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Exploring key electronic and non-linear optical amplitude with bilateral intramolecular charge transfer into thieno[3,2-*b*]thiophene-based compounds: a DFT approach†

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Organic molecules are considered important NLO materials in the modern age because of their potential optoelectronic features. Therefore, a series of organic molecules (DBTD1–DBTD7) with a D– π –A architecture was designed from the reference compound (DBTR) by structurally tailoring it with effective donors and making significant modifications to the π -spacers. All derivatives and references were initially optimized at the M06 functional with a 6-311G(d,p) basis set. Furthermore, the optimized structures (DBTR and DBTD1–DBTD7) were used to determine the density of states (DOS), frontier molecular orbitals (FMOs), natural bond orbital (NBO), UV-Visible, transition density matrix (TDM) analyses, and the most significant NLO properties. Among the compounds, DBTD4 had the smallest band gap (2.882 eV) which was further supported by DOS analysis. The global reactivity parameters were also related to the HOMO–LUMO band gap, where the compound with the lowest band gap showed a lower hardness value and higher softness value. NBO analysis was used to explain the molecular stability and hyperconjugation. DBTD4 and DBTD5 demonstrated comparable low bandgaps with the highest comparable NLO parameter values. However, we also observed efficient NLO characteristics for DBTD6 and DBTD7. The highest μ_{tot} value is observed in DBTD6 as 10.362 D, whereas the highest α_{tot} (1.48×10^{-22} and 1.47×10^{-22} esu) is observed in DBTD6 and DBTD7. Further, β_{tot} (6.68×10^{-28} and 6.23×10^{-28} esu) in DBTD4 and DBTD5 and γ_{tot} (6.20×10^{-33} and 6.59×10^{-33} esu) values are observed in DBTD5 and DBTD6. To acquire favorable NLO responses in molecules, structural modelling utilizing efficient donor units played significant role. Thus, current research insights encourage researchers to develop efficient NLO materials for optoelectronic applications.

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Introduction

Nonlinear optics is an emerging topic in physical sciences that explores the interaction of intense light with materials. It has significant scientific and technological implications.^{1,2} NLO materials and their application in photometric materials have been important in the past several decades.³ Non-linear optic

studies first, second, and third-order characteristics such as photorefractive, optoelectronic, and frequency doubling. The development of NLO substances is an area of experimental and theoretical investigations.⁴

Among these, understanding and finding new and better NLO materials to fulfil the needs of advanced technologies is the most important research field. Some years ago, several manufactured polymers, coloured compounds, and inorganic or organic diodes were semi-conductors.⁵ Although each kind of NLO material has its merits and demerits. Currently, organic based non-linear optical materials are more popular in the field of NLO research due to high non-linear optical parameters and the versatility of structures.^{6,7} Recent findings show that organic chemicals are valued for their excellent non-linear optical characteristics, attributed to their application in contemporary technologies, and can be synthesized cost-effectively.⁸ Therefore, currently much experimental and theoretical attention has been paid to π -conjugated molecular systems, based on donor moiety (D) linked through a π -bridge to an acceptor moiety (A)

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due to their exceptional NLO response.^{9,10} The design strategy of NLO in the D- π -A arrangement is characterized as a push-pull mechanism, which further enhances the internal charge transfer. Theoretical and experimental investigations confirm that careful selection of the π -spacer, donor, and acceptor units plays a crucial role in tuning NLO efficiency.¹¹ The literature

reports indicate that π -conjugated moieties provide an efficient charge transfer pathway under an electric field.^{12,13}

Moreover, it is observed that incorporating π -conjugated linkers of an appropriate length enhances the NLO characteristics.^{7,14} By enhancing charge transfer between donor and acceptor units, non-centrosymmetric π -linker derivatives

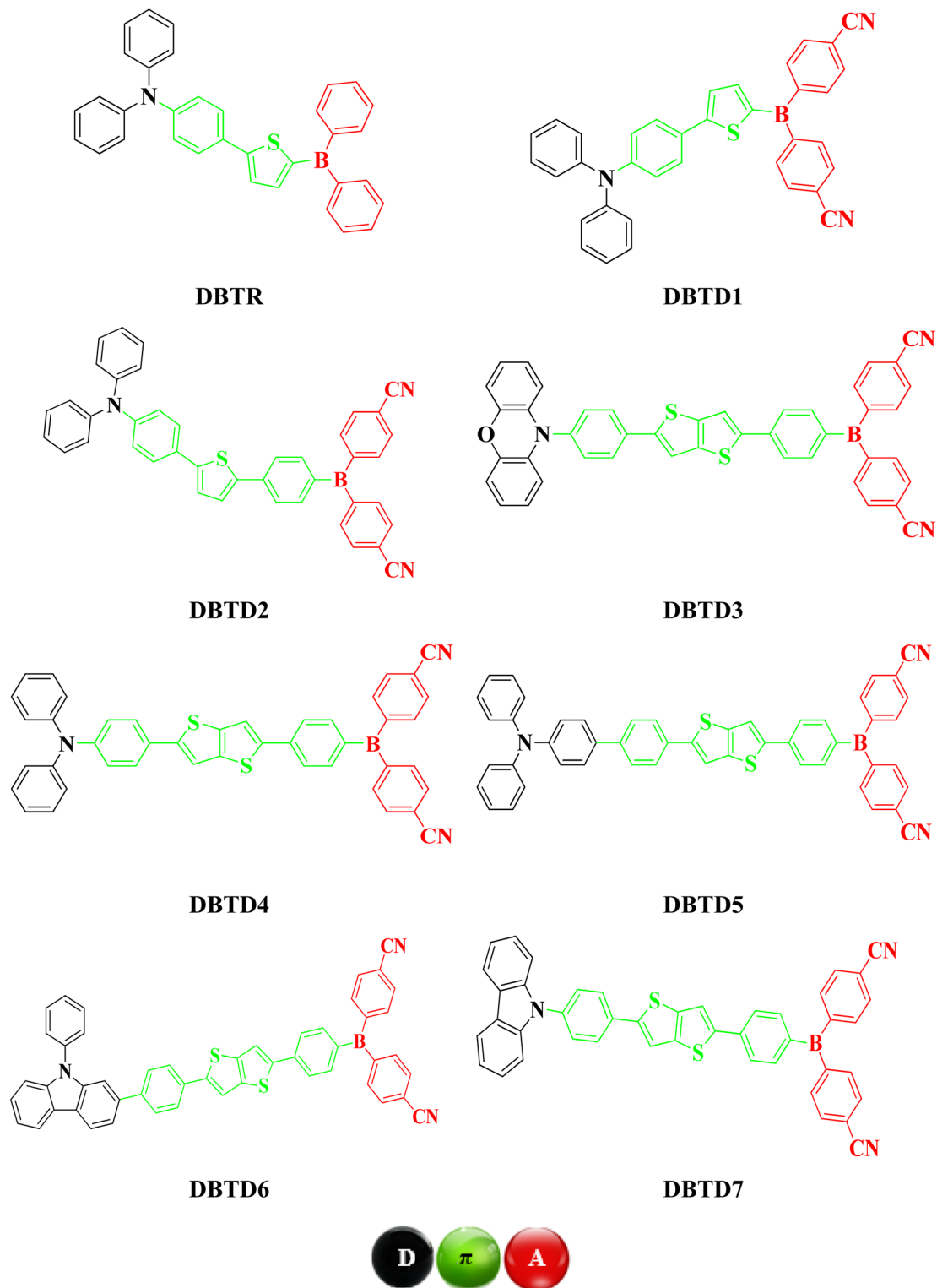


Fig. 1 Structures of reference and synthesized compounds (DBTR–DBTD7).



exhibiting first, second and third-order nonlinear polarizabilities can be achieved.^{15,16} To design the push-pull mechanism various schemes, for example, donor-acceptor, donor- π -acceptor, acceptor- π -donor- π -acceptor, and donor-donor- π -acceptor, are reported.¹⁷ Among them D- π -A framework, which facilitates photoinduced charge transfer, has been extensively explored in the design of high-performance NLO materials. At higher wavelengths, the push-pull mechanism can increase the asymmetric electronic distribution, improve the NLO response, lower the HOMO/LUMO band gap, and increase the light penetration range.^{18,19} The diphenylborane moiety enhances electron affinity and charge transfer process through its vacant p-orbital, lowering the LUMO level and boosting NLO and optoelectronic performance.²⁰ The thiophene-based π -spacers enhance charge transport due to their high conjugation efficiency and rigidity, which aid in stabilizing the charge-separated states.²¹ Similarly, phenoxazine, triphenylamine, and carbazole units, with their strong electron-donating capabilities, enhance charge asymmetry and boost hyperpolarizability.²²

Given these advantages, we designed a series of D- π -A derivatives, **DBTD1-DBTD7**, from a synthesized compound 4-(5-(diphenylboranyl)thiophen-2-yl)-*N,N*-diphenylaniline²³ utilized as reference compound for current study. This reference compound consists of a diphenylborane acceptor and a diphenylaniline donor, bridged by phenylthiophene π -spacers to promote efficient charge transfer. To further enhance intramolecular charge transfer (ICT), and electronic properties, the phenylthiophene unit in **DBTR** was replaced with tolyl thiophene and diphenyl thieno thiophene rings. Additionally, strategic donor modifications and central core adjustments were introduced to optimize π -conjugation and improve NLO efficiency.

This study explores the structure-property relationship of entitled compounds with distinct π -spacers and donor groups. A literature review revealed that their NLO properties remain unexplored. Our structural modifications significantly enhance the NLO response, making these compounds promising candidates for optoelectronic applications such as electro-optic switching, second-harmonic generation, and photonic devices. Furthermore, these findings provide

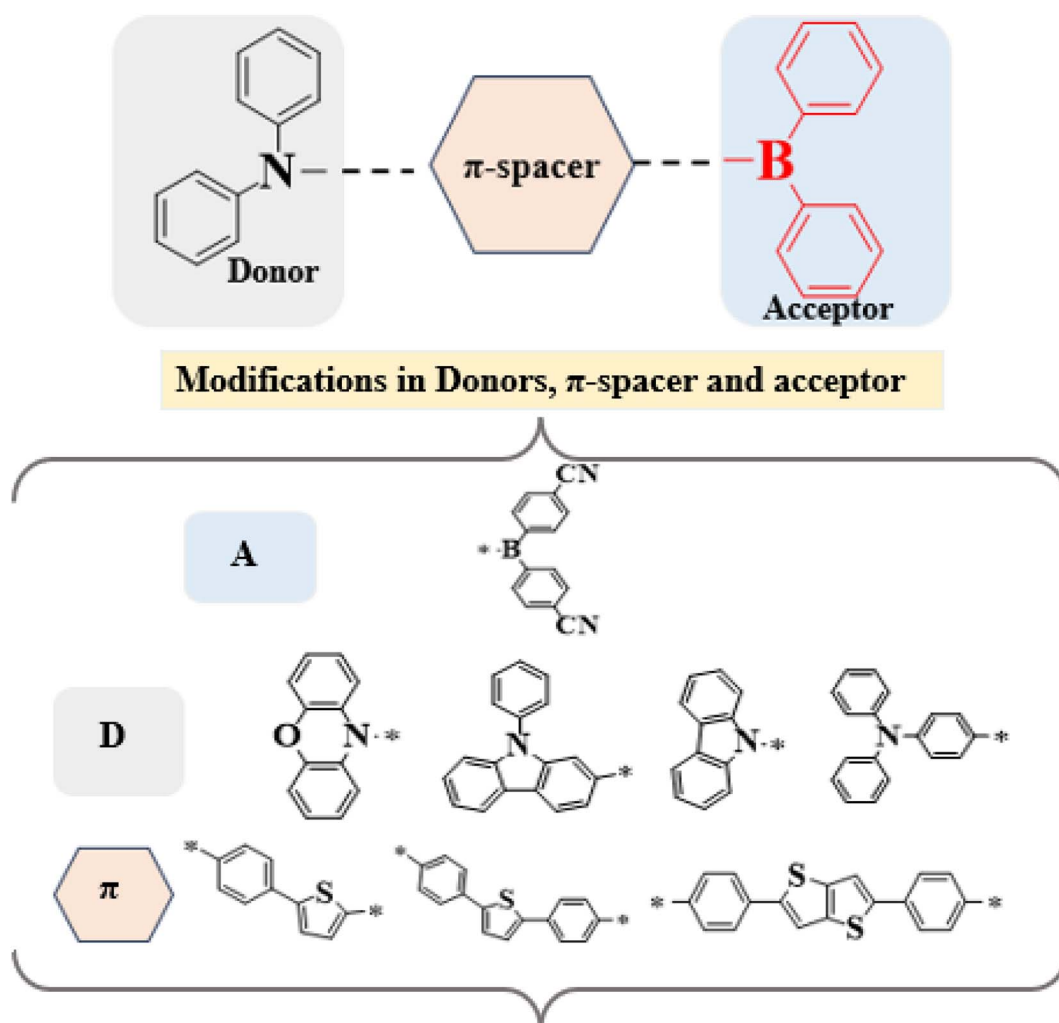


Fig. 2 Schematic representation of (DBTR-DBTD7).



valuable insights into the rational design of next-generation NLO materials, contributing to the development of more efficient and high-performance optoelectronic devices. The chemical structures and schematic representation of the reference and investigated compounds are presented in Fig. 1 and 2.

Computational procedure

New thieno[3,2-*b*]thiophene-based compounds were designed with different structural modifications *via* various donor moieties with the configuration of the D- π -A framework. All the quantum calculations were performed using the Gaussian 16 program.²⁴ The designed and reference compounds were optimized using the M06 functional²⁵ with the 6-311G(d,p) basis set.²⁶ Various electronic and optical properties were investigated using the same DFT functional and basis set. These include the density of states (DOS), frontier molecular orbital (FMO), transition density matrix (TDM), natural population (NPA), UV-Visible, natural bond orbital analyses (NBO), global reactivity parameter (GRPs), and the most important non-linear optical studies. DOS was performed using pyMolyze software²⁷ utilizing Gaussian log files, to investigate the electronic density relating to the FMO results. A natural bond orbital analysis was carried out to estimate the stabilization energy of the designed compound. UV-Visible absorption spectra in tetrahydrofuran solvent and gas were obtained to analyse the optical properties using Origin 8.0,²⁸ and GaussSum²⁹ software from Gaussian outputs. The solvent effect was investigated using the conductor-like polarizable continuum model (CPCM).³⁰ Using Multiwfn 3.7,³¹ we constructed a transition density matrix to investigate the electronic transitions. Chemcraft,³² Avogadro,³³ and Gaussview³⁴ were used for various type of data interpretation from Gaussian outputs. The convergence graphs of reference and designed compounds are shown in Fig. S6.†

Dipole moment³⁵ was calculated by using following eqn (1)

$$\mu = (\mu_x^2 + \mu_y^2 + \mu_z^2)^{\frac{1}{2}} \quad (1)$$

Average polarizability (α)³⁶ was examined using the eqn (2)

$$(\alpha) = 1/3(a_{xx} + a_{yy} + a_{zz}) \quad (2)$$

The magnitude of β_{tot} ³⁷ is defined in eqn (3).

$$\beta_{\text{tot}} = (\beta_x^2 + \beta_y^2 + \beta_z^2)^{\frac{1}{2}} \quad (3)$$

Eqn (4) was utilized to get the second hyperpolarizability (γ_{tot})³⁸

$$\gamma_{\text{total}} = \sqrt{\gamma_x^2 + \gamma_y^2 + \gamma_z^2} \quad (4)$$

where,

$$\gamma_i = \frac{1}{15} \sum_j (\gamma_{ijj} + \gamma_{jij} + \gamma_{jjj}) \quad i, j = \{x, y, z\}$$

Koopman's theorem³⁹ was used to calculate global reactivity descriptors using the following eqn (5)–(9):

$$X = \frac{[\text{IP} + \text{EA}]}{2} \quad (5)$$

$$\eta = \frac{[\text{IP} - \text{EA}]}{2} \quad (6)$$

$$\mu = \frac{E_{\text{HOMO}} + E_{\text{LUMO}}}{2} \quad (7)$$

$$\sigma = \frac{1}{2\eta} \quad (8)$$

$$\omega = \frac{\mu^2}{2\eta} \quad (9)$$

Results and discussions

The designed compounds comprised an acceptor, π -conjugated linkers, and a donor. The compound characteristics were modified by structural modifications, including adding donors and π -conjugated linkers. Generally, theoretical and experimental studies indicate that unifying strong donor and acceptor groups at the periphery of an adequately conjugated system may result in a significant NLO response.⁴⁰ The present study uses **DBTR** with IUPAC name, 4-(5-(diphenylboranyl)thiophen-2-yl)-*N,N*-diphenylaniline as reference compound with D- π -A configuration. **DBTR** underwent modification to design different derivatives. Literature shows that cyano moieties are widely used in FREAs due to their strong electron-withdrawing ability, which lowers FMOs energy levels and enhances charge transfer.⁴¹ Therefore, to improve the acceptor character of the molecule, **DBTD1** was designed, in which electrons with drawing cyano units were added to the terminal acceptors. To, further enhance the resonance, a benzene ring was integrated between the terminal acceptor and thiophene spacer in **DBTD2**. Additionally, conjugation was improved by substituting a single thiophene ring with fused double thiophene ring, leading to the design of **DBTD3**. Building on this structure, various donors were introduced at the left terminal of **DBTD3**, resulting in additional derivatives (**DBTD4**–**DBTD7**) with 4,4'-((4-(5(4(9-*H*-carbazol-9-yl)phenyl)thieno[3,2-*b*]thiophenyl)phenyl)boranediyl)dibenzonitrile as the central core. A double thiophene ring was incorporated as a π -spacer to further enhance the NLO characteristics of these derivatives. The following parameters for entitled compounds were calculated (i) linear polarizability, (ii) first hyperpolarizability, (iii) second hyperpolarizability, (iv) optical properties, and (v) NBO transitions and HOMO–LUMO energies through DFT/TD-DFT approaches. The optimized structures obtained *via* DFT analysis is shown in Fig. 3, respectively. Their Cartesian coordinates are presented in Tables S1–S8.† Moreover, the IUPAC names of examined compounds are given in Table S59.† Nevertheless, the current research provides theoretical insights into the design of



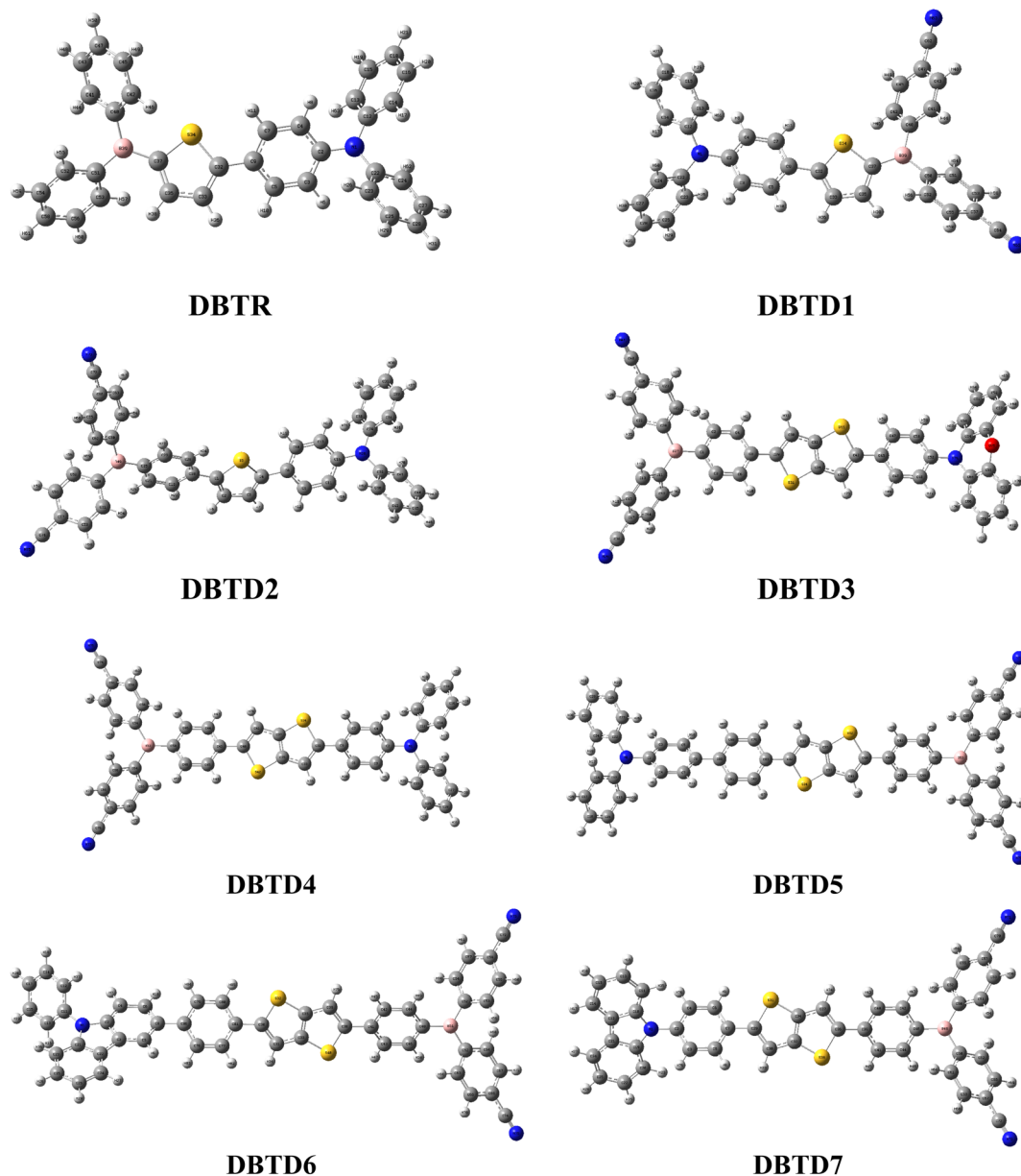


Fig. 3 Optimized structures of reference and designed compounds.

efficient NLO materials, but their practical applicability requires further experimental synthesis and characterization.

Frontier molecular orbitals (FMOs)

Frontier molecular orbitals (FMOs) analysis is a significant method for studying electronic transitions, molecular reactivities, chemical stability, and optical characteristics in organic systems.⁴² This parameter is crucial for determining the transfer of charge probability among compounds.⁴³ In a conjugated framework, the lowest unoccupied molecular orbital (LUMO) and the highest occupied molecular orbital (HOMO) determine the different interactions and reactions between molecules.^{44,45} According to band theory, valence and conduction bands are another way to describe the HOMO and LUMO of

compounds.^{46,47} HOMOs demonstrate the electron-donating capacity of nucleophiles, whereas LUMOs indicate the electron-accepting ability of electrophiles from nucleophiles.⁴⁸ Molecules with higher energy gap (ΔE) values are difficult because they resist changes in electronic configuration, resulting in low reactivity and high kinetic stability.⁴⁹ On the other hand, substances with low band gaps are thought of as soft molecular structures that are unstable and likely to give a reaction.⁵⁰ These substances exhibit efficient nonlinear optical properties and have high polarizability.⁵¹ Transferring electrons from HOMO to LUMO is crucial to comprehend a molecule's electronic properties and optical behavior.⁵²

The visual depiction of entitled chromophores depicts electronic distribution patterns in HOMO/LUMO during intramolecular charge transfer using red and blue areas in Fig. S1.†



In the same way, the pictographs of charge transferred in other orbitals (HOMO-1/LUMO+1 and HOMO-2/LUMO+2) are represented in Table S10 and Fig. S1.†

The energy findings of FMOs as illustrated in Table S9† represented that E_{HOMO} of reference compound and designed derivatives are calculated as -5.631, -5.711, -5.576, -5.662, -5.553, -5.583, -5.786 and -5.915 eV while E_{LUMO} values are found as -2.106, -2.613, -2.663, -2.667, -2.671, -2.690, -2.688 and -2.706 eV. We calculated the band gaps of the studied compounds using the HOMO-LUMO energies. The band gaps of DBTR-DBTD7 are computed as follows: 3.525, 3.098, 2.913, 2.995, 2.882, 2.893, 3.098 and 3.209 eV. Here, we observed that all the designed compounds showed smaller band gaps than the reference compound. The insertion of cyano units in derivatives, replacement of thiophene with thieno[3,2-*b*]thiophene, and enhanced resonance effect cause this

reduction and enhance charge transfer. **DBTD4** has the smallest band gap because of the introduction of a donor (*N*-methyl-*N*-phenylaniline) with the highest electron-donating effect. The data of FMOs along with molecular orbital energies also tells us about internal charge transfer between orbitals.⁵³ The compound **DBTD4** with lowest band gap has efficient charge transfer from donor to acceptor through π -bridge. The decreasing trend of band gaps for tailored compounds and reference are as follows: **DBTR** (3.525) > **DBTD7** (3.209) > **DBTD6** (3.098) > **DBTD1** (3.098) > **DBTD3** (2.995) > **DBTD2** (2.913) > **DBTD5** (2.893) > **DBTD4** (2.882) in eV. **DBTD6**, **DBTD1**, and **DBTD7** had the largest band gaps among all the derivatives. This might be due to a varied chemical framework and various associated donors with a lower electron-donating impact and hence does not cause effective charge transfer.

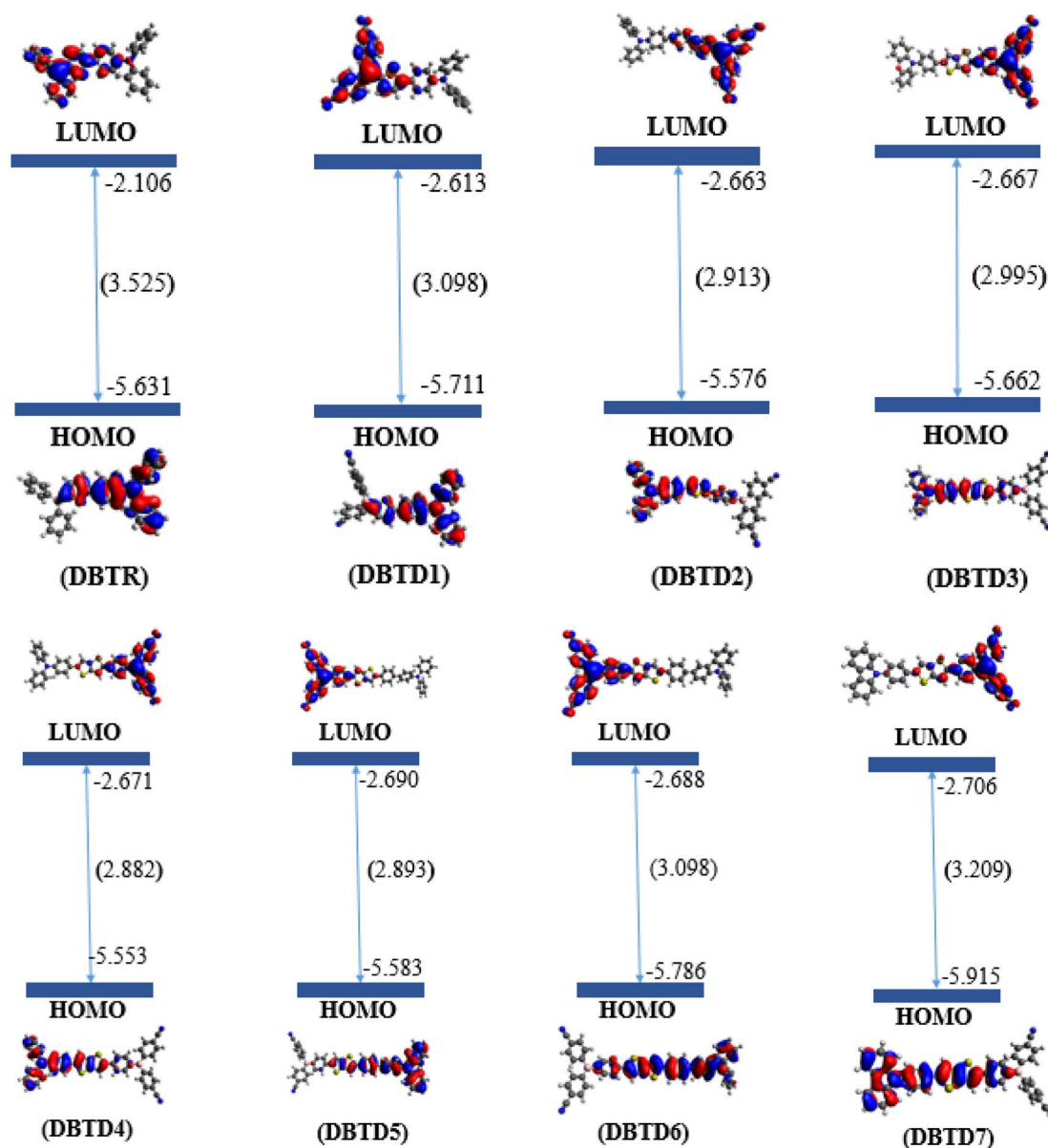


Fig. 4 Pictorial representation of band gap of designed derivatives.



The pictograms of FMO indicate the contribution of charge density. In **DBTR**, the electronic clouds related to the HOMO are mainly present on the donor and π -spacer. In contrast, the acceptor primarily contributes to the LUMO and a minor contribution from the π -spacer. However, a diverse charge density distribution was observed for the tailored compounds (**DBTD1–DBTD7**). The HOMOs of all the developed compounds demonstrated that their electronic clouds were primarily centred on the donor portion, with a modest contribution to the π -spacer.

On the other hand, the LUMO has focused electron clouds toward the acceptor part while leaving a tiny area on the π -spacer portion. This suggests that π -spacers facilitate the charge transfer between the donor and acceptor units. All investigated compounds may be considered significant NLO components owing to this charge flow. The pictorial representation of band gaps is presented in Fig. 4.

Density of states (DOS)

To support the FMOs data and explore the electron delocalization seen in the HOMOs and LUMOs of the compounds, a density of states (DOS) analysis was carried out.^{45,54} By measuring the density of states per unit of energy and volume, we can determine the number of electrons occupying a particular energy level. The energy distribution of electron scattering and the general distribution of energy levels concerning energy are disclosed through DOS calculations.⁵⁵ The π -conjugated bridge, which functions as a facilitator between electron-donating and electron-accepting moieties, is essential for improving the electron transfer processes within the molecule.⁵⁶ One important component of the band structure is the band gap, which significantly affects the optical and electrical properties of the compounds.⁵⁷ Left side peaks values in the graphical representation of DOS denote the HOMOs connected to the valence band. In contrast, right side values along the x -axis reflect the LUMOs connected to the conduction band.⁵⁸ The Fermi level, also referred to as the difference between the lowest and highest occupied single-electronic states, is the area with zero intensity between the HOMOs and LUMOs peaks. The Fermi level in n -type semiconductors is located near the conduction band (HOMO), but in p -type semiconductors, it is more likely to be located near the valence band.⁵⁹

To investigate the DOS, we further divided the reference compound and its derivatives into three segments: donor, π -spacer, and acceptor parts, represented by red, green, and blue lines, respectively. The density of states (DOS) examination at the M06 functional was conducted using the 6-311G(d,p) basis set to elucidate the findings of **DBTR** and **DBTD1–DBTD7**. The alteration of donors and π -spacers resulted in a change in the electronic distribution pattern. This transition can be explained very well by studying the percentages of moieties in HOMOs and LUMOs.⁶⁰ Regarding **DBTR**, the findings indicated that the donor moiety showed contributions of 57.4% towards HOMO and 2.4% towards LUMO. Similarly, the acceptor moiety displayed substantial contributions of 53.8% towards the LUMO and 2.3% towards the HOMO. The π -spacer contributed to the

LUMO by 43.7% and 40.3% towards the HOMO, respectively. In same way, in the case of **DBTD1–DBTD7**, the notable electronic charge distributions for the LUMO were 71.0, 74.1, 73.6, 73.4, 72.4, 72.8, 72.9% for the acceptor moiety. The electron density in the HOMO is distributed over the acceptor is noted as 2.5, 0.8, 1.1, 0.7, 0.3, 1.0, and 1.2%. Similarly, the charge contributions on the donors for the LUMO are as follows: 1.4, 0.3, 0.1, 0.2, 0.1, 0.0, 0.1%, and 27.7, 25.5, 26.3, 26.4, 27.6, 27.1, and 27.0% on π -spacer, respectively. Whereas in HOMO, the charge contributions on donor are 58.1, 49.5, 25.2, 42.9, 50.1, 17.8, 35.4% and 39.4, 18.7, 73.7, 56.4, 49.5, 81.2, 63.5% on π -spacer, accordingly. The acceptor significantly contributes to the LUMO. In the HOMO, all compounds displayed a significant contribution towards the donor, except for **DBTD3**, **DBTD4**, **DBTD6**, and **DBTD7**. However, these compounds have a major contribution to the π -spacers. The recorded images for studied compounds are presented in Fig. 5, and their percentages are presented in Table S11.†

Global reactivity parameters (GRPs)

The global reactivity parameters are calculated using the HOMO and LUMO energies, such as the ionization potential (IP),⁶¹ chemical potential (μ),⁶² electron affinity (EA),⁶³ electronegativity (X),⁶⁴ global electrophilicity index (ω),⁵¹ global hardness (η),⁶⁵ global softness (σ)⁶⁶ and charge transfer index (ΔN_{max}).⁶⁷ The chemical hardness [$\eta = -(E_{\text{LUMO}} - E_{\text{HOMO}})/2$] of a compound determines its activity and stability. It measures the resistance to changes in electron density or charge.⁶⁵ Similarly, the ability of an atom in a species to attract electrons to itself is known as electronegativity [$X = -(E_{\text{HOMO}} + E_{\text{LUMO}})/2$].⁶⁴ The values of GRPs are depicted in Table 1.

The association between the hardness (η), energy gaps, stability of a molecule, and chemical potential (μ) is direct. However, the relationship between reactivity and softness is inverse.⁴⁵

Based on the GRPs, we observed that the electron affinity values were higher in derivatives than the reference chromophore. The hardness values of **DBTR–DBTD7** are 1.762, 1.549, 1.456, 1.497, 1.441, 1.446, 1.549, and 1.604 eV. The softness values of the tailored compounds were greater than those of the reference compound: 0.284, 0.323, 0.343, 0.334, 0.347, 0.346, 0.323, and 0.312 eV⁻¹. Another parameter used to examine the stability and reactivity of the substances under study is the chemical potential (μ). It varies directly with stability and inversely with the compound reactivity. A D- π -A system is predicted to be more reactive and less stable because its μ values are more negative.⁶⁸ The chemical potential values are presented in Table 2. Moreover, the decreasing trend in hardness was **DBTR** > **DBTD7** > **DBTD6** = **DBTD1** > **DBTD3** > **DBTD2** > **DBTD5** > **DBTD4**. The order of increasing softness was **DBTR** < **DBTD7** < **DBTD6** = **DBTD1** < **DBTD3** < **DBTD2** < **DBTD5** < **DBTD4**. The values of the global electrophilicity index of **DBTR–DBTD7** were 4.246, 5.592, 5.826, 5.791, 5.867, 5.915, 5.795, and 5.790 eV, respectively. Since the compound (**DBTD4**) has the lowest hardness (1.441 eV) and the highest softness (0.347 eV⁻¹), which matches its lower energy band gap (2.882 eV), it is



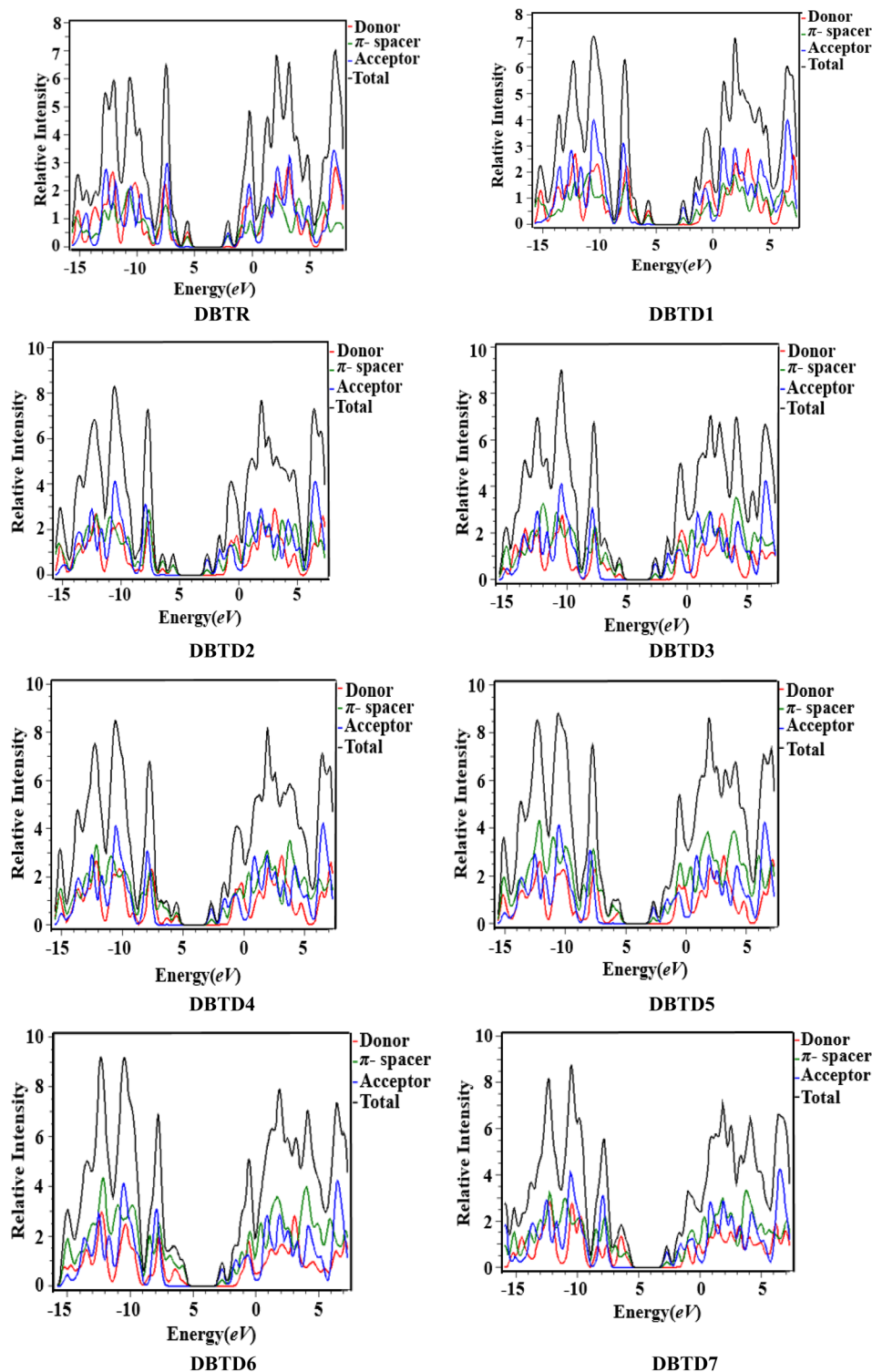


Fig. 5 DOS pictographs of reference compound and derivatives (DBTR–DBTD7).

expected to be the most advantageous designed compound with good NLO properties.

UV-Visible analysis

The nature of a molecular transition and its charge-transfer properties are understood by UV-Vis investigation.⁶⁹ A

conductor-like polarizable continuum model⁷⁰ is used to compute absorption bands to comprehend the system's optical characteristics in tetrahydrofuran solvent through DFT. This model considers the polarity of the solvent and its capacity to stabilize the $n \rightarrow \pi^*$ and $\pi \rightarrow \pi^*$ transition states at suitable energy levels, which generally advances our knowledge of



Table 1 Global reactivity parameters of the reference compound and its studied compounds (DBTR, DBTD1–DBTD7)^a

Compounds	IP	EA	X	η	μ	ω	σ	ΔN_{\max}
DBTR	5.631	2.106	3.868	1.762	−3.868	4.246	0.284	2.195
DBTD1	5.711	2.613	4.162	1.549	−4.162	5.592	0.323	2.687
DBTD2	5.576	2.663	4.119	1.456	−4.119	5.826	0.343	2.828
DBTD3	5.662	2.667	4.164	1.497	−4.164	5.791	0.334	2.781
DBTD4	5.553	2.671	4.112	1.441	−4.112	5.867	0.347	2.854
DBTD5	5.583	2.690	4.136	1.446	−4.136	5.915	0.346	2.859
DBTD6	5.786	2.688	4.237	1.549	−4.237	5.795	0.323	2.736
DBTD7	5.915	2.706	4.310	1.604	−4.310	5.790	0.312	2.686

^a Ionization potential (*I*), electron affinity (*A*), electronegativity (*X*), global hardness (η), chemical potential (μ), global electrophilicity (ω), global softness (σ), and maximum charge transfer index (ΔN_{\max}).

Table 2 Wavelength (λ), excitation energy (*E*), oscillator strength (f_{os}), and nature of the molecular orbital contributions of DBTR and DBTD1–DBTD7 in the solvent phase^a

Comp.	DFT λ (nm)	<i>E</i> (eV)	f_{os}	MO contributions
DBTR	430.097	2.883	0.921	H → L (95%) H−1 → L (2%) H → L+1 (2%)
DBTD1	485.242	2.555	0.758	H → L (96%)
DBTD2	502.265	2.469	0.741	H → L (93%) H−1 → L (3%) H → L+1 (3%)
DBTD3	490.657	2.527	0.984	H → L (93%) H−1 → L (3%) H → L+1 (3%)
DBTD4	502.591	2.467	0.885	H → L (92%) H−1 → L (4%) H → L+1 (3%)
DBTD5	488.088	2.540	1.024	H−1 → L (14%) H → L (81%), H → L+1 (3%)
DBTD6	473.276	2.620	1.248	H → L (87%) H−1 → L (8%) H → L+1 (3%)
DBTD7	459.116	2.701	1.163	H → L (86%) H−1 → L (9%) H → L+1 (3%)

^a MO = molecular orbital, H = HOMO, L = LUMO, f_{os} = oscillator strength, DFT = density functional theory.

optimal behavior.⁷¹ The λ_{\max} value indicates the amount of photon energy required to excite an electron from the HOMO to the LUMO. The f_{os} reflects the potential of transition, and the *E* represents the energy required for a transition.⁷² Broader absorption at greater λ_{\max} , high f_{os} , and low excitation energy are expected to result in efficient intramolecular charge transfer (ICT).⁷³ UV-Visible analysis of the modified compounds was performed at the M06/6-311G(d,p) level in gaseous and tetrahydrofuran, to investigate their absorption shifts. The separate and combined absorption spectra in tetrahydrofuran solvent and gaseous phase are shown in Fig. S2–S4.†

The values of λ_{\max} , along with the associated aspects, such as transition energy, oscillation strength, and molecular contributions in the solvent and gaseous phases, are presented in Tables

2 and 3. In contrast, Tables S12–S27† indicate the lowest six transitions. The data show that all of the compounds under investigation have absorption levels in the visible region of the electromagnetic spectrum, except **DBTD2**, **DBTD3**, and **DBTD4** in tetrahydrofuran solvent. They displayed low absorption values in the ultraviolet region, which might be due to the presence of double thiophene, which can lead to increased steric hindrance, conjugational distortion, and complex geometry alterations. When the oscillation strength was preferred over the wavelength, the absorption maxima (λ_{\max}) in THF solvent were as follows: **DBTD5** (488.088) > **DBTD1** (485.242) > **DBTD6** (473.276) > **DBTD7** (459.116) > **DBTR** (430.097) > **DBTD4** (394.013) > **DBTD2** (387.693), and **DBTD3** (377.333) (Tables S12–S19†). Here, **DBTD5** has the highest wavelength, which might be due to the presence

Table 3 Wavelength (λ), excitation energy (*E*), oscillator strength (f_{os}), and nature of the molecular orbital contributions of DBTR and DBTD1–DBTD7 in the gaseous phase^a

Comp.	DFT λ (nm)	<i>E</i> (eV)	f_{os}	MO contributions
DBTR	413.391	2.999	0.836	H → L (95%)
DBTD1	472.555	2.624	0.667	H → L (97%)
DBTD2	506.203	2.449	0.600	H → L (96%)
DBTD3	496.056	2.499	0.810	H → L (96%) H−1 → L (2%)
DBTD4	509.468	2.434	0.710	H → L (95%) H−1 → L (3%)
DBTD5	502.754	2.466	0.667	H → L (90%) H−1 → L (7%)
DBTD6	481.380	2.576	0.988	H → L (90%) H−1 → L (6%)
DBTD7	463.579	2.675	0.867	H−1 → L (10%) H → L (86%) H → L+1 (2%)

^a MO = molecular orbital, H = HOMO, L = LUMO, f_{os} = oscillator strength, DFT = density functional theory.



of auxochrome, leading to a redshift with lowest excitation energy of 2.540 eV.

However, there are variations in the wavelength values if we favor wavelength over oscillation strength in the THF solvent. In this case, the absorption maxima in THF decreased in the following order in nm: **DBTD4** (502.591) > **DBTD2** (502.265) > **DBTD3** (490.657) > **DBTD5** (488.088) > **DBTD1** (485.242) > **DBTD6** (473.276) > **DBTD7** (473.276), and **DBTR** (459.116).

At the lowest excitation energy of 2.467 eV, the maximum absorption value was recorded for **DBTD4** at 502.591 nm. Because of its narrowest band gap, it can take less energy to transition, and more electrons might move into excited states, accounting for its highest wavelength.

All derivatives exhibited higher absorption values in the gaseous phase than the reference compound. The compounds exhibited absorption values in the range from 413.391 nm to 509.468 nm. The **DBTD4** had the highest wavelength with an excitation energy of 2.434 eV in the gas phase. This might be due to the development of a more efficient push-pull mechanism. These findings were obtained when we prioritized wavelengths above oscillation strength. The gaseous phase exhibited a declining trend in λ_{\max} values measured in nm, as follows: **DBTD4** (509.468) > **DBTD2** (506.203) > **DBTD5** (502.754) > **DBTD3** (496.056) > **DBTD6** (481.380) > **DBTD1** (472.555) > **DBTD7** (463.579), and **DBTR** (413.391). When we prefer oscillation frequency to wavelength, the proposed compounds exhibit absorption values ranging from 370 to 496 nm regarding oscillation strength preference. All the substances exhibited UV absorption values except for **DBTD3**, **DBTD7**, **DBTD1**, and **DBTR**. (Tables S20–S36†).

In summary, **DBTD4** had the strongest affinity for decreasing the transition energy, the smallest energy difference, and exhibited a bathochromic shift. These features make it a desirable optically active material for nonlinear optics applications.

Light-harvesting efficiency (LHE) is a crucial factor that determines the optical performance of the designed compounds. It can be evaluated using the oscillator strength (f_{os}) obtained from TD-DFT calculations. A higher LHE value indicates an improved photocurrent response of the compound. The LHE for the designed chromophores was calculated using eqn (10) (ref. 74) and the results are presented in Table S36.†

$$\text{LHE} = 1 - 10^{-f} \quad (10)$$

Oudar and Chemla introduced a two-state model³⁵ that has been widely used in the literature to analyze the NLO response, incorporating the key contributions of the ground and excited states in a sum-over-states approach. This model establishes a relationship between charge transfer transitions and second-order polarizability, which serves as the foundation for the push-pull architecture in designing high-performance NLO compounds.

$$\beta_{CT} = (\Delta\mu_{gm} \times f_{gm})/E_{gm}^3 \quad (11)$$

Eqn (11) represents the two-state model for calculating first hyperpolarizability (β_{CT}). In this model, β_{CT} depends on three

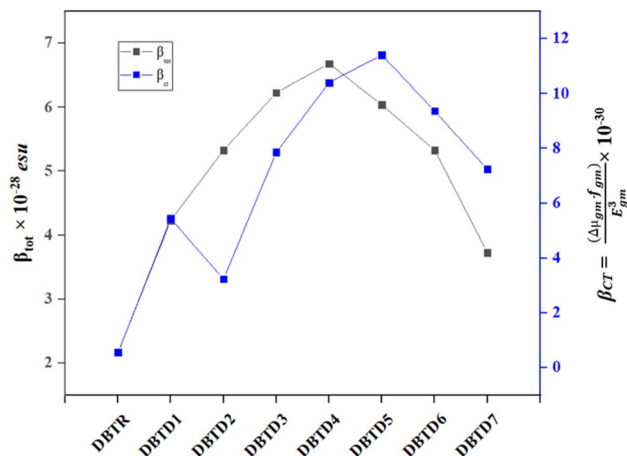


Fig. 6 Correlation between β_{tot} values and the corresponding β_{CT} values for the studied compounds.

key factors: the dipole moment difference between the ground and excited states ($\Delta\mu_{gm}$), the oscillator strength of the transition (f_{gm}), and the cube of the transition energy (E_{gm}^3). According to this model, β_{CT} is directly proportional to $\Delta\mu_{gm}$ and f_{gm} , meaning that a larger dipole moment difference and higher oscillator strength enhance the NLO response. Conversely, β_{CT} is inversely proportional to E_{gm}^3 , indicating that a lower transition energy leads to better NLO performance. Therefore, molecules with strong NLO properties should have a significant dipole moment difference, strong oscillator strength, and lower transition energy. In this study, these parameters for entitled compounds were determined and are presented in Table S36.†

The results indicate that among the studied compounds, reference compound exhibits the highest E_{gm}^3 value at 3.862 eV, while the lowest value is observed for **DBTD4** at 2.435 eV. This suggests that **DBTD4** compound might exhibit a stronger NLO response due to its lower transition energy. Additionally, to further analyze the NLO properties, the relationship between the two-state model (β_{CT}) and the total hyperpolarizability (β_{tot}) for all compounds (**DBTR**, **DBTD1–DBTD7**) is illustrated in Fig. 6. The variations in the graph reflect the influence of different π -spacers and terminal acceptor units on β_{tot} , which align with the trends predicted by the two-state model (β_{CT}), supporting their potential as efficient NLO materials.

Transition density matrix (TDM)

One widely used method for estimating electronic transitions in organic compounds is TDM analysis.^{75,76} In the excited state, TDM assists in determining the (a) donor-acceptor moiety interface, (b) electronic charge promotions, and (c) vitality of charge separation or charge compensation whereby electron-hole pairs are localized or delocalized.⁵⁸ It is possible to visually depict the distribution of charges using an electrostatic potential map.⁵⁶ Owing to minimal contribution in transition of hydrogen atoms their impact is neglected in these TDM maps. The M06/6-311G(d,p) functional was used for TDM analysis of the studied compounds. The investigated compounds were



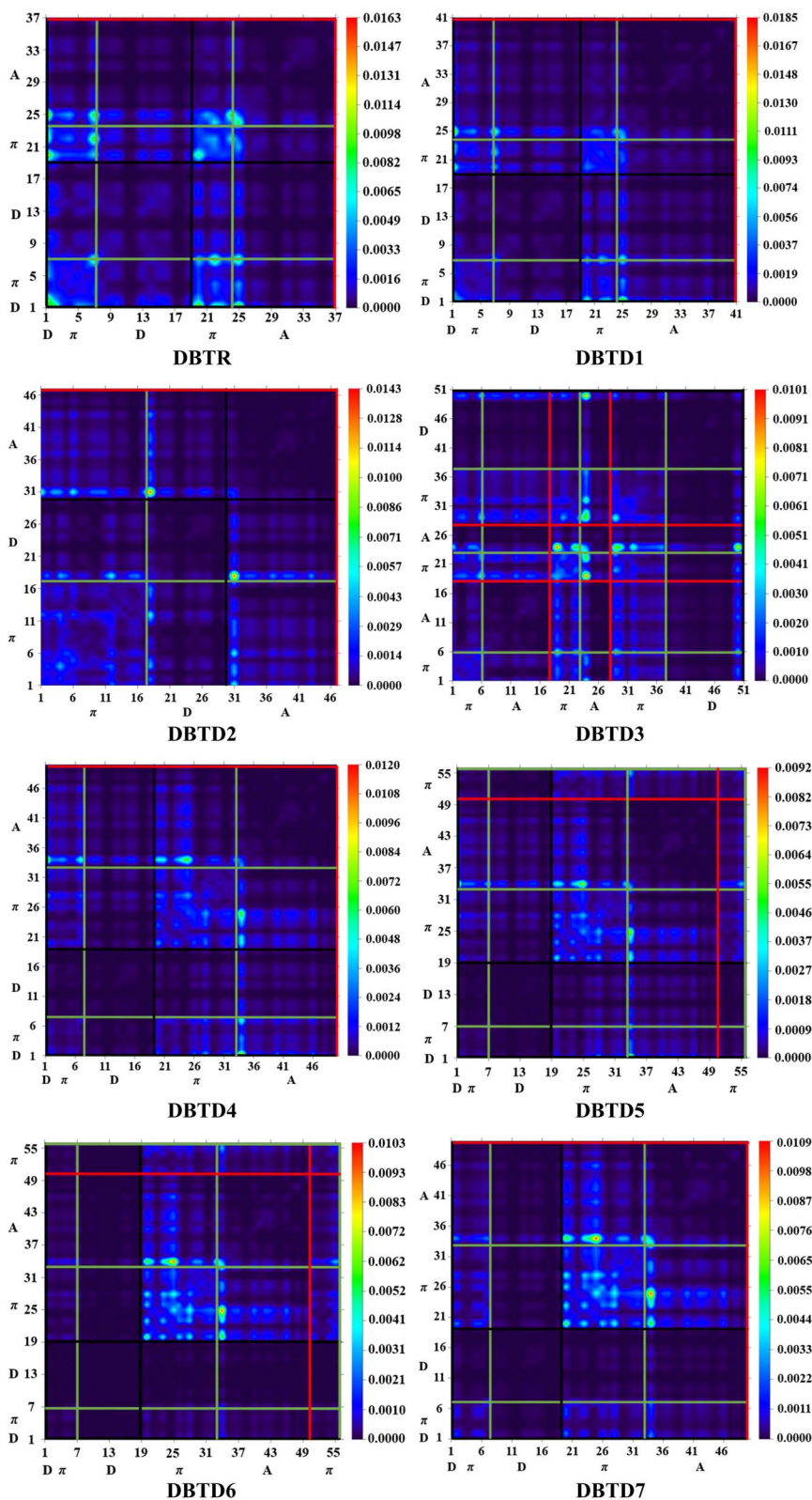


Fig. 7 TDM of reference compound DBTR and designed compounds (DBTD1–DBTD7).

divided into the following parts: donor, π -spacer, and acceptor. The TDM heat maps are shown in Fig. 7, where the colored dots represent the electron flow pattern. Here, the vertical and

horizontal lines define the electron coherence regions. The donor is represented by a black line, whereas the π -spacer and acceptor are represented by green and red lines, respectively.



The electron densities in the reference compound (**DBTR**) and derivatives (**DBTD1–DBTD3**) were present along the diagonal and off-diagonal portions. For further derivatives (**DBTD4–DBTD7**), electron densities are present along the diagonal section. We observed that charge was transferred from the donor to the acceptor through the π -bridge. The radiant portions in the TDM pictographs show that the electron density was concentrated at that spot. In the studied compounds, bright portions were mainly present on the acceptor and π -spacer and showed minor contribution towards the donor. This outcome is

important for the development of NLO materials. This TDM analysis also supports the FMO and DOS investigations.

Natural bond orbital (NBO) analysis

NBO analysis examines the electron-transfer pathway, bonding characteristics, bond-to-bond interactions, charge distribution, and hyper-conjugative association relationships.⁷⁷ Moreover, the intramolecular delocalization is precisely shown by NBO analysis.⁷⁸ To explain the uniform depiction of the donor– π –acceptor framework and the charge density shift from fully

Table 4 Selected natural bond orbital analysis of reference and designed compounds

Compounds	Donor(<i>i</i>)	Type	Acceptor(<i>j</i>)	Type	$E^{(2)}$ [kcal mol ⁻¹]	$E(j) - E(i)$ [a.u.]	$F(i,j)$ [a.u.]
DBTR	C2–C4	π	C7–C9	π^*	24.54	0.3	0.077
	C35–C37	π	C32–C33	π^*	15.88	0.3	0.064
	C37–B39	σ	C35–C37	σ^*	7.21	1.17	0.082
	S34–C37	σ	C33–C35	σ^*	0.52	1.22	0.023
	S34	LP (2)	C32–C33	π^*	27.12	0.27	0.078
DBTD1	N1	LP (1)	C22–C23	σ^*	3.55	0.85	0.052
	C50–C52	π	C55–C57	π^*	26.37	0.28	0.077
	C32–C33	π	C32–C33	π^*	0.55	0.3	0.012
	C57–C64	σ	C64–N65	σ^*	8.85	1.61	0.107
	C37–B39	σ	S34–C37	σ^*	0.52	0.8	0.018
DBTD2	S34	LP (2)	C32–C33	π^*	27.7	0.27	0.079
	N63	LP (1)	C47–C62	σ^*	11.87	1.06	0.1
	C18–C20	π	C23–C25	π^*	26.26	0.3	0.079
	C23–C25	π	C23–C25	π^*	0.64	0.29	0.012
	C66–C72	σ	C72–N73	σ^*	8.86	1.61	0.107
DBTD3	B49–C50	σ	C25–B49	σ^*	0.57	1.08	0.022
	N28	LP (1)	C13–C15	π^*	25.54	0.29	0.08
	N73	LP (1)	C66–C72	σ^*	11.87	1.06	0.1
	C11–C12	π	C14–C17	π^*	26.02	0.28	0.076
	C1–C3	π	C1–C3	π^*	0.64	0.29	0.012
DBTD4	C40–N41	σ	C26–C40	σ^*	8.2	1.58	0.102
	C42–S44	σ	C42–C43	σ^*	0.52	1.25	0.023
	N74	LP (1)	C50–C52	π^*	32.39	0.32	0.094
	N39	LP (1)	C17–C38	σ^*	11.87	1.06	0.1
	C43–C45	π	C48–C50	π^*	26.1	0.3	0.079
DBTD5	C48–C50	π	C48–C50	π^*	0.64	0.29	0.012
	C71–C76	σ	C76–N77	σ^*	8.86	1.61	0.107
	C32–S34	σ	C32–C33	σ^*	0.52	1.25	0.023
	N1	LP (1)	C2–C4	π^*	28.23	0.29	0.084
	N75	LP (1)	C61–C74	σ^*	11.87	1.06	0.1
DBTD6	C43–C45	π	C48–C50	π^*	25.9	0.3	0.079
	C35–C37	π	C40–C41	π^*	17.08	0.32	0.067
	C61–C74	σ	C74–N75	σ^*	8.86	1.61	0.107
	C81–H86	σ	C81–C82	σ^*	0.96	1.12	0.029
	S42	LP (2)	C35–C37	π^*	25.38	0.26	0.075
DBTD7	N75	LP (1)	C61–C74	σ^*	11.87	1.06	0.1
	C41–C43	π	C46–C48	π^*	25.88	0.3	0.079
	C21–C22	π	C21–C22	π^*	0.69	0.29	0.013
	C74–N75	σ	C69–C74	σ^*	8.21	1.58	0.102
	C30–S32	σ	C30–C31	σ^*	0.54	1.25	0.023
DBTD7	N1	LP (1)	C2–C3	π^*	36.4	0.31	0.097
	N73	LP (1)	C59–C72	σ^*	11.87	1.06	0.1
	C39–C41	π	C44–C46	π^*	25.71	0.3	0.079
	C44–C46	π	C44–C46	π^*	0.57	0.29	0.011
	C72–N73	σ	C67–C72	σ^*	8.21	1.58	0.102
DBTD7	C2–C4	σ	N1–C10	σ^*	0.5	1.16	0.022
	N1	LP (1)	C19–C20	π^*	40.32	0.31	0.101
	N71	LP (1)	C57–C70	σ^*	11.87	1.06	0.1



bound to half-filled non-bonded orbitals, the NBO analysis is useful.⁷⁹ Delocalization reactions are calculated using a second-order perturbation technique.⁸⁰ The following eqn (12) is used to determine the stabilizing energy $E^{(2)}$ that aids in delocalization $i \rightarrow j$, transition donor (i), and acceptor (j).⁸¹

$$E^2 = q_i \frac{(F_{ij})^2}{(\epsilon_j - \epsilon_i)} \quad (12)$$

Herein, $E^{(2)}$ suggests the stabilization energy, $F(i \rightarrow j)$ represents the diagonal, q_i represents donor orbital occupancy, and ϵ_j, ϵ_i denote the off-diagonal NBO Fock or Kohn–Sham mediums.⁸²

Different transitions are observed in the natural bond orbital analysis, such as $\pi \rightarrow \pi^*$, $\sigma \rightarrow \sigma^*$, $LP \rightarrow \pi^*$, and $LP \rightarrow \sigma^*$. The π – π^* transition is the most significant of all the transitions. Charge transfer in the $\pi \rightarrow \pi^*$ transition is more dominant than that in $\sigma \rightarrow \sigma^*$. The stabilization energy is the most important parameter to consider in NBO analysis. Table 4 illustrates the selected NBO values, whereas the other values are presented in Tables S37–S44.†

In the case of reference compound (**DBTR**), the value of the highest stabilization energy for π (C2–C4) $\rightarrow \pi^*$ (C2–C4) transition is 24.4 kcal mol^{−1} while the lowest energy for π (C35–C37) $\rightarrow \pi^*$ (C32–C33) is about 15.88 kcal mol^{−1}. The weak transitions $\sigma \rightarrow \sigma^*$ show greatest and lowest stabilization energy as 7.21 and 0.52 for σ (C37–B39) $\rightarrow \sigma^*$ (S35–C37) and σ (S34–C37) $\rightarrow \sigma^*$ (C33–C35). The lone pair transitions show values as 27.12 and 3.55 kcal mol^{−1} for LP (S34) $\rightarrow \pi^*$ (C32–C33) and LP (N1) $\rightarrow \sigma^*$ (C22–B23).

The stabilization energy values for **DBTD1** are 26.37, 0.55, 8.85, 0.52, 27.7, and 11.87 kcal mol^{−1} with corresponding transition as π (C50–C52) $\rightarrow \pi^*$ (C55–C57), π (C32–C33) $\rightarrow \pi^*$ (C32–C33), σ (C57–C64) $\rightarrow \sigma^*$ (C64–N65), σ (C37–B39) $\rightarrow \sigma^*$ (S34–C37), LP (S34) $\rightarrow \pi^*$ (C32–C33), LP (N63) $\rightarrow \sigma^*$ (C47–C62). We observed that **DBTD1** had a higher stabilization energy than the reference compound.

The designed compound (**DBTD2**) has a maximum stabilization energy of 26.26 kcal mol^{−1} for the π (C18–C20) $\rightarrow \pi^*$ (C23–C25) transition and a minimum of 0.64 kcal mol^{−1} for the π (C23–C25) $\rightarrow \pi^*$ (C23–C25) transition. The weak $\sigma \rightarrow \sigma^*$ transitions exhibit the highest and lowest stabilization energies of 8.86 and 0.57, respectively, for σ (C66–C72) $\rightarrow \sigma^*$ (C72–N73) and σ (B49–C50) $\rightarrow \sigma^*$ (C25–B49). The lone pair transitions for LP (N28) $\rightarrow \pi^*$ (C13–C15) and LP (N73) $\rightarrow \sigma^*$ (C66–C72) gave values of 25.54 and 11.87, respectively.

For the π (C11–C12) $\rightarrow \pi^*$ (C14–C17) transition, the proposed compound (**DBTD3**) had a maximum stabilization energy of 26.02 kcal mol^{−1}, whereas for the π (C1–C3) $\rightarrow \pi^*$ (C1–C3) transition, it had a minimum stabilization energy of 0.64 kcal mol^{−1}. The stabilization energies of the weak transitions $\sigma \rightarrow \sigma^*$ are the lowest at 0.52 and the highest at 8.2 for σ (C42–S44) $\rightarrow \sigma^*$ (C42–C43) and σ (C40–N41) $\rightarrow \sigma^*$ (C26–C40). For LP (N74) $\rightarrow \pi^*$ (C50–C52) and LP (N39) $\rightarrow \sigma^*$ (C17–C38), the lone pair transitions provided values of 32.39 and 2.34, respectively.

Regarding **DBTD4**, there is a minimum stabilization energy of 0.64 kcal mol^{−1} for the π (C48–C50) $\rightarrow \pi^*$ (C48–C50)

transition and a maximum stabilization energy of 26.1 kcal mol^{−1} for the π (C43–C45) $\rightarrow \pi^*$ (C48–C50) transition. When discussing weak transitions σ (C71–C76) $\rightarrow \sigma^*$ (C76–N77) and σ (C32–S34) $\rightarrow \sigma^*$ (C32–C33), the highest and lowest stabilization energies of the transitions $\sigma \rightarrow \sigma^*$ are 8.86 and 0.52 kcal mol^{−1}. Analysing further transitions for LP (N1) $\rightarrow \pi^*$ (C2–C4) and LP (S42) $\rightarrow \sigma^*$ (C33–C35), the values are 28.23 and 11.87 kcal mol^{−1}.

With respect to **DBTD5**, the π (C35–C37) $\rightarrow \pi^*$ (C40–C41) transition had a minimum stabilization energy of 17.08 kcal mol^{−1}, whereas the π (C43–C45) $\rightarrow \pi^*$ (C48–C50) transition had a maximum stabilization energy of 25.9 kcal mol^{−1}. The weak transitions σ (C61–C74) $\rightarrow \sigma^*$ (C74–N75) and σ (C81–H86) $\rightarrow \sigma^*$ (C81–C82) have the greatest and lowest stabilization energy, respectively, at 8.86 and 0.96 kcal mol^{−1}. The results of the analysis of lone pair transitions for LP (S42) $\rightarrow \pi^*$ (C35–C37) and LP (N75) $\rightarrow \sigma^*$ (C61–C74) are 25.38 and 11.87 kcal mol^{−1}, respectively.

A comparison of the stabilization energies for all the electronic transitions occurring in the **DBTD6** molecule shows that the energy transition π (C41–C43) $\rightarrow \pi^*$ (C46–C48) has the maximum energy stabilized at 25.88 kcal mol^{−1}. The transition π (C21–C22) $\rightarrow \pi^*$ (C21–C22) requires a small energy of only 0.69 kcal mol^{−1}. Quantitatively, we found that the maximum energy of $\sigma \rightarrow \sigma^*$ transition is 8.21 kcal mol^{−1} for (C74–N75) and (C69–C74). Conversely, the minimum energy in the case of σ (C30–S32) $\rightarrow \sigma^*$ (C30–C31) was 0.54 kcal mol^{−1}. It was noted that $LP \rightarrow \pi^*$ held the maximum value of 36.4 kcal mol^{−1} in LP (N1) and (C2–C3) transition, while a minimum value of 11.87 kcal mol^{−1} was found in LP (N73) $\rightarrow \sigma^*$ (C59–C72) transition.

Analysis of the **DBTD7** molecule showed that the π (C39–C41) $\rightarrow \pi^*$ (C44–C46) transitions have the largest stabilization energy of 25.71 kcal mol^{−1}, while the π (C44–C46) $\rightarrow \pi^*$ (C44–C46) transitions have the lowest value, 0.57 kcal mol^{−1}. The weak interaction transitions σ (C72–N73) $\rightarrow \sigma^*$ (C67–C72) and σ (C2–C4) $\rightarrow \sigma^*$ (N1–C10) have the highest and lowest stabilization energies of 8.21 and 0.5 kcal mol^{−1}, respectively. Furthermore, the LP (N71) $\rightarrow \sigma^*$ (C57–C70) and LP (N1) $\rightarrow \pi^*$ (C19–C20) transitions, which have stabilization energies of 11.87 and 40.32 kcal mol^{−1}, respectively, are caused by resonance effects in the molecule. For **DBTD7**, many other transitions are shown in Table S37.†

The above discussion indicates that the NBO findings of all the chemicals examined are interrelated. However, it should be noted that, unlike other designed compounds, **DBTD1** gave the highest stabilization energies. This could be due to the presence of cyano groups, as they often enhance electron delocalization, leading to molecular stability. The stability of the compound is also evidenced by the charges of the NBOs. We can observe that ICT and hyperconjugation play a vital role in developing materials with strong NLO characteristics.

Natural population analysis (NPA)

The distribution of charges on an atom significantly influences numerous characteristics of the chemical system, including



dipole moment, chemical reactivity, and electrostatic interactions between atoms and compounds. A more accurate interpretation of the conjugated system depends on calculating the atomic charges of the selected compounds.^{83,84} Fig. S5† shows the charges of **DBTR** and **DBTD1–DBTD7** calculated using the M06/6-311G(d,p) functional. To gain a better understanding of their molecular structures, it is important to calculate the atomic charge of the reference **DBTR** and the developed compounds (**DBTD1–DBTD7**), as boron, sulfur, and hydrogen atoms have a net positive charge in NBO analysis. However, because of steric effects, carbon primarily carries negative charges, while it may also carry positive charges because of its accepting and donating properties. In comparison to all other atoms in all the molecules under investigation (**DBTR–DBTD7**), nitrogen has the highest electronegativity. Thus, it exhibited the strongest intensity peak in the negative direction. There is an unequal distribution of charges when strongly electronegative elements are present.⁸⁵ NPA analysis is a useful method to gain information about reactivity since all of the derivatives have effective charge transfer mechanisms and are useful NLO materials. The graphs of the NPA investigation are depicted in Fig. S5.†

Non-linear optical (NLO) properties

The development of NLO research has increased following the emergence of lasers. It is considered the most promising technology for the advancement of biomedicine, photonics, and optoelectronics.^{86,87} The electrical characteristics of a compound, such as its polarizability (linear response) and hyperpolarizability (non-linear response), dictate its nonlinear features. Therefore, the linear and NLO potentials of materials are estimated by evaluating both polarizability and hyper-polarizability. In most cases, obtaining significant hyperpolarizability requires linear polarizability.⁸⁸ By altering the compound's non-centrosymmetric characteristics, the NLO qualities are improved.⁸⁹ A push–pull architecture that relies on the donor and acceptor, connected by a π -spacer, has been designed to establish the NLO response.⁹⁰ The NLO qualities are ascertained by quantum chemical approaches employing metrics such as the first hyper-polarizability (β_{tot}),³⁷ second hyper-polarizability (γ_{tot}), linear polarizability (α),³⁶ and dipole moment (μ_{tot}),³⁵ using eqn (1)–(4) mentioned earlier. To check the solvent effect on polarizability of investigated compounds the NLO analysis was performed in both the gas and solvent phase. The total values of the computed NLO parameters in tetrahydrofuran solvent and are presented in Tables S45 and S51,† respectively. However, in gaseous phase gas phase these values are presented in Tables S52–58,† respectively.

An essential component for determining the polarizability of organic compounds is the dipole moment (μ).⁹¹ For entitled molecules, their μ values are calculated in Debye (D), in solvent and gas phase shown in Tables S45 and S52,† respectively. As observed here, all developed compounds exhibited higher values than the reference compound. The highest μ_{tot} was found in **DBTD6** and **DBTD3** (10.362 and 10.223D, respectively)

in solvent phase because of the strong donor moieties, which may be due to their non-centrosymmetric structure. Urea was used as a reference compound for comparison analysis in NLO studies, according to the literature data its dipole moment is 1.3732D.¹⁶ The μ_{tot} in Debye for all other compounds and references in tetrahydrofuran are as the following: **DBTD1** (9.920D), **DBTD4** (9.610), **DBTD2** (9.090), **DBTD5** (8.828), **DBTD7** (5.402), **DBTR** (2.328). All the designed derivatives had higher values than urea. Higher values of μ_{tot} indicate high polarizability. For all the studied compounds, larger polarizability values were analyzed along the x -axis (**DBTD1** = -9.573 , **DBTD2** = 8.874 , **DBTD3** = 10.167 , **DBTD4** = 9.610 , **DBTD5** = -8.827 , **DBTD6** = -10.258 , and **DBTD7** = -5.401 in D).

For **DBTR** and **DBTD1–DBTD7**, Table S46† displays the $\langle\alpha\rangle$ values in solvent phase. It also reveals crucial 3-D tensors for polarizability. The linear polarizability value of **DBTR** was the lowest, and all other designed compounds had the highest values. Along the z -axis all compounds showed major contributing tensor in the $\langle\alpha_{\text{tot}}\rangle$ value compared with the x - and y -axes with comparable values. The $\langle\alpha\rangle$ values for entitled molecules in declining order are represented as follows in solvent phase: **DBTD5** (1.48×10^{-22}) > **DBTD6** (1.47×10^{-22}) > **DBTD4** (1.31×10^{-22}) > **DBTD7** (1.29×10^{-22}) = **DBTD3** (1.29×10^{-22}) > **DBTD2** (1.19×10^{-22}) > **DBTD1** (1.02×10^{-22}) and **DBTR** (0.916×10^{-22}) in esu. The computed linear polarizability in gas phase is presented in Table S53.†

According to the literature, the FMO bandgap has a distinct impact on the polarizability of compounds; that is, the narrower the bandgap, the higher the polarizability values, and *vice versa*.⁵⁰ The compounds to be studied in our case have a similar pattern: **DBTD4** has the highest β_{tot} value (6.68×10^{-28} esu) with the smallest band gap of 2.882 eV due to the extended conjugated system and strong push–pull framework. A decreasing trend in all entitled compounds in solvent phase is as follows: **DBTD4** (6.68×10^{-28}) > **DBTD3** (6.23×10^{-28}) > **DBTD5** (6.04×10^{-28}) > **DBTD2** (5.33×10^{-28}) = **DBTD6** (5.33×10^{-28}) > **DBTD1** (4.23×10^{-28}) > **DBTD7** (3.73×10^{-28}) and **DBTR** (2.17×10^{-28}) Table S47.† The calculated β_{tot} in gas phase of all examined compounds is presented in Table S55.†

In the second hyper-polarizability, the values of γ_{tot} are the highest for the **DBTD5** and **DBTD4** compounds for solvent phase, which have the lowest band gaps. All values of γ_{tot} were comparable to each other and greater than that of the reference compound (Table S48†). The order of declining trend of γ_{tot} in solvent medium in esu is following: **DBTD5** (6.59×10^{-33}) > **DBTD4** (6.20×10^{-33}) > **DBTD6** (5.32×10^{-33}) > **DBTD3** (5.07×10^{-33}) > **DBTD7** (3.54×10^{-33}) > **DBTD2** (4.42×10^{-33}) > **DBTD1** (2.75×10^{-33}) and **DBTR** (1.43×10^{-33}). The second hyper-polarizability of examined compounds in gas phase is presented in Table S56.†

We computed frequency-dependent first hyperpolarizability coefficients for electro-optic Pockel's effect (EOPE) with $\beta(\omega;\omega,0)$ and second-harmonic generation of first hyperpolarizability (SHG) with $\beta(-2\omega;\omega,\omega)$ at commonly used laser wavelengths of 532 nm and 1064 nm to give insights for experimental studies. At 532 nm, the EOPE values fall between 2.84×10^{-28} to 10.1×10^{-28} , whereas at 1064 nm, the response is in the range of 20.5



$\times 10^{-28}$ to 773×10^{-28} in solvent phase. We found significant EOPE values at 1064 nm, which might be attributed to an enhancement of resonance. Similarly, the SHG values are high for all derivatives except **DBTD5** and **DBTD6** at both 532 and 1064 nm as shown in the Table S49.† Therefore, it can be concluded that increased wavelength produces significant SHG responses.

At 532 nm and 1064 nm, third order nonlinear optical response coefficients, such as the electric field induced second harmonic generation (ESHG) $\gamma(-2\omega; \omega, \omega, 0)$ and the dc-Kerr effect $\gamma(-\omega; \omega, 0, 0)$ were also calculated. Higher dc-Kerr effect and ESHG values are mostly seen at 1064 nm, showing that the entitled compounds can significantly improve their performance at higher wavelengths as depicted in Table S50.† might be due to increased polarizability.

Overall, it has been noted that all derivatives have polarizable characteristics and a smaller band gap than **DBTR**. Among all derivatives, **DBTD4**, **DBTD5**, and **DBTD6** had the highest NLO values. Furthermore, these compounds have low hardness and high softness values, indicating that these are more polarizable. Therefore, the entitled compounds are considered as good NLO materials for optoelectronic devices. The frequency dependent second order and third order optical response in gas phase is illustrated in Tables S56–S57.†

Conclusion

A quantum chemical study was performed on the designed molecules (**DBTD1–DBTD7**) to clarify the relationship among the donor, π -spacer, and acceptor; all the designed compounds showed a notable reduction in the band gaps compared to the reference compound. The smallest band gap is obtained for **DBTD4**, which has the highest softness value, as illustrated in GRPs among all compounds (**DBTD1–DBTD7**). Similarly, their UV-Visible spectra were obtained in the range of bathochromic shifts with the lowest transition energies. Additionally, NBO investigations revealed the stabilization energies of the compounds. According to the TDM heat maps, the charge is transported from the donor to the acceptor *via* the π -spacer along the diagonal and off-diagonal portions. Moreover, it observed that the efficient NLO response is shown not only by **DBTD4** but also by **DBTD5** and **DBTD6**. The highest μ_{tot} value was observed for **DBTD6** at 10.362 D, whereas the highest values of α_{tot} , β_{tot} , and γ_{tot} were observed for **DBTD4** and **DBTD5**. Consequently, these compounds exhibited better NLO performance. Therefore, the above mentioned significant findings suggest that these compounds may be used in different application of photonic fields.

Data availability

All data generated or analyzed during this study are included in this published article and its ESI Files.†

Conflicts of interest

There is no conflict of interest.

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References

- 1 R. W. Boyd, *J. Opt. Soc. Am. B*, 2011, **28**, A38–A44.
- 2 J. Zyss, *Molecular Nonlinear Optics: Materials, Physics, and Devices*, Academic press, 2013.
- 3 A. Karakaş, A. Elmali, H. Ünver and I. Svoboda, *Spectrochim. Acta, Part A*, 2005, **61**, 2979–2987.
- 4 M. Akram, M. Adeel, M. Khalid, M. N. Tahir, M. U. Khan, M. A. Asghar, M. A. Ullah and M. Iqbal, *J. Mol. Struct.*, 2018, **1160**, 129–141.
- 5 A. Datta, *J. Phys. Chem. C*, 2009, **113**, 3339–3344.
- 6 S. Basu, *Ind. Eng. Chem. Prod. Res. Dev.*, 1984, **23**, 183–186.
- 7 S. Muhammad, *J. Mol. Graphics Modell.*, 2015, **59**, 14–20.
- 8 S. Muhammad, H. Xu, Z. Su, K. Fukuda, R. Kishi, Y. Shigeta and M. Nakano, *Dalton Trans.*, 2013, **42**, 15053–15062.
- 9 M. Panneerselvam, A. Kathiravan, R. V. Solomon and M. Jaccob, *Phys. Chem. Chem. Phys.*, 2017, **19**, 6153–6163.
- 10 A. Shanthi, C. Krishnan and P. Selvarajan, *J. Cryst. Growth*, 2014, **393**, 7–12.
- 11 R. Canton-Vitoria, Y. Sayed-Ahmad-Baraza, M. Pelaez-Fernandez, R. Arenal, C. Bittencourt, C. P. Ewels and N. Tagmatarchis, *npj 2D Mater. Appl.*, 2017, **1**, 13.
- 12 N. S. Abdel-Kader, S. A. Abdel-Latif, A. L. El-Ansary and A. G. Sayed, *J. Mol. Struct.*, 2021, **1223**, 129203.
- 13 M. R. S. A. Janjua, *J. Iran. Chem. Soc.*, 2017, **14**, 2041–2054.
- 14 M. R. S. A. Janjua, M. U. Khan, B. Bashir, M. A. Iqbal, Y. Song, S. A. R. Naqvi and Z. A. Khan, *Comput. Theor. Chem.*, 2012, **994**, 34–40.
- 15 M. Khalid, M. U. Khan, I. Shafiq, R. Hussain, K. Mahmood, A. Hussain, R. Jawaria, A. Hussain, M. Imran and M. A. Assiri, *Arabian J. Chem.*, 2021, **14**, 103295.
- 16 P. N. Prasad and D. J. Williams, *Introduction to Nonlinear Optical Effects in Molecules and Polymers*, Wiley, New York, 1991.
- 17 C. Yao, Y. Yang, L. Li, M. Bo, J. Zhang, C. Peng, Z. Huang and J. Wang, *J. Phys. Chem. C*, 2020, **124**, 23059–23068.
- 18 M. Haroon, R. Mahmood and M. R. S. A. Janjua, *J. Cluster Sci.*, 2017, **28**, 2693–2708.
- 19 M. R. S. A. Janjua, C.-G. Liu, W. Guan, J. Zhuang, S. Muhammad, L.-K. Yan and Z.-M. Su, *J. Phys. Chem. A*, 2009, **113**, 3576–3587.
- 20 S. Muhammad, M. R. S. A. Janjua and Z. Su, *J. Phys. Chem. C*, 2009, **113**, 12551–12557.



- 21 A. Ammasi and X.-H. Ju, *J. Phys. Chem. Solids*, 2024, **193**, 112155.
- 22 F. Liu, Y. Yang, H. Wang, J. Liu, C. Hu, F. Huo, S. Bo, Z. Zhen, X. Liu and L. Qiu, *Dyes Pigm.*, 2015, **120**, 347–356.
- 23 L. Wan, Z. Cheng, F. Liu and P. Lu, *Mater. Chem. Front.*, 2023, **7**, 4420–4444.
- 24 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov and D. J. Fox, *Gaussian 16, Revision A.03*, Gaussian, Inc., Wallingford, CT, 2016.
- 25 Y. Zhao and D. G. Truhlar, *Theor. Chem. Acc.*, 2008, **120**, 215–241.
- 26 M. P. Andersson and P. Uvdal, *J. Phys. Chem. A*, 2005, **109**, 2937–2941.
- 27 M. Khalid, I. Shafiq, S. C. Ojha, A. A. Braga, T. Ahamad and M. Arshad, *J. Photochem. Photobiol., A*, 2023, **445**, 115091.
- 28 K. J. Stevenson, *J. Am. Chem. Soc.*, 2011, **133**, 5621.
- 29 N. M. O'boyle, A. L. Tenderholt and K. M. Langner, *J. Comput. Chem.*, 2008, **29**, 839–845.
- 30 M. Cossi, N. Rega, G. Scalmani and V. Barone, *J. Comput. Chem.*, 2003, **24**, 669–681.
- 31 T. Lu and F. Chen, *J. Comput. Chem.*, 2012, **33**, 580–592.
- 32 G. A. Zhurko and D. A. Zhurko, *ChemCraft, version 1.6*, 2009, <https://www.chemcraftprog.com>.
- 33 M. D. Hanwell, D. E. Curtis, D. C. Lonie, T. Vandermeersch, E. Zurek and G. R. Hutchison, *J. Cheminf.*, 2012, **4**, 1–17.
- 34 R. Dennington, T. A. Keith, and J. M. Millam, *GaussView, Version 6*, Semichem Inc., Shawnee Mission, KS, 2016.
- 35 J.-L. Oudar and D. S. Chemla, *J. Chem. Phys.*, 1977, **66**, 2664–2668.
- 36 A. Alparone, *Chem. Phys.*, 2013, **410**, 90–98.
- 37 A. Plaquet, M. Guillaume, B. Champagne, F. Castet, L. Ducasse, J.-L. Pozzo and V. Rodriguez, *Phys. Chem. Chem. Phys.*, 2008, **10**, 6223–6232.
- 38 F. Ullah, K. Ayub and T. Mahmood, *New J. Chem.*, 2020, **44**, 9822–9829.
- 39 T. Koopmans, *Physica*, 1934, **1**, 104–113.
- 40 M. G. Papadopoulos, A. J. Sadlej and J. Leszczynski, *Non-linear Optical Properties of Matter*, Springer, 2006.
- 41 H. Wang, C. Zhao, Z. Burešová, F. Bureš and J. Liu, *J. Mater. Chem. A*, 2023, **11**, 3753–3770.
- 42 S. Uzun, Z. Esen, E. Koç, N. C. Usta and M. Ceylan, *J. Mol. Struct.*, 2019, **1178**, 450–457.
- 43 I. Shafiq, M. Khalid, M. Muneer, M. A. Asghar, R. Baby, S. Ahmed, T. Ahamad, S. F. de Alcântara Morais and A. A. Braga, *Mater. Chem. Phys.*, 2023, **308**, 128154.
- 44 M. N. Arshad, A.-A. M. Al-Dies, A. M. Asiri, M. Khalid, A. S. Birinji, K. A. Al-Amry and A. A. Braga, *J. Mol. Struct.*, 2017, **1141**, 142–156.
- 45 M. N. Tahir, M. Khalid, A. Islam, S. M. A. Mashhadi and A. A. Braga, *J. Mol. Struct.*, 2017, **1127**, 766–776.
- 46 M. Srncic and E. I. Solomon, *J. Am. Chem. Soc.*, 2017, **139**, 2396–2407.
- 47 M. U. Khan, M. Ibrahim, M. Khalid, A. A. C. Braga, S. Ahmed and A. Sultan, *J. Cluster Sci.*, 2019, **30**, 415–430.
- 48 M. Khalid, M. Ali, M. Aslam, S. H. Sumrera, M. U. Khan, N. Raza, N. Kumar and M. Imran, *Int. J. Pharm. Sci. Res.*, 2017, **8**, 457–469.
- 49 A. J. Garza, O. I. Osman, N. A. Wazzan, S. B. Khan, A. M. Asiri and G. E. Scuseria, *Theor. Chem. Acc.*, 2014, **133**, 1–8.
- 50 M. Khalid, H. M. Lodhi, M. U. Khan and M. Imran, *RSC Adv.*, 2021, **11**, 14237–14250.
- 51 P. K. Chattaraj and D. R. Roy, *Chem. Rev.*, 2007, **107**, PR46–PR74.
- 52 M. Khalid, I. Shafiq, K. Mahmood, R. Hussain, M. F. Ur Rehman, M. A. Assiri, M. Imran and M. S. Akram, *Sci. Rep.*, 2023, **13**, 1395.
- 53 A. Mahmood, M. I. Abdullah and M. F. Nazar, *Bull. Korean Chem. Soc.*, 2014, **35**, 1391–1396.
- 54 T. Hassan, R. Hussain, M. U. Khan, U. Habiba, Z. Irshad, M. Adnan and J. Lim, *Mater. Sci. Semicond. Process.*, 2022, **151**, 107010.
- 55 P. Goszeczycki, K. Stadnicka, M. Z. Brela, J. Grolik and K. Ostrowska, *J. Mol. Struct.*, 2017, **1146**, 337–346.
- 56 A. Mahmood, A. Irfan, F. Ahmad and M. R. S. A. Janjua, *Comput. Theor. Chem.*, 2021, **1204**, 113387.
- 57 A. Mahmood, M. I. Abdullah and S. U.-D. Khan, *Spectrochim. Acta, Part A*, 2015, **139**, 425–430.
- 58 M. Khalid, I. Shafiq, M. Zhu, M. U. Khan, Z. Shafiq, J. Iqbal, M. M. Alam, A. A. C. Braga and M. Imran, *J. Saudi Chem. Soc.*, 2021, **25**, 101305.
- 59 C. Zhao, C. G. Tang, Z.-L. Seah, Q.-M. Koh, L.-L. Chua, R.-Q. Png and P. K. Ho, *Nat. Commun.*, 2021, **12**, 2250.
- 60 M. N. Arshad, I. Shafiq, M. Khalid, M. Asad, A. M. Asiri, M. M. Alotaibi, A. A. Braga, A. Khan and K. A. Alamry, *Polymers*, 2023, **15**, 1508.
- 61 R. G. Parr and W. Yang, *J. Am. Chem. Soc.*, 1984, **106**, 4049–4050.
- 62 P. Politzer, in *Chemical Applications of Atomic and Molecular Electrostatic Potentials: Reactivity, Structure, Scattering, and Energetics of Organic, Inorganic, and Biological Systems*, Springer, 1981, pp. 7–28.
- 63 K. Fukui, *Science*, 1982, **218**, 747–754.
- 64 R. G. Parr, R. A. Donnelly, M. Levy and W. E. Palke, *J. Chem. Phys.*, 1978, **68**, 3801–3807.
- 65 R. G. Parr and R. G. Pearson, *J. Am. Chem. Soc.*, 1983, **105**, 7512–7516.
- 66 R. Parthasarathi, J. Padmanabhan, M. Elango, V. Subramanian and P. Chattaraj, *Chem. Phys. Lett.*, 2004, **394**, 225–230.
- 67 J. Padmanabhan, Z. Arif, P. Singh and R. Parthasarathi, *ACS Omega*, 2021, **6**, 21514–21524.
- 68 R. G. Pearson, *J. Mol. Struct.: THEOCHEM*, 1992, **255**, 261–270.
- 69 Z. Zhao, C. Xu, Y. Ma, K. Yang, M. Liu, X. Zhu, Z. Zhou, L. Shen, G. Yuan and F. Zhang, *Adv. Funct. Mater.*, 2022, **32**, 2203606.
- 70 V. Barone and M. Cossi, *J. Phys. Chem. A*, 1998, **102**, 1995–2001.



- 71 W. Rahmalia, J.-F. Fabre, T. Usman and Z. Mouloungui, *Spectrochim. Acta, Part A*, 2014, **131**, 455–460.
- 72 D. Liu, B. Kan, X. Ke, N. Zheng, Z. Xie, D. Lu and Y. Liu, *Adv. Energy Mater.*, 2018, **8**, 1801618.
- 73 R. Saleem, A. Farhat, R. A. Khera, P. Langer and J. Iqbal, *Comput. Theor. Chem.*, 2021, **1197**, 113154.
- 74 R. H. Xie, in *Handbook of Advanced Electronic and Photonic Materials and Devices*, Elsevier, 2001, pp. 267–307.
- 75 M. Chahkandi and H. A. R. Aliabad, *Chem. Phys.*, 2019, **525**, 110418.
- 76 M. Chahkandi, A. Keivanloo Shahrestanaki, M. Mirzaei, M. Nawaz Tahir and J. T. Mague, *Acta Crystallogr., Sect. B: Struct. Sci., Cryst. Eng. Mater.*, 2020, **76**, 591–603.
- 77 N. Sadlej-Sosnowska, *J. Org. Chem.*, 2001, **66**, 8737–8743.
- 78 R. T. Ulahannan, C. Y. Panicker, H. T. Varghese, R. Musiol, J. Jampilek, C. Van Alsenoy, J. A. War and S. Srivastava, *Spectrochim. Acta, Part A*, 2015, **151**, 184–197.
- 79 M. Szafran, A. Komasa and E. Bartoszak-Adamska, *J. Mol. Struct.*, 2007, **827**, 101–107.
- 80 M. Khalid, M. U. Khan, E.-T. Razia, Z. Shafiq, M. M. Alam, M. Imran and M. S. Akram, *Sci. Rep.*, 2021, **11**, 19931.
- 81 F. Weinhold and C. R. Landis, *Valency and Bonding: a Natural Bond Orbital Donor-Acceptor Perspective*, Cambridge University Press, 2005.
- 82 A. E. Reed, L. A. Curtiss and F. Weinhold, *Chem. Rev.*, 1988, **88**, 899–926.
- 83 Z. Demircioğlu, Ç. A. Kaştaş and O. Büyükgüngör, *Mol. Cryst. Liq. Cryst.*, 2017, **656**, 169–184.
- 84 Z. Demircioğlu, G. Kaştaş, Ç. A. Kaştaş and R. Frank, *J. Mol. Struct.*, 2019, **1191**, 129–137.
- 85 L. Li, C. Wu, Z. Wang, L. Zhao, Z. Li, C. Sun and T. Sun, *Spectrochim. Acta, Part A*, 2015, **136**, 338–346.
- 86 N. Z. Hakim, B. E. Saleh and M. C. Teich, *J. Lightwave Technol.*, 1991, **9**, 318–320.
- 87 D. F. Eaton, *Science*, 1991, **253**, 281–287.
- 88 C. Qin and A. E. Clark, *Chem. Phys. Lett.*, 2007, **438**, 26–30.
- 89 A. Mahmood, *J. Cluster Sci.*, 2019, **30**, 1123–1130.
- 90 M. Khalid, M. U. Khan, I. Shafiq, R. Hussain, A. Ali, M. Imran, A. A. Braga, M. Fayyaz ur Rehman and M. S. Akram, *R. Soc. Open Sci.*, 2021, **8**, 210570.
- 91 L. Kara Zaitri and S. M. Mekelleche, *Mol. Phys.*, 2020, **118**, 1618508.

