expected since the anisotropy and spectrum data sets overlap considerably.

4. Phase I spectrum

Given that the spectra in the different declination bands agree with the flux fitted in the common declination band within statistical uncertainties, we combine the vertical and inclined into a single spectrum. While in section 3 we restricted the vertical and inclined spectra to their common declination band $[-84.8^\circ, +24.8^\circ]$, we now combine them at all observed declinations, reaching a total exposure of $104\,900\text{ km}^2\text{sr yr}$ corresponding to the sum of the exposure of the vertical and inclined datasets. The combined spectrum is calculated using Eq. 3 after shifting the inclined spectrum with the parameters calculated in section 3. We show the combined spectrum and the corresponding fit in Fig. 4, and the spectrum data are provided in [5]. The shaded band represents the systematic uncertainty of the flux, which is dominated by the systematic uncertainty of the energy scale. The flux exhibits the firmly established features of the ankle and suppression, as well as the instep, which we unveiled with a significance of 3.9σ in 2020 [3]. Table 2 contains the fitted spectral parameters and their statistical and systematic uncertainties.

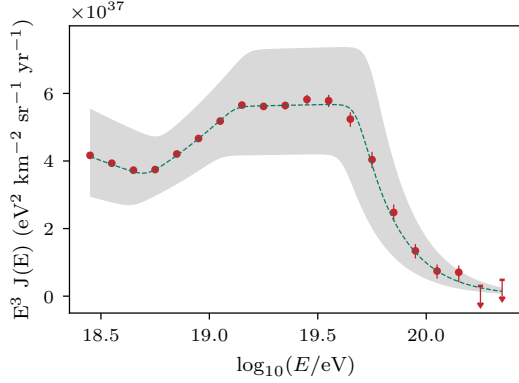


Figure 4: Combined spectrum using events arriving with zenith angle up to 80° observed at declinations from the south celestial pole up to $+44.8^\circ$. The systematic flux uncertainty is shown as a shaded band.

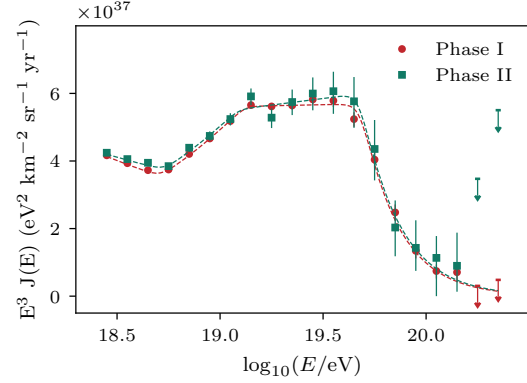


Figure 5: Preliminary spectrum using events arriving with zenith angle up to 60° observed during the second phase of Auger. The spectrum of Phase I is shown for comparison.

To assess the significance of the instep with the larger dataset currently available, we sampled energies from a reference model, which contained a slow suppression instead of the instep. We fitted the sampled data with the reference model and alternative one containing the instep and built a test statistic as the ratio of their respective likelihoods. The estimated significance is based on the number of times the sampled test statistic is greater than the one calculated by fitting the observed data. The test statistic was larger than the observed value ($t_{\text{obs}} \sim 35$) in only two out of 10^8 simulations, corresponding to a significance of 5.5σ .

5. Preliminary Phase II spectrum

In this section, we present a preliminary spectrum built from vertical events observed during the second Phase of the Auger Observatory, using data recorded from April 1, 2023, to March 1, 2025. To provide this first view of the Phase II spectrum, we reconstructed the events using the Phase I setup. For example, we applied the same lateral distribution function, attenuation with zenith angle, energy calibration, and detector response model used to unfold the spectrum. We also use the same method as in Phase I to select spectrum events, excluding periods of unstable data-taking, and to calculate the exposure. The exposure during the data-taking considered was $(9200 \pm 300) \text{ km}^2 \text{ sr yr}$, about 10% of the Phase I spectrum. Reusing the Phase I analysis was possible because the upgrade of the Observatory was designed so that water-Cherenkov detectors were backwards compatible.

We show in Fig. 5 the spectra of Phase I and II. We evaluated the consistency of Phase I and Phase II spectra with a Fisher's test that combines chi-square tests for bins with many events with exact binomial tests for bins with few events [15]. We assumed a systematic uncertainty of 1% uncorrelated between the two spectra, a conservative estimate given the 3% exposure uncertainty.

The observed Fisher's test statistic was $t = 38.0$ for 36 degrees of freedom, corresponding to a p-value $p = 0.30$ estimated numerically with Monte Carlo simulations.

We fitted the Phase II spectrum with the Phase I model, and obtained spectral parameters that are in agreement with the values in Table 2 within the statistical uncertainties. We will customize the event reconstruction to accommodate the improvements of the water-Cherenkov detectors in Phase II, but this early analysis already shows that minor changes to the existing setup will be required. In the future, we will use Phase I and II events to build a single spectrum with the combined exposure achieved during the complete lifetime of the Pierre Auger Observatory.

6. Conclusions

We presented an updated spectrum measured with the 1500 m Surface Detector Array of the Pierre Auger Observatory using events arriving with zenith angles up to 80° . The data set comprises 19 years of Phase I data taken from 2004 to 2023 with an exposure of $(104\,900 \pm 3000) \text{ km}^2 \text{ sr yr}$, enabling the discovery-level observation of the spectrum in step at $(13 \pm 1 \pm 2) \text{ EeV}$ with the spectral index increasing from $2.51 \pm 0.03^{\text{stat}} \pm 0.05^{\text{sys}}$ to $2.99 \pm 0.03^{\text{stat}} \pm 0.10^{\text{sys}}$. We did not find any statistically significant dependence of the flux with declination from the south celestial pole up to a $+44.8^\circ$, but a small trend consistent with the well-established dipolar anisotropy in arrival directions. Finally, we presented for the first time a spectrum obtained with events measured during Phase II of the Auger Observatory, showing its consistency with the Phase I spectrum.

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